Flutter Evaluation and Control of an Airfoil Solved in the Laplace Domain

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Abstract

Aeroelasticity is concerned with the interaction of aerodynamic forces and the resulting structural deformations for a structure in an airflow. Flutter, which is one example of an aeroelastic phenomena, may result in the catastrophic failure of an aircraft. The traditional methods of predicting the flutter velocity, namely the U - g and p - k, have been used with great success. However, in recent years, new methods have been suggested which approximate the unsteady aerodynamic forces and moments by rational functions. The resulting equations are then solved in the Laplace domain. The goals of this work are twofold. Firstly, to solve the equations of motion for an airfoil, in both incompressible and transonic flow, in the Laplace domain using rational functions to approximate the unsteady aerodynamics. Secondly, to implement active control for the airfoil in incompressible flow, with the aim of increasing the critical flutter speed. The solution of the aeroelastic equations of motion in the Laplace domain proved to be a powerful tool in the analysis of flutter.

Sommaire

L'aéroélasticité traite de l'interaction des forces aérodynamiques et des déformations structurelles résultantes d'un corps dans un écoulement d'air. Le flottement, qui est un exemple de phénomène aéroélastique, peut provoquer d'importants dommages à un avion Les méthodes traditionnelles, U-g et p-k, qui sont capables de prédire la vitesse critique de flottement, ont été utilisées avec grand succès. Cependant, récemment, de nouvelles méthodes ont vu le jour, estimant les forces et moments instationnaires par des fonctions rationnelles. Les équations résultantes sont résolues dans le domaine de Laplace. Cette thèse a deux objectifs. Tout d'abord, résoudre les équations du mouvement pour un profil en écoulement incompressible et en écoulement transsonique dans le domaine de Laplace, en utilisant ces fonctions rationnelles. Ensuite, inclure un contrôle actif du profil en écoulement incompressible, dans le but d'augmenter le vitesse critique de flottement. On a démoutié aussi que la solution des équations aéroélastiques dans le domaine de Laplace représente un outil puissant pour l'analyse du flottement.

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Nomenclature

symbol

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a	non-dimensional distance measured from midchord to elastic
	axis (positive towards trailing edge)
a + ib	cigenvalue of form $\beta + iw$, a and b give damping and
	frequency, respectively
a _T	control of damping input to system by flap
a _m	numerator coefficients Pade approximate
$[\Lambda_j]$	linear coefficients of approximating rational function
b	semi-chord of airfoil
b _j	nonlinear lag terms of approximating rational function
b _l	denominator coefficients Pade approximate
С	chord
c + id	cigenvalue of form $\beta + iw$, c and d give damping and
	frequency, respectively
c _ð	non-dimensional distance measured from the midchord to
	the flap hinge line (positive towards trailing edge)
C(k)	exact Theodorsen function $F + iG$
$ar{C}(k)$	approximate Theodorsen function $ar{F}+iar{G}$
[C]	damping matrix (structural and/or aerodynamic)
C_{l_h}	lift coefficient due to plunge
$C_{l_{\alpha}}$	lift coefficient due to pitch
C_{m_h}	moment coefficient due to plunge

C_{m_0}	moment coefficient due to pitch
<i>c.g.</i>	center of gravity
D(p)	denominator of Pade approximant
e.a.	elastic axis
$err(\imath k)$	error function
F	real part
F_o,\ldots,F_m	coefficients of a m^{th} order polynomial
g	structural damping coefficient
G	imaginary part
h	plunge displacement
i	complex variable $\sqrt{-1}$
Ι	wing mass moment of inertia about center of gravity
I_{α}	wing mass moment of inertia about elastic axis
I_{β}	flap mass moment of inertia about flap hinge line
[<i>I</i>]	identity matrix
k	reduced frequency wb/U
k_{c}	reduced frequency wc/U
k _{cr}	flutter reduced frequency wb/U
Kh	spring constant in bending at elastic axis
Kα	spring constant in pitching at elastic axis
K_{β}	spring constant in pitching at flap hinge line
[K]	stiffness matrix (structural and/or aerodynamic)
L	lift force
т	mass of airfoil per unit span

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М	free stream Mach number
M_{α}	moment about elastic axis
M_B	moment about flap hinge line
[M]	inertia matrix (structural and/or aerodynamic)
M(ik)	normalizing function
[<i>N</i>]	approximation matrix, a function of reduced frequency
N(p)	numerator of Pade approximant
р	complex eigenvalue $\beta + iw$
Р	parameter gives nature of rational function
РD	proportional-derivative control
PID	proportional-integral-derivative control
q	dynamic pressure $1/2\rho U^2$
$q_{\rm c}$	control generalized coordinate parameter, i.e. eta
Q(ik)	tabular data of unsteady aerodynamics
Q(ik)	approximate data of unsteady aerodynamics
Q_h	total unsteady, nonconservative forces
Q_{lpha}	total unsteady, nonconservative moments
R	order-of-fit
r _a	radius of gyration about elastic axis
r_{β}	radius of gyration about flap hinge line
8	Laplace variable, <i>iw</i>
t	time variable
[T]	transfer function (1×2 matrix)
U	free stream velocity

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Ľ	non-dimensional velocity U/bw_{α}
.r	general Laplace solution multiplier (flow regime: incompressible
	or transonic)
.r.,	non-dimensional distance measured from elastic axis to center of mass
n	pitch angle
۵ ₀	complex constant pitch angle
3	flap angle
ζ	damping ratio $-\beta/\sqrt{\beta^2+w^2}$ ($p-k$ and Laplace methods)
γ	damping ratio $-g/2$ ($U - g$ method)
ζ_h	viscous damping for plunging motion
ζα	viscous damping for pitching motion
λ	flutter eigenvalue
λ_g	$\mu(1+ig)w_{\alpha}^2b^2/U^2$
μ	airfoil-air mass ratio $m/\pi \rho b^2$
ξ	non-dimensional displacement h/b
ξo	complex constant non-dimensional displacement h/b
ρ	air density
v	oscillation frequency
wh	uncoupled bending frequency
w_{lpha}	uncoupled torsional frequency
Wβ	uncoupled torsional frequency
w _r	reference frequency, normally taken as no control flutter frequency
[],	structural quantity
[]_c	control quantity

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1 Introduction

1.1 Motivation

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Aeroelasticity is concerned with the interaction of aerodynamic forces and the resulting structural deformations for a structure in an airflow. Flutter, which is one example of an aeroelastic phenomena, is a flow induced vibration that may result in the catastrophic failure of an aircraft. The traditional methods of predicting the flutter velocity, namely the l' - g and p - k methods, have been used with great success. However, in recent years, new methods have been suggested which approximate the unsteady aerodynamic forces and moments by rational functions. The resulting equations are then solved in the Laplace domain. This allows for a more efficient solution to the flutter problem, and also, allows for a direct extension of the aeroelastic equations for the implementation of active control.

The goals of this work are twofold. Firstly, to solve the equations of motion for an airfoil in both incompressible and transonic flow in the Laplace domain using rational functions to approximate the unsteady aerodynamics. These results will be compared with U - g and p - k generated results. Secondly, to implement active control for the airfoil in incompressible flow, with the aim of increasing the critical flutter speed.

1.2 Flutter Definition

In the field of acroelasticity (Fung 1955; Ashley, Bisplinghoff and Halfman 1955; Ashley and Bisplinghoff 1962), no topic has received more attention than that of flutter. Flutter is a violent self-induced acrodynamic vibration that may affect airplane wings, tail surfaces and control tabs. The following example will give a more physical representation of flutter. Consider a cantilever airfoil, with no sweepback and no aileron, rigidly attached to the side wall of a wind tunnel. The wing is then confined to exhibit motion only in the bending and pitching directions. With the air flow turned off, the airfoil is given an initial small amplitude disturbance, and then allowed to freely oscillate. The amplitude of the oscillation will decay until it is totally damped out. When the air flow is turned on, this original oscillation will decay much quicker. With subsequent increases in air speed the system will show improved decay characteristics; however, there exists a specific speed where the total system damping will begin to decrease. The air speed where the wing shows constant amplitude motion is referred to as the critical flutter speed. At speeds slightly greater, a small disturbance force will trigger violent oscillations, that in most cases lead to immediate catastrophic failure of the airfoil.

The dramatic and destructive consequences of flutter have gained the interest of much of the aeronautical community. One industry representative went so far as to state that: "The flutter problem is now generally accepted as a problem of primary concern in the design of current aircraft structures. Stiffness criteria based on flutter requirements are, in many instances, the critical design criteria." (Head 1958). In recent years, airplane design has resulted in the emergence of lighter, more flexible aircraft, with an increase in maximum speed, resulting in an increase in the likelihood of flutter occurring. The transonic regime tends to be the most critical with respect to flutter; however, there are still many other flutter problems that require attention.

Accordynamic structures are, by definition, continuous, having an infinite number of degrees-of-freedom. However, most theoretical analysis of such structures discretize the structure into a finite number of degrees-of-freedom, using, for example, either the finite element method or the assumed mode method. To obtain accurate estimates of the flutter velocity it is necessary to consider a large number of modes in the flutter analysis. However, it has been found that only two degrees-of-freedom, namely bending and pitching, are required to give a good physical representation of flutter. It is the coupling between these two modes that produces flutter. A rigid airfoil allowed only to have bending motion does not flutter. In general, a rigid airfoil with only a torsional degree-of-freedom will not flutter. It is the coupling of the two degrees-of-freedom that allows flutter to possibly exist. Past experience has shown that near coincidence of the bending and torsion frequencies is one of the main factors in the cause of flutter, and that the character of the torsional mode plays a more fundamental part in its occurrence.

The only source of external energy causing the flutter is that of the airstream. Hence the airstream must be the key component in the instability of the airfoil. In fact, a complete analysis of the interaction of the structural forces coupled with the aerodynamic forces is required to determine when the system instability will occur. The aerodynamic forces increase with velocity. However, the stiffness in the bending and torsional directions is independent of the air speed. Any damping that is introduced to the system comes from the aerodynamic forces: the structural damping is often so negligible that it is omitted in most aeroelastic analysis. Hence, there exists a critical speed at which the aerodynamic

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damping is insufficient to give system stability, and flutter occurs.

1.3 Unsteady Aerodynamics

The interaction of an airstream with an airfoil produces a behaviour that is very complicated. The calculation of the resulting aerodynamic forces presents, possibly, the major difficulty in the analysis of flutter. The analysis of flutter has seen a gradual progression in the complexity of the aerodynamic forces, this has considered the aerodynamics initially as steady, then quasi-steady, and finally the true unsteady motion of the aerodynamics has been considered.

When an airflow reacts with an airfoil, disturbances are introduced that result in the flow following the contour of the airfoil. According to thin airfoil theory (Fung 1955), the airfoil can be represented by a continuous distribution of vorticity. The strength and configuration of this vortex sheet produces the aerodynamic properties of the flow for the desired airfoil. However, this is not quite correct, not only does there exist the bound vortices on the airfoil, but there also exist free vortices in the wake of the airfoil. These free vortices are being produced continuously by the oscillating airfoil, shed at the trailing edge of the wing and carried downstream by the flow. These free vortices also produce vertical components of velocity on the wing, thus altering the aerodynamic properties of the wing. Steady flow analysis completely neglects the frequency of oscillation of the airfoil, and hence the free vortices produced by this oscillation.

Quasi-steady flow analysis realizes that vortices are being produced in the wake, but assumes they have no effect. This is true if these free vortices are swept quickly downstream. A major simplification of applying a quasi-steady analysis is that the aerodynamics can be considered independent of the airfoil's frequency of oscillation, which is generally not the case. However, if this assumption is made, it eliminates a time consuming iterative process. The strength of this theory is that it gives qualitative information about the system, without an extreme computational burden; but caution must be taken, since the method has made many simplifying assumptions.

Unsteady flow analysis accounts for the free vortices in the wake. The aerodynamics are now a function of the frequency of oscillation, k = wb/U (or $k_c = wc/U$), where w is frequency, c = 2b is the chord and U is the free stream velocity. Thus, for a given velocity, the frequency of oscillation is required, before the aerodynamic forces can be determined. This results in an iterative procedure. If the frequency of oscillation is small, steady or quasi-steady flow analysis may suffice, but for the most part this is not the case and an unsteady flow analysis must be used. Due to the interaction of the flow with the airfoil geometry, lag effects result in the flow; the circulation around the airfoil is not developed instantaneously due to a change in incidence, but there is a phase lag between the motion and the resulting aerodynamic forces. These lag effects can lead to an amplification of the small oscillatory motions.

The main difficultly in solving the aeroelastic equations of motion rests with the calculation of the unsteady aerodynamics. The unsteady aerodynamics are sometimes written as Q_h and Q_{α} , where $Q_h = -L$ and $Q_{\alpha} = M_{\alpha}$ are the total unsteady, nonconservative forces and moments acting about the elastic axis, respectively.

The quantities Q_h and Q_o (Lee 1984) can be expressed in this form

$$Q_h = -qc(C_{l_h}\frac{\xi}{2} + C_{l_\alpha}\alpha) , \qquad (1)$$

$$Q_{\alpha} = qc^2 (C_{m_h} \frac{\xi}{2} + C_{m_{\alpha}} \alpha) , \qquad (2)$$

where, C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} are the nondimensional derivatives of the lift and moment coefficients, with respect to plunging, h, and pitching, α , respectively. The dynamic pressure, q, is a function of velocity, and given by

$$q = \frac{1}{2}\rho V^2 , \qquad (3)$$

where ρ is the air density.

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The determination of the unsteady aerodynamics is not an easy task. Theodorsen has determined analytically the unsteady aerodynamics for incompressible flow on a twodimensional airfoil (Theodorsen 1935; Garrick and Theodorsen 1941). The derivation of these equations is based on potential flow theory and the Kutta condition at the trailing edge, and assumes inviscid, irrotational flow. All derivations assume small amplitude oscillations, and thus, this allows for the linearization of the equations. The unsteady aerodynamics, Q_h and Q_{α} , are written as a function of the Theodorsen Function, C(k), and are expressed as:

$$Q_{h} = -L = -\pi\rho b^{2} \left(\ddot{h} + U\dot{\alpha} - ba\ddot{\alpha}\right) - 2\pi\rho U b C(k) \left(\dot{h} + U\alpha + b(0.5 - a)\dot{\alpha}\right) , \qquad (4)$$

and

$$Q_{\alpha} = M_{\alpha} = \pi \rho b^{2} \left(b a \ddot{h} - U b (0.5 - a) \dot{\alpha} - b^{2} (.125 + a^{2}) \ddot{\alpha} \right) + 2\pi \rho U b^{2} (a + 0.5) C(k) \left(\dot{h} + U \alpha + b (0.5 - a) \dot{\alpha} \right) .$$
(5)

The Theodorsen Function C(k), composed of a real part F and an imaginary part G, is a function of reduced frequency, k,

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$$C(k) = F + \iota G , (6)$$

$$k = \frac{wb}{U} , \qquad (7)$$

where w is frequency, b is a characteristic length, the semi-chord, and U is the velocity. Note that the equations are both a function of velocity and frequency. This will result in an iterative type solution even for a fixed velocity.

For the transonic regime, readily solvable equations do not exist. This is due to the nonlinearity and irregularity of the associated unsteady transonic aerodynamics. The analysis of transonic flow started with linear potential theory in the 1950's and has progressed to the use of the nonlinear small disturbance equations, nonlinear Euler and the Navier-Stokes equations. Several methods, such as the relaxation methods (STRANS2 and UTRANS2) developed by Traci, Albano and Farr in 1974, and the indirect method (LTRAN2) developed by Ballhaus and Goorjian in 1977, can be used to determine the unsteady aerodynamic coefficients C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} as a function of reduced frequency, k, for transonic flow. LTRAN2 was based on the transonic small disturbance equations, and utilized a time integration (indicial) finite difference method. However, it gave accurate results only for low values of reduced frequency. Rizzetta and Chin furthered the work of Ballhaus and Goorjian in 1979, by integrating the LTRAN2 program and the structural equations of motion. Since then, the work by Houwink and van der Vooren, and later by Couston and Angelini, have taken the original code and produced much improvement in the accuracy and range of results for increased reduced frequencies.

It is now routine to calculate the aerodynamic derivatives, C_{l_h} , $C_{l_{\alpha}}$, C_{m_h} and $C_{m_{\alpha}}$

as a function of reduced frequency, k, to give the unsteady aerodynamics required for the aeroelastic equations of motion. The unsteady aerodynamics are calculated by performing computer simulations of pure bending and pure twisting in small disturbance flows. These equations cannot be linearized because the flow behaviour is so irregular in the transonic regime. Finding these unsteady aerodynamics for this regime is important since this regime tends to be the most critical for the flutter of wings. Lee (1984) gives the results

for M = 0.80, 0.85 and 0.875.

1.4 Flutter Solution

Prior to 1938, it was thought that the solution to the flutter problem could be found by flight testing alone. It was supposed that flutter could be determined by noting a reduction in the system damping with increase in velocity. Unfortunately, in February of that year, during a carefully planned flight test, an aircraft went down killing all the crew and scientists. Such testing had a number of shortcomings that affected reliability and safety (Weissenburger and Zimmerman 1964). In particular, it was shown that flutter can occur very quickly and without much warning. Thus, it was realized that to determine the onset of flutter, a theoretical analysis must be used at the initial design stage and for the modification of the existing aircraft, in conjunction with flight testing still being the ultimate test of when flutter will occur.

for the transonic unsteady aerodynamic derivatives as a function of reduced frequency, k,

New theoretical techniques consider an eigenvalue analysis of the flutter equations of motion. A two-degree-of-freedom rigid airfoil is represented as a flat plate, which is flexibly mounted and undergoing bending and twisting (torsion) motion (see Figure 1). Bending (plunge) is denoted by h, positive downward, and the twisting (pitch) about the elastic axis is denoted by α , positive in the clockwise direction. The elastic axis is defined as the point in which a load force would produce only pure bending (no twist). The aeroelastic equations of motion, which are a function of inertia forces, elastic forces and aerodynamic forces (Fung 1955; Ashley and Bisplinghoff 1962), can be derived by summing forces and moments, about the elastic axis:

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$$m\ddot{h} + m(x_{\alpha}b)\ddot{\alpha} = -K_{h}h - L , \qquad (8)$$

$$(I + m(x_{\alpha}b)^{2})\ddot{\alpha} + m(x_{\alpha}b)\ddot{h} = -K_{\alpha}\alpha + M_{\alpha} .$$
⁽⁹⁾

where *m* is the mass per unit span, *I* is the wing mass moment of inertia about the centre of gravity, c = 2b is the chord length, $x_{\alpha}b$ is the distance from the centre of gravity to the elastic axis, *ab* is the distance measured from the midchord to elastic axis, both positive towards the trailing edge, K_h is the spring constant in bending, K_{α} is the spring constant in twisting, and *L* and M_{α} are the lift force and moment acting about the elastic axis, respectively.

The calculated roots of the system aeroelastic equations give an indication of stability. This form of analysis gives an alternative indication of the onset of flutter. There exists two well established techniques for determining flutter; in particular for dealing with the iterative nature of the solution, namely the p-k (British) and U-g (American) methods, see for example Lee (1984).

In the U - g method, harmonic oscillatory motion is assumed, such that the generalized coordinates, ξ and α , can be written in the following form:

 $\xi = \xi_o e^{-iwt}$

$$\alpha = \alpha_o e^{-iwt} , \qquad (10)$$

where ξ_o and α_o are complex constants, the absolute values of which represent the amplitudes, and their arguments give phase angles. A theoretical structural damping component, g, is then introduced into the equations of motion, in order to produce steady harmonic system motion. The rational behind the U - g method is that the system will require just a sufficient amount of this fictitious hysteretic structural damping, g, to maintain steady harmonic motion. If the system requires the input of negative damping (-g) to make it harmonic, it is stable, likewise if it requires positive damping (+g) to make it harmonic, it is unstable. An expression for the damping ratio (Lee 1984) can be written as

$$\gamma = -\frac{g}{2} . \tag{11}$$

Hence, a positive value of damping ratio indicates a stable system.

In the p - k method, the generalized coordinates, ξ and α , take on a more general form

$$\xi = \xi_c e^{-pt}$$

$$\alpha = \alpha_o e^{-pt}$$
(12)

where ξ_o and α_o are complex constants, and $p = \beta + iw$, where β gives an indication of damping and w gives frequency. The p - k method solves for the critical flutter velocity by utilizing an iteration procedure of frequency for each eigenvalue. The damping ratio (Lee 1984), a function of β and w, can be given as

$$\zeta = -\frac{\beta}{\sqrt{\beta^2 + w^2}} \ . \tag{13}$$

Flutter occurs when there is insufficient aerodynamic damping to keep the system stable. Hence, a negative value for the system damping ratio would indicate an instability. Both methods agree at the critical instability point, but it has been found that the American approach can largely over estimate the magnitude of the relative damping ratio for values of subcritical velocity (Jackson and Lawrence 1968). These methods have been used to great success in flutter analysis, but they tend to be time consuming iterative processes. In recent years, the focus of flutter analysis has turned away from the time domain analysis to the frequency domain or Laplace domain.

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The Laplace domain employs rational function approximations of the unsteady aerodynamic forces. When the unsteady aerodynamics are approximated this eliminates much of the iterative work that characterizes the previous methods. The aerodynamic forces are usually found in tabular form for simple harmonic motion at discrete values of reduced frequency. The tabular data for the unsteady aerodynamics has been, for sometime, readily available for the incompressible regime, in the form of the Theodorsen Function, C(k). It is not until recently, that advancements in transonic numerical methods for computing unsteady aerodynamic forces in the transonic regime, have made it possible to consider this regime for further aeroelastic analysis. Transonic codes can now routinely calculate the required unsteady aerodynamic coefficients C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} at various Mach numbers.

The idea of using rational functions to approximate unsteady aerodynamics is not a new concept. Jones (1940) is credited with first using rational Laplace transform functions to approximate unsteady aerodynamics. He considered an approximation of the Theodorsen Function, C(k), of this form,

$$C(k) = 1 - \frac{0.165ik}{ik + 0.0455} - \frac{0.335ik}{ik + 0.32} .$$
⁽¹⁴⁾

The importance of such approximations were not fully realized until its versatility in the

analysis of active control systems was discovered. However, for the proper analysis of the aeroelastic equations of motion, the governing rational functions have to be somewhat more robust, than that first introduced by Jones. A rational function is required that gives good approximation to the unsteady aerodynamics as a function of reduced frequency, with a compatible amount of complexity. The approximation of the incompressible unsteady aerodynamics, namely the Theodorsen Function, C(k), as a function of reduced frequency, k, may not pose much of a difficulty; however, problems may be encountered in the transonic regime, where the unsteady transonic aerodynamic coefficients C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} vary in a very irregular manner with reduced frequency. Thus, in order for the Laplace method to be viable, the approximations of these tabular unsteady aerodynamics can not be considered lightly.

A function Q(ik) is required that approximates Q(ik), as a function of reduced frequency, k, and gives good results over the entire specified range. The tabular data Q(ik) of the unstandy aerodynamic forces may be written in terms of real and imaginary parts as a function of reduced frequency,

$$Q(ik) = F + iG av{15}$$

where F and G are the real and imaginary parts, respectively.

Many rational functions can be used to approximate the unsteady aerodynamics (Poirel 1988). The Nasa-Langley rational function has been shown to be fairly robust in its ability to approximate given data (Adams and Tiffany 1988), and is given by

$$\bar{Q}(p) = A_0 + A_1 p + \dots + A_P p^{\bar{P}} + \sum_{l=1}^R A_{l+P} \frac{p}{p+b_l} , \qquad (16)$$

where p = ik. The coefficients A_0 , A_1 , ... and b_1 , b_2 , ... are chosen to ensure minimum approximation error to the given tabular data. This is the form of the approximation

employed by Tiffany and Adams (1988), and the same method that is used in this thesis. The associated error distribution seems to be less, and it has a built in robustness that lets it consider a larger initial frequency range. The physical interpretation of the equation is that the first three terms stem from quasi-steady influences, and the other terms give the unsteady effects. The lag terms b_j give an approximation to the time delays (lag in the development of the circulation about the airfoil) inherent in unsteady aerodynamics.

In the past, these lag terms were not optimized, but arbitrarily selected from the range of reduced frequencies for which tabular data of the unsteady aerodynamic forces were available (Abel 1979). Present methods, that optimize these nonlinear lag terms, give better results with a decreased order in the approximating equation, and thus a solution of less complexity and increased efficiency. Optimization of the rational function coefficient to ensure minimum approximation error was performed by utilizing a least square method to determine the *linear coefficients* and a sequential simplex method to determine the non-linear lag terms. The method employed to solve for the optimal nonlinear lag terms is a sequential simplex method (Mead and Nelder 1965; Nelson and Olsson 1975) developed by Nelder and Mead in 1965. This method is simple, robust and requires no derivatives, and hence lends itself quite well to finding the minimum of a non-linear objective function of more than one independent variable. The validity of the Laplace solution depends greatly on how well the unsteady aerodynamics can be approximated.

The advantage of approximating the unsteady aerodynamics is that it allows the acroelastic equations of motion to be written in the Laplace domain. The equations. in this form, can be solved quite easily by readily available computer subroutines, that utilize matrix algebra and eigenvalue analysis. Abel (1979) has demonstrated how the acroelastic equations of motion in the Laplace domain can be solved. This requires the expansion of the equations to obtain a m^{th} order polynomial of the following form

$$[F_m]s^m + [F_{m-1}]s^{m-1} + [F_{m-2}]s^{m-2} + \dots + [F_1]s + [F_0] = 0 , \qquad (17)$$

where the coefficients $[F_m]$, $[F_{m-1}]$, $[F_{m-2}] \cdots [F_0]$ are functions of dynamic pressure (velocity).

This can be easily placed in a typical eigenvalue problem of the given form

$$s\{x\} = [A]\{x\},$$
(18)

to give a series of $m \times n$ first order equations, where

$$\{x\} = \{ s^{m-1} , s^{m-2} \cdots s^1 , s^0 \}.$$
⁽¹⁹⁾

These equations can be solved quite easily by linear matrix techniques.

1.5 Flutter-Suppression Systems

Flutter-suppression systems have received increased attention in the last two decades (Kass and Thompson 1971; Guruswamy, Olsen, Striz and Yang 1980; Karpel 1981). They have been found to be an effective way of increasing the flutter speed. There are two types of flutter-suppression systems: passive and active flutter control. Passive control includes increasing the structural stiffness and/or mass balancing. Proper mass balancing has a direct effect on the system inertia coupling, which is often a predominant factor in the occurrence of flutter. Active control uses a control surface which is deflected in response to the wing motion, resulting in a change in the aerodynamic forces on the wing. The first application of such a system in flight was on the B-52 in the mid 1970's. Much work has been done in wind tunnels and flight tests to show the feasibility of implementing active control (Nissim and Abel, 1978; Abel 1979; Messina and O'Connell, 1979). Increases in critical flutter speed of 20-25 % have been obtained through these methods.

The problem associated with applying an active control device to an airfoil hinges on developing a set of aeroelastic equations, where the form of the unsteady aerodynamics and the control law are compatible. This can be achieved by approximating the unsteady aerodynamics by a rational function, and thus allows for both the unsteady aerodynamics and the control law to be written in the Laplace domain.

The aeroelastic equations for a three-degree-of-freedom rigid airfoil including the incorporation of active control, a flap, (see Figure 10) are given by numerous references (Fung 1955; Ashley and Bisplinghoff 1962). The flap is deflected by an angle β , in response to the wing motion, a function of the plunge, h/b, and pitch, α , with the aim of causing an increase in the flutter speed. Control laws of the following form have been considered (Nissim and Abel 1978),

$$\{ \beta \} = [T] \left\{ \begin{array}{c} h/b \\ \alpha \end{array} \right\} , \qquad (20)$$

where [T] is a transfer function matrix of size 1×2 .

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The transfer function used was a function of h/b and α and their first derivatives,

$$[T] = [t_{11}, t_{12}] + \frac{a_T}{w_R} s[t_{11}^*, t_{12}^*] .$$
(21)

This is referred to as a *damping type transfer function* by Nissim and Abel (1978), where a_T is a control of the amount of damping introduced by the control surface (flap), and w_R is a reference frequency, normally taken as the no-control flutter frequency. This transfer function can be considered as a type of proportional-derivative (PD) control.

In the late 1970's, considerable work was done on flutter control systems to determine realistic transfer function parameters (Nissim and Abel 1978; Nissim 1977; Nissim 1971). Nissim and Abel (1978) used a theory based on the aerodynamic energy concept to determine realistic transfer function parameter values. The aerodynamic energy concept considers the work done by the aerodynamic forces on the wing per cycle of oscillation, referred to as P. Hence, the transfer function [T] was determined to give a stable system, or a negative value of work P. Nissim and Abel determined the optimal parameters for the transfer function for an airfoil of specific geometric configuration. A simple transformation was also derived that allowed for changes in airfoil geometry.

With the introduction of the control law into the aeroelastic equations of motion, the equations can be expanded, as before, to obtain a m^{th} order polynomial of the following form

$$[F_m] s^m + [F_{m-1}] s^{m-1} + [F_{m-2}] s^{m-2} + \dots + [F_1] s + [F_0] = 0 , \qquad (22)$$

where the coefficients $[F_m]$, $[F_{m-1}]$, $[F_{m-2}] \cdots [F_0]$ are functions of dynamic pressure (velocity). The solution follows the same methodology as shown earlier.

1.6 Thesis Overview

This study considers the analysis of the aeroelastic equations of motion of a rigid airfoil, flexibly mounted, undergoing bending and twisting motion, as it pertains to flutter. Initially, the formulation of the equations of motion for a two-degree-of-freedom system are considered in chapter 2. Chapter 3 discusses the unsteady aerodynamics and their associated difficulties for both the incompressible and transonic regimes. An overview of eigenvalue analysis methods, techniques that are commonly used for all forms of analyt-

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ical flutter calculation, is described in detail in chapter 4. In chapter 5 the traditional solutions to the flutter problem, namely the p - k and U - g, are derived for the general case of incompressible flow.

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In chapter 6 the Laplace method is introduced. This method requires the approximation of the unsteady forces and moments by a rational function. The rational function employed for this thesis is of the Nasa-Langley form. It requires the use of a two stage function optimization (Adams and Tiffany 1988). A least square method is used to determine the linear coefficients and a sequential simplex method (Mead and Nelder 1965) to determine the nonlinear lag terms. A general Laplace solution is derived to consider not only incompressible flow, but also that of transonic flow.

Chapter 7 is devoted mainly to the presentation and discussion of the results for the approximation of the unsteady aerodynamics, for both the incompressible and transonic regimes. The following chapter, chapter 8, is devoted to a comparison of the traditional methods, namely the U - g and p - k, with the Laplace method. Comparisons are done in the transonic regime by using previous U - g and p - k results given by Lee (1984).

In Chapter 9, the incompressible equations of motion for a three-degree-of-freedom airfoil are formulated to consider control. Active control is considered with the aim of obtaining an increase in the critical flutter speed. The values suggested by Abel and Nissim, who employed the Aerodynamic Energy Concept (Abel and Nissim 1978), are used as a starting point to find the optimal transfer function parameters for the control law employed. The simplex method is then used to determine if these are indeed the optimal values. Chapter 10 is devoted to a discussion of the consequences of implementing active control for the incompressible flow regime.
The final chapter, chapter 11, summarizes the conclusions.

The Appendices contain additional notes, tables and figures, as well as, complete computer program listings (U - g, p - k and Laplace methods, rational function optimization, incompressible control using Laplace method).

2 Equations of Motion 2DOF

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Consider a two-degree-of-freedom rigid airfoil, represented as a flat plate, flexibly mounted undergoing bending and twisting (torsion) motion (see Figure 1). Bending (plunge) is denoted by h, positive downward, and the twisting (pitch) about the elastic axis is denoted by α , positive in the clockwise direction. The elastic axis is defined as the point in which a load force would produce only pure bending (no twist). The aeroelastic equations of motion, which are a function of inertia forces, elastic forces and aerodynamic forces (Fung 1955; Ashley and Bisplinghoff 1962), can be derived by summing forces and moments, about the elastic axis:

$$m\ddot{h} + m(x_{o}b)\ddot{\alpha} = -K_{h}h - L , \qquad (23)$$

$$(I + m(x_{\alpha}b)^{2})\ddot{\alpha} + m(x_{\alpha}b)\ddot{h} = -K_{\alpha}\alpha + M_{\alpha} .$$
⁽²⁴⁾

where *m* is the mass per unit span, *I* is the wing mass moment of inertia about the centre of gravity, c = 2b is the chord length, $x_{\alpha}b$ is the distance from the centre of gravity to the elastic axis, *ab* is the distance measured from the midchord to elastic axis, both positive towards the trailing edge, K_h is the spring constant in bending, K_{α} is the spring constant in twisting, and *L* and M_{α} are the lift force and moment acting about the elastic axis, respectively.

Past experience has shown that structural damping can often be assumed to be negligible for wing flutter analysis, and thus, it is not included.

Introducing a non-dimensional generalized displacement coordinate

$$\xi = \frac{h}{b} \tag{25}$$

$$K_h = w_h^2 m , \qquad (26)$$

$$K_{\alpha} = w_{\alpha}^2 I_{\alpha} , \qquad (27)$$

where

$$I_{\alpha} = J + m(x_{\alpha}b)^2 , \qquad (28)$$

gives

$$mb\ddot{\xi} + m(x_{\alpha}b)\ddot{\alpha} = -w_h^2 mb\xi - L , \qquad (29)$$

$$I_{\alpha}\ddot{\alpha} + m(x_{\alpha}b^2)\ddot{\xi} = -w_{\alpha}^2 I_{\alpha}\alpha + M_{\alpha} .$$
(30)

The above two equations may be simplified by dividing through by mb and mb^2 , respectively to give

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + w_h^2\xi = \frac{-L}{mb} , \qquad (31)$$

$$\frac{I_{\alpha}}{mb^2}\ddot{\alpha} + x_{\alpha}\ddot{\xi} + w_{\alpha}^2 \frac{I_{\alpha}}{mb^2}\alpha = \frac{M_{\alpha}}{mb^2} .$$
(32)

Finally, using $I_{\alpha} = r_{\alpha}^2 m b^2$, where r_{α} is the radius of gyration about the elastic axis, the complete aeroelastic equations become

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + w_h^2\xi = \frac{-L}{mb} , \qquad (33)$$

$$r_{\alpha}^{2}\ddot{\alpha} + x_{\alpha}\ddot{\xi} + w_{\alpha}^{2}r_{\alpha}^{2}\alpha = \frac{M_{\alpha}}{mb^{2}}.$$
(34)

3 Unsteady Aerodynamics

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The calculation of the unsteady aerodynamics forces presents, possibly, the major difficulty in the analysis of flutter. The unsteady aerodynamics are sometimes written as Q_h and Q_{α} , where $Q_h = -L$ and $Q_{\alpha} = M_{\alpha}$ are the total unsteady, nonconservative forces and moments acting about the elastic axis, respectively.

The quantities Q_h and Q_{α} can be expressed in this form

$$Q_h = -qc(C_{l_h}\frac{\xi}{2} + C_{l_a}\alpha) , \qquad (35)$$

$$Q_{\alpha} = qc^2 (C_{m_h} \frac{\xi}{2} + C_{m_\alpha} \alpha) , \qquad (36)$$

where, C_{l_h} , C_{l_α} , C_{m_h} and C_{m_α} are the derivatives of the lift and moment coefficients, with respect to h, plunging, and α , pitching, respectively. The dynamic pressure, q, is a function of velocity, and given by

$$q = \frac{1}{2}\rho V^2 . \tag{37}$$

The solution to the flutter equations of motion is complicated by the nature of the unsteady aerodynamics. They tend to vary in a very nonlinear manner with reduced frequency, which results in an iterative type solution.

3.1 Incompressible Two-Dimensional Aerodynamics

In the incompressible regime, Q_h and Q_α take on exact forms, expressed as a function of the Theodorsen Function C(k) (Theodorsen 1935; Garrick and Theodorsen 1941),

$$Q_{h} = -L = -\pi\rho b^{2} \left(\ddot{h} + U\dot{\alpha} - ba\ddot{\alpha} \right) -$$

$$2\pi\rho UbC'(k)\left(\dot{h} + U\alpha + b(0.5 - a)\dot{\alpha}\right) , \qquad (38)$$

and

$$Q_{\alpha} = M_{\alpha} = \pi \rho b^{2} \left(b a \ddot{h} - U b (0.5 - a) \dot{\alpha} - b^{2} (.125 + a^{2}) \ddot{\alpha} \right) + 2\pi \rho U b^{2} (a + 0.5) C(k) \left(\dot{h} + U \alpha + b (0.5 - a) \dot{\alpha} \right) .$$
(39)

The Theodorsen Function, C(k), (see Table 1), composed of a real part F and an imaginary part G, is a function of reduced frequency, k,

$$C(k) = F + iG \tag{40}$$

$$k = \frac{wb}{U} \tag{11}$$

where w is frequency, b is a characteristic length, the semi-chord, and U is the velocity. Note that the equations are a function of both velocity and frequency. This will result in an iterative type solution even for a fixed velocity.

The derivation of these equations is based on potential flow theory and the Kutta condition at the trailing edge, and assumes incompressible, inviscid and irrotational flow. Furthermore, all derivations assume small amplitude oscillations for h and α , so the flow remains potential and unseparated. This allows for the linearization of the equations.

Theodorsen Function C(k)

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k	F	G
10.00	0.5006	-0.0124
6.000	0.5017	-0.0206
4.000	0.5037	-0.0305
3.000	0.5063	-0.0400
2.000	0.5129	-0.0577
1.500	0.5210	-0.0736
1.000	0.5394	-0.10 03
0.800	0.5541	-0.1165
0.500	0.5979	-0.15 07
0.400	0.6250	-0.1650
0.300	0.665 0	-0.1793
0.200	0.7276	-0.1886
0.100	0.8320	-0.1723
0.050	0.9090	-0.1305
0.025	0.9545	-0.0872
0.010	0.9824	-0.0482
0.000	1.0000	-0.0000

Pable 1	Theodorsen	Function	C(h) -	$F(h) \perp$	iC(h)	where k	_	20h / 1	T T
rable r	Theodorsen	runction	$C(\kappa) =$	$\mathbf{r}(\kappa) +$	$\iota G(\kappa),$	where κ	=	<i>wo</i> / ($\mathcal{O}_{\mathbb{C}}$

3.2 Transonic Two-Dimensional Aerodynamics

As yet, no such readily solvable equations for Q_h and Q_α exist for the transonic regime. Due to the associated nonlinearities and irregularity in this regime, it has been very difficult to model the unsteady aerodynamics. However, the computation of transonic unsteady aerodynamics has seen much development in the last decade.

Before the development of transonic numerical codes, aeroelastic analysis in this regime was considered difficult, if not impossible. Fortunately, much progress has been made in this field. The numerical methods available consider the uncoupled motion of an airfoil pitching and plunging in small disturbance flow. It also assumes the flow is inviscid and irrotational. The equations cannot be linearized, because the behaviour of flow in the transonic regime is so nonlinear. The codes now available allow for the routine calculation of the required unsteady aerodynamics, C_{l_h} , C_{l_α} , C_{m_h} and C_{m_α} (Lee 1984). The aerodynamic derivatives for a NACA64A006 airfoil at M = 0.85 are shown graphically in Figures 2a to 2d. Each aerodynamic derivative is composed of a real and an imaginary part, where the real part represents the total forces and moments in phase with the airfoil motion, and the imaginary part represents those ninety degrees out of phase with the motion. Refinements in existing codes have allowed for more time efficient calculations, as well as an increase in the range of reduced frequency, k, and for different values of M number.

Note:

The techniques presently available to generate the unsteady aerodynamics, give

the forces and moments only for purely oscillatory motion at discrete values of reduced frequency (tabular data). In order to obtain solutions for decaying and growing motions (not simply harmonic, iw), the concept of analytic continuation is used. These results are extended off the imaginary axis, by iteration (or by using analytic approximating functions) of the given tabular unsteady aerodynamics. This procedure is not completely valid, but because flutter occurs along the imaginary axis of the complex plane, the approximation is sufficient.

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4 Eigenvalue Analysis

The numerical techniques that are used in this study to solve the aeroclastic equations of motion rely on the equations being placed in an eigenvalue format. The procedure that each method utilizes to obtain this form is slightly different, but the end result is the same.

If a standard 2DOF system is considered in eigenvalue form, the solution will generate two complex conjugate pairs (four eigenvalues) of the following form

$$\lambda_1 = a \pm ib ,$$

$$\lambda_2 = c \pm id . \tag{42}$$

The roots determine the stability of the system at a given value of dynamic pressure (velocity). The real part of the roots, a and c, give an indication of the damping, β , and the imaginary part, b and d, give frequency, w. Hence, the roots having negative imaginary parts are not considered (a system can not have a negative frequency). The damping ratio, a function of β and w, is given as

$$\zeta = -\frac{\beta}{\sqrt{\beta^2 + w^2}} \,. \tag{43}$$

Flutter occurs when there is insufficient aerodynamic damping to keep the system stable. Hence, a positive real part, a or c, indicates an instability (this gives a negative value for the system damping ratio). It is difficult to determine which root will produce the instability, therefore a complete tracking of both roots is required.

Flutter analysis is commonly considered by plotting both the system frequency and the damping ratio versus dynamic pressure or velocity. Commonly a non-dimensional value of velocity is used, given by

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$$\bar{U} = \frac{U}{bw_{\alpha}} . \tag{44}$$

For example, see Figure 3, (Abel 1979) where the dynamic instability occurs at a dynamic pressure of about 5.0. Another predominant characteristic of flutter is the coalescence of the two frequencies. It is not a complete coalescence, but the two frequencies become very close to one another at the onset of flutter. This can be seen from the graph of frequency versus dynamic pressure (see Figure 3).

5.1 U-g Method

In the U - g method, harmonic oscillatory motion is assumed such that the generalized coordinates, ξ and α , can be written in the following form:

$$\xi = \xi_o e^{iwt} ,$$

$$\alpha = \alpha_o e^{iwt} ,$$
(15)

where ξ_o and α_o are complex constants, the absolute values of which represent the amplitudes, and their arguments give phase angles.

A structural damping coefficient g is introduced into (33) and (34) by multiplying the third term (stiffness) of the two equations by the factor (1 + ig), a hysteretic damping. A sufficient amount of this structural damping, g, will be required to produce steady harmonic motion. For the case of incompressible aerodynamics, the appropriate equations for Q_h and Q_{α} , (38) and (39), respectively, are then substituted in to the equations, to give:

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + (1+ig)w_{h}^{2}\xi = \frac{Q_{h}}{mb}$$

$$= -\frac{\pi\rho b}{m} \left(b\ddot{\xi} + U\dot{\alpha} - ba\ddot{\alpha}\right) - \frac{2\pi\rho U}{m}C(k)\left(b\dot{\xi} + U\alpha + b(0.5-a)\dot{\alpha}\right), \quad (46)$$

$$x_{\alpha}\ddot{\xi} + r_{\alpha}^{2}\ddot{\alpha} + (1+ig)r_{\alpha}^{2}w_{\alpha}^{2}\alpha = \frac{Q_{\alpha}}{mb^{2}}$$
$$= \frac{\pi\rho}{m} \left(b^{2}a\ddot{\xi} - Ub(0.5-a)\dot{\alpha} - b^{2}(0.125+a^{2})\ddot{\alpha}\right) + \frac{2\pi\rho U}{m}(a+0.5)C(k)\left(b\dot{\xi} + U\alpha + b(0.5-a)\dot{\alpha}\right)$$
(47)

Multiplying through by $\mu = m/\pi\rho b^2$ (a non-dimensional airfoil mass) and $c^2/4U^2$ and making the following substitutions

$$\xi = \xi_o e^{iwt}$$
 and $\alpha = \alpha_v e^{iwt}$,

 $\lambda = iw ,$ $k_c = \frac{wc}{U} = 2k ,$ $\lambda_g = \mu (1 + ig) \frac{w_o^2 b^2}{U^2} ,$ (48)

the preceding equations can be expressed as

$$\left(\frac{1}{4}\mu k_c^2 - \lambda_g \left(\frac{w_h}{w_\alpha}\right)^2 + \frac{1}{4}k_c^2 - ik_c C(k)\right) \xi_o +$$

and

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$$\left(\frac{1}{4}x_{\alpha}\mu k_{c}^{2} - i\frac{1}{2}k_{c} - \frac{a}{4}k_{c}^{2} - 2C(k) - i(\frac{1}{2} - a)k_{c}C(k)\right) \alpha_{o} = 0, \qquad (49)$$

$$\left(\frac{1}{4}x_{\alpha}\mu k_{c}^{2} - \frac{1}{4}ak_{c}^{2} + i(a + \frac{1}{2})C(k)k_{c}\right) \xi_{o} + \left(\frac{1}{4}r_{\omega}^{2}\mu k_{c}^{2} - \lambda_{y}r_{\alpha}^{2} - \frac{i}{2}(\frac{1}{2} - a)k_{c} + \frac{1}{4}(\frac{1}{8} + a^{2})k_{c}^{2} + 2(a + \frac{1}{2})C(k) + i(a + \frac{1}{2})(\frac{1}{2} - a)C(k)k_{c}\right) \alpha_{o} = 0$$

$$(50)$$

The above equations can be placed in matrix form as

$$[A] \{x\} = \lambda_g [B] \{x\}, \qquad (51)$$

where [A], [B] are 2×2 matrices and

 $\{x\} = \{\xi_o, \alpha_o\}^T.$

This typical eigenvalue problem can be solved to obtain λ_g for different values of reduced frequency, k, namely

$$\lambda_{g_1} = \bar{a} + ib ,$$

$$\lambda_{g_2} = \bar{c} + i\bar{d} , \qquad (52)$$

Note that

$$\lambda_{g_{1,2}} = \mu (1 + ig) \frac{w_{\alpha}^2 b^2}{U^2} = \mu \frac{w_{\alpha}^2 b^2}{U^2} + i\mu g \frac{w_{\alpha}^2 b^2}{U^2} .$$
(53)

The only unknowns in these equations are g, the structural damping coefficient, and U, the free stream velocity. Considering λ_{g_1} , U_1 and g_1 can be solved by the following equations

$$\lambda_{g_1}=\bar{a}+i\bar{b}\;.$$

Thus

$$U_1 = \sqrt{\frac{\mu w_\alpha^2 b^2}{\bar{a}}} = w_\alpha b \sqrt{\frac{\mu}{\bar{a}}} , \qquad (54)$$

$$g_1 = \frac{U_1^2 \bar{b}}{w_{\alpha}^2 b^2 \mu} = \frac{\bar{b}}{\bar{a}} , \qquad (55)$$

and the frequency,

$$w_1 = \frac{U_1 k}{b} . \tag{56}$$

An expression for the damping ratio (Lee 1984) can be written as

$$\gamma_1 = -\frac{g_1}{2} \ . \tag{57}$$

A similar analysis for λ_{g_2} , will give U_2 , g_2 , w_2 and γ_2 . Thus, each value of reduced frequency, k, will give two velocities and their respective associated frequency and damping ratio.

The rational behind the l'-g method is that the system will require just a sufficient amount of this fictitious hysteretic structural damping, g, to maintain steady harmonic motion. If the system requires the input of negative damping (-g) to make it harmonic, it is stable; likewise if it requires positive damping (+g) to make it harmonic, it is unstable. If the system response is harmonic, hence g = 0, this indicates a stability boundary.

5.2 p-k Method

In the p-k method the generalized coordinates, ξ and α , take on a more general form

$$\xi = \xi_o e^{-pt}$$

$$\alpha = \alpha_o e^{-pt}$$
(58)

where ξ_o and α_o are complex constants, and $p = \beta + iw$, where β gives an indication of damping and w gives the frequency.

Substituting the appropriate equations for Q_h and Q_α in (33) and (34), again, for the particular case of incompressible aerodynamics, gives

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + w_{h}^{2}\xi = \frac{Q_{h}}{mb}$$

$$= -\frac{\pi\rho b}{m} \left(b\ddot{\xi} + U\dot{\alpha} - ba\ddot{\alpha} \right) - \frac{2\pi\rho U}{m} C(k) \left(b\dot{\xi} + U\alpha + b(0.5 - a)\dot{\alpha} \right) , \quad (59)$$

$$x_{\alpha}\ddot{\xi} + r_{\alpha}^{2}\ddot{\alpha} + r_{\alpha}^{2}w_{\alpha}^{2}\alpha = \frac{Q_{\alpha}}{mb^{2}}$$

$$= \frac{\pi\rho}{m} \left(b^{2}a\ddot{\xi} - Ub(0.5 - a)\dot{\alpha} - b^{2}(0.125 + a^{2})\ddot{\alpha} \right) + \frac{2\pi\rho U}{m} (a + 0.5)C(k) \left(b\dot{\xi} + U\alpha + b(0.5 - a)\dot{\alpha} \right) . \quad (60)$$

Grouping common powers of ξ and α , and using the substitution $\mu = m/\pi\rho b^2$, the above equations can be rewritten as

$$\ddot{\xi}(1+\frac{1}{\mu}) + \dot{\xi}(\frac{2C(k)U}{b\mu}) + \xi(w_h^2) + \ddot{\alpha}(x_\alpha - \frac{a}{\mu}) + \dot{\alpha}\left(\frac{U}{b\mu} + 2C(k)(\frac{1}{2}-a)\frac{U}{b\mu}\right) + \alpha(\frac{2C(k)U^2}{b^2\mu}) = 0 , \qquad (61)$$

and

$$\ddot{\xi}(x_{\alpha} - \frac{a}{\mu}) + \dot{\xi}\left(-2(a + \frac{1}{2})C(k)\frac{U}{b\mu}\right) + \ddot{\alpha}\left(r_{\alpha}^{2} + \frac{\frac{1}{8} + a^{2}}{\mu}\right) + \dot{\alpha}\left(\frac{U}{b\mu}(\frac{1}{2} - a) - 2(a + \frac{1}{2})(\frac{1}{2} - a)C(k)\frac{U}{b\mu}\right) + \alpha\left(r_{\alpha}^{2}w_{\alpha}^{2} - 2(a + \frac{1}{2})C(k)\frac{U^{2}}{b^{2}\mu}\right) = 0.$$
(62)

These two equations can then be expressed in matrix form as

$$[M] \left\{ \begin{array}{c} \ddot{\xi} \\ \ddot{\alpha} \end{array} \right\} + [C] \left\{ \begin{array}{c} \dot{\xi} \\ \dot{\alpha} \end{array} \right\} + [K] \left\{ \begin{array}{c} \xi \\ \alpha \end{array} \right\} = 0 , \qquad (63)$$

where [M], [C] and [K] are the structural and aerodynamic mass matrix, damping matrix and stiffness matrix, respectively.

A solution to this system of equations is done by considering an eigenvalue analysis of the equations rewritten in the following form

$$\begin{bmatrix} M & C \\ 0 & I \end{bmatrix} \begin{cases} \ddot{z} \\ \dot{z} \end{cases} + \begin{bmatrix} 0 & K \\ -I & 0 \end{bmatrix} \begin{cases} \dot{z} \\ z \end{cases} = 0 , \qquad (64)$$

where

$$I = \left[\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array} \right]$$

and



Equation (61) can be written as

$$[D] \{\dot{p}\} + [E] \{p\} = 0, \qquad (65)$$

where

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$$p = \left\{ \dot{z} \ z \right\}^T$$

and [D] and [E] are 4×4 matrices, that are a function of the Theodorsen Function, C(k), and velocity, U. By letting $p = p_o e^{\lambda t}$, the resulting eigenvalue problem will give a solution of two complex conjugate pairs,

$$\lambda_{1,2} = \bar{a} \pm i b ,$$

$$\lambda_{3,4} = \bar{c} \pm i \bar{d} ,$$
(66)

where \bar{a} and \bar{c} are damping, and \bar{b} and \bar{d} represent frequency, w.

The solution to this problem is an iterative one. Initially, a specific velocity is chosen, U, and a reduced frequency, k, guessed. This then gives a frequency, w, by the equation

$$k=\frac{wb}{U}.$$

The choice of k, will give C(k).

When solving the eigenvalue problem, only the positive imaginary parts of the complex conjugate pairs are considered. The roots having negative imaginary parts are omitted from the analysis; a system cannot have a negative frequency. Considering only one of the roots, say $\lambda = \bar{a} + i\bar{b}$, the frequency \bar{b} is compared to w. If they are equal, a solution has been found, if not, it is required to calculate a new reduced frequency by

$$k = \frac{\bar{b}b}{U} \tag{67}$$

and the process is continued till convergence occurs. This same process is repeated for the other root. The idea is to create a table of velocity and its two associated solutions. The roots determine the stability, or instability, of the system at a given value of dynamic pressure (velocity). The real parts of the roots, \bar{a} and \bar{c} , give an indication of damping, β , and the imaginary parts, \bar{b} and \bar{d} , give frequency, w. The damping ratio (Lee 1984), a function of β and w, can be given as

$$\zeta = -\frac{\beta}{\sqrt{\beta^2 + w^2}} \,. \tag{68}$$

Flutter occurs when a negative value for the system damping is obtained.

6 Laplace Solution

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More efficient algorithms for determining flutter speeds can be found by solving the equations of motion in the Laplace domain. The main problem in this approach is that of finding a suitable rational function approximation to a given tabular data set of the unsteady aerodynamics as a function of reduced frequency. This eliminates the costly iterative process, that is characteristic of the U - g and p - k methods.

The tabular data for the unsteady aerodynamics has been, for sometime, readily available for incompressible flow, in the form of the Theodorsen Function, C(k). It is not until recently, that advancements in transonic numerical methods for computing unsteady aerodynamic forces in the transonic regime, have made it possible to consider this regime for further aeroelastic analysis. Transonic codes can now routinely calculate the required unsteady aerodynamic coefficients C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} at various Mach numbers.

6.1 Rational Function Approximation

The tabular data Q(ik) for the unsteady aerodynamic forces may be written in terms of its real and imaginary parts as a function of reduced frequency.

$$Q(ik) = F + iG , (69)$$

where F and G are the real and imaginary parts, respectively.

A function Q(ik) is required that approximates Q(ik), as a function of reduced frequency, k, and gives good results over the entire specified range. There are many possible methods available that use rational functions to approximate unsteady aerodynamic given in tabular form. One method employed by Poirel (1988) uses a Pade approximant in the following form

$$\bar{Q}(p) = \frac{N(p)}{D(p)} = \frac{\sum_{m=0}^{M} a_m(p)^m}{\sum_{l=1}^{L} b_l(p)^l} ,$$
(70)

where

p = ik.

The Pade approximant (Baker and Gammel 1970; Vepa 1976) is simply a fraction consisting of two polynomials of order M and L, where the coefficients a_m and b_l are chosen to give minimum error to the given unsteady aerodynamic tabular data. It leads to useful results, but accuracy in the range of interest depends upon a fairly complicated weighting scenario. This weighting at the point of interest tends to result in more error at the extremes of the desired frequency range. This method is accurate for a specified decreased overall reduced frequency range, but the approximation is not as robust when one is concerned with a larger overall range. Another method used to approximate given tabular data is that given by Nasa-Langley, and is often referred to as Roger's Approximation (Roger 1977)

$$\bar{Q}(p) = A_0 + A_1 p + \dots + A_P p^P + \sum_{l=1}^R A_{l+P} \frac{p}{p+b_l} .$$
(71)

The value of \bar{P} depends on the nature of the function, and R is the order-of-fit of the approximant. The coefficients A_0, A_1, \ldots and b_1, b_2, \ldots are chosen to ensure minimum approximation error to the given tabular data. With the given parameters \bar{P} and R made equivalent to 2 and 4, respectively, this gives

$$\bar{Q}(p) = A_0 + A_1 p + A_2 p^2 + A_3 \frac{p}{p+b_1} + A_4 \frac{p}{p+b_2} + A_5 \frac{p}{p+b_3} + A_6 \frac{p}{p+b_4} .$$
(72)

This is the form of the approximation employed by Tiffany and Adams (1988), and the method that is used in this thesis. The associated error distribution seems to be less, and

it has a built in robustness that lets it consider a large frequency range. The physical interpretation of the equation is that the first three terms stem from quasi-steady influences, and the other terms give the unsteady effects. The lag terms b_j give an approximation to the time delays (lag in the development of the circulation about the airfoil) inherent in unsteady aerodynamics.

In the past, these lag terms were not optimized, but arbitrarily selected from the range of reduced frequencies for which tabular data of the unsteady aerodynamic forces were available (Abel 1979). Present methods, that optimize these nonlinear lag terms, give better results with a decreased order in the approximating equation, and thus a solution of less complexity and increased efficiency.

6.2 Constant Optimization

A method is required that can optimally choose the coefficients of a governing rational function, to give the best possible approximation of the unsteady aerodynamics. This optimization must keep the computational burden to a minimum , have the ability to distinguish between global minimum and local minimum, as well as converge to the same solution from varying starting values. The method must also require no derivatives, since the objective function is of considerable complexity. The result is a two stage function optimization. Optimization of the minimization of the objective function (error function) was performed by utilizing a least square method to determine the *linear coefficients* and a sequential simplex method to determine the nonlinear lag terms.

Linear Coefficients : A_0 , A_1 , ... A_6

The linear coefficients are determined by utilizing a least square method, based on the minimization of a given error function. Many different forms of error functions are investigated (Appendix A) to determine the one that led to the best overall approximation of the unsteady aerodynamics. They consider slight modifications in error calculation between real and approximate data and incorporate various forms of normalizing techniques. The one that seemed to give the best results (Adams and Tiffany 1988) is

$$err(ik) = \sum \frac{|\bar{Q}(ik) - Q(ik)|^2}{M(ik)}$$
, (73)

where

$$M(ik) = max\left\{1, |Q(ik)|^2\right\}$$

The quantities $\bar{Q}(ik)$ and Q(ik) represent the rational function approximation and the true tabular data at a given reduced frequency, respectively. The term M(ik) is used to normalize the aerodynamic data, such that certain points do not receive larger than normal weighting. Note, the error function is defined as the total normalized sum for all the given values of the unsteady aerodynamic tabular data.

In order to obtain a minimum error, the derivative of the coefficients A_j , j = 0 to 6, with respect to the error function is taken

$$\frac{\partial err(ik)}{\partial(A_j)} = 0 . (74)$$

The result is a set of linear algebraic equations (see Appendix B), that can be written in the following form,

$$[A] \{x\} = [B] , (75)$$

where

$$[A] = \sum \frac{1}{M(ik)} \begin{bmatrix} 1 & 0 & -k^2 & B_1 & \cdots & B_4 \\ 0 & k^2 & 0 & B_1b_1 & \cdots & B_4b_4 \\ -k^2 & 0 & k^4 & -k^2B_1 & \cdots & -k^2B_4 \\ B_1 & B_1b_1 & -k^2B_1 & B_1B_1(1+\frac{b_1b_1}{k^2}) & \cdots & B_1B_4(1+\frac{b_1b_4}{k^2}) \\ B_2 & B_2b_2 & -k^2B_2 & B_2B_1(1+\frac{b_2b_1}{k^2}) & \cdots & B_2B_4(1+\frac{b_2b_4}{k^2}) \\ B_3 & B_3b_3 & -k^2B_3 & B_3B_1(1+\frac{b_3b_1}{k^2}) & \cdots & B_3B_4(1+\frac{b_3b_4}{k^2}) \\ B_4 & B_4b_4 & -k^2B_4 & B_4B_1(1+\frac{b_4b_1}{k^2}) & \cdots & B_4B_4(1+\frac{b_1b_4}{k^2}) \end{bmatrix}$$

$$[B] = \sum \frac{1}{M(ik)} \left[F, kG, -k^2F, FB_1 + \frac{GB_1b_1}{k}, \cdots, FB_4 + \frac{GB_4b_4}{k} \right]^T$$
$$\{x\} = \{A_0, A_1, A_2, \cdots, A_6\},$$

and

$$B_j = \frac{k^2}{k^2 + b_j^2}$$
, $j = 1$ to 4.

These equations can be solved quite easily by readily available computer subroutines, for matrix manipulation and eigenvalue analysis.

The objective error function is calculated by considering the approximate function, Q(ik), and the actual tabular data, Q(ik). Thus, given a set of lag terms, b_l , (l = 1 to 4), the coefficients A_j , (j = 0 to 6), can be found. The values for the *linear coefficients* obtained will result in a minimum error, but it will not be an optimal. There exists a particular set of values for the nonlinear lag terms that will result in an overall optimized minimal error function; this must be found as follows.

Nonlinear Coefficients : b_1 , b_2 , b_3 , b_4

The method employed to solve for the optimal nonlinear lag terms is a sequential simplex method (Mead and Nelder 1965; Nelson and Olsson 1975) developed by Nelder and Mead in 1965. This method is simple, robust and requires no derivatives, and hence lends itself quite well to finding the minimum of a nonlinear objective function of more than one independent variable (see Appendix C). The only assumption made is that the function does have a minimum and that the surface is continuous. Its only downfall is that it is not very efficient in the number of function evaluations that it requires. This method requires the function evaluation at (n + 1) vertices for a function of n variables. The largest function value is then replaced by a new point that is decided by reflection, extension or contraction onto the space. This process continues until convergence occurs. The simplex procedure has the ability to adapt itself to any given contour, it does not get easily fooled by a possible dead end. The beauty of the method is that it approaches the function minimum by moving away from the higher function values (see Appendix C for more details).

This method is quite self-contained (FORTRAN code: Olson 1974; O'Neil 1971), requiring the user to supply only the objective (error) function and the initial start and step values. This method also has the ability to consider self imposed boundary conditions. The sequential simplex procedure is a direct method that requires no derivatives of the objective function. In this study, since the objective function is quite complicated, it would be very difficult to supply such derivatives.

Powell's method may also be considered as an attractive alternative to the simplex method. This is a direction set method that produces N mutually conjugate directions

to decide on its path of solution. Problems with large computational burdens involving a large number of function evaluations will surely be faster using Powell's method, but the simplex adds an added robustness that seems to be lacking in Powell's method (Fletcher and Powell 1963; Flannery, Press, Teukolsky and Vetterling 1986). Powell's method has a tendency to give incorrect values if initial guess values result in a simplex that covers more than one valley.

With any method, there is still a chance of false convergence. The results obtained can be checked by choosing different starting points and repeating the process, or/and by letting the process continue after the initial convergence to see if it is indeed a minimum.

Boundary Conditions

The rational function used to approximate any given tabular set of aerodynamic data must obey several important boundary conditions. In order to be solved analytically, none of the lag terms can be equivalent to one another. If two lag terms are equal, the resulting equation $\bar{Q}(ik)$ could be reduced to a six-term equation versus the original seven-term equation. The lag terms must also be greater than zero (non-negative) to ensure system stability (Poirel 1988). This forces the poles of the resulting function to be in the left half-plane of the Laplace domain. If poles are found to be in the right half-plane, this implies an unstable governing function. The lag terms usually take on values in the range of reduced frequencies over which tabular data is available. Note, that in previous work, before these terms were optimized, they were chosen arbitrarily over this given reduced frequency range. Boundary conditions can be handled quite easily in the computer algorithm by replacing an incorrect (or undesirable) value with a new (allowable) one, i.e. a negative number, with, say, 0.00001.

Constraints

It is often required to place constraints on the approximating function. For example, constraints could be imposed at the reduced frequency k=0.0, forcing the approximating rational function to agree exactly with the tabular data at this point. This constraint results in a better fit at k=0.0 (an exact fit), but such a condition leads to an overall poorer fit. It may also be desirable to place constraints at a reduced frequency close to the expected flutter frequency. However, any imposed constraints lead to a loss in the degrees-of-freedom for the least-squares solution. In order to obtain an accurate fit to the given tabular data with imposed constraints, it is necessary to choose an adequate order of the approximating function. An increase in the order of the function (increase in number of lag terms) will definitely give better results, but with an increase in computational time.

An analysis was considered to determine the optimal number of lag terms required to give the desired accuracy with a compatible amount of computational time. This was done by looking at the approximation error obtained by the rational function, and by comparing flutter results obtained from the Laplace method to those obtained from the U - g and p - k analysis; both for the incompressible and transonic regime. If necessary, constraints are applied to the aerodynamic derivatives at a reduced frequency of k = 0.0, in hope of achieving improved results.

6.3 Aeroelastic Equations in the Laplace Domain

In this section the equations of motion for both incompressible and transonic flow will be transformed into a general Laplace domain solution (see Table 2) of the following form

$$([M] s2 + [C] s + [K] + x [N])\tilde{z}(s) = 0 , \qquad (76)$$

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where [M], [C] and [K] are the mass matrix, damping matrix and stiffness matrix, respectively. The matrix [N] consists of all the terms that are a function of reduced frequency, k, and is approximated by the Nasa-Langley rational function. The constant value x is simply a factor-multiplier that depends on the flow regime, incompressible or transonic. The term \hat{z} is the laplace operator of z where

$$z = \{ \xi \ \alpha \}^T .$$

The approximation matrix [N] given by

$$[N] = \left[\bar{Q}(p)\right] = [A_0] + [A_1]p + [A_2]p^2 + [A_3]\frac{p}{p+b_1} + [A_4]\frac{p}{p+b_2} + [A_5]\frac{p}{p+b_3} + [A_6]\frac{p}{p+b_4}.$$
 (77)

can be easily written in the Laplace domain by using the expression for reduced frequency

$$k=\frac{wb}{U},$$

along with

C

$$s = iw$$
 and $p = ik$,

to give

$$p = \frac{c}{2U}s = \frac{b}{U}s . ag{78}$$

The approximate rational function in Laplace form is then

$$\left[\bar{Q}(s)\right] = \left[A_{0}\right] + \frac{b}{U}\left[A_{1}\right]s + \left(\frac{b}{U}\right)^{2}\left[A_{2}\right]s^{2} + \left[A_{3}\right]\frac{s}{s + \frac{U}{b}b_{1}} + \left[A_{4}\right]\frac{s}{s + \frac{U}{b}b_{2}} + \left[A_{5}\right]\frac{s}{s + \frac{U}{b}b_{3}} + \left[A_{6}\right]\frac{s}{s + \frac{U}{b}b_{4}}.$$
(79)

6.3.1 Incompressible Aerodynamics

The general equations of motion of a two degree-of-freedom airfoil (Eq. 33,34) are given by

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + w_{h}^{2}\xi = \frac{Q_{h}}{mb} ,$$
$$r_{\alpha}^{2}\ddot{\alpha} + x_{\alpha}\ddot{\xi} + w_{\alpha}^{2}r_{\alpha}^{2}\alpha = \frac{Q_{\alpha}}{mb^{2}} .$$

For the particular case of incompressible aerodynamics these can be written as

$$\ddot{\xi}(1+\frac{1}{\mu}) + \dot{\xi}(\frac{2C(k)U}{b\mu}) + \xi(w_h^2) + \ddot{\alpha}(\frac{U}{b\mu} + 2C(k)(\frac{1}{2}-a)\frac{U}{b\mu}) + \alpha(\frac{2C(k)U^2}{b^2\mu}) = 0 , \qquad (80)$$

and

$$\ddot{\xi}(x_{\alpha} - \frac{a}{\mu}) + \dot{\xi}\left(-2(a + \frac{1}{2})C(k)\frac{U}{b\mu}\right) + \ddot{\alpha}(r_{\alpha}^{2} + \frac{\frac{1}{8} + a^{2}}{\mu}) + \dot{\alpha}\left(\frac{U}{b\mu}(\frac{1}{2} - a) - 2(a + \frac{1}{2})(\frac{1}{2} - a)C(k)\frac{U}{b\mu}\right) + \alpha\left(r_{\alpha}^{2}w_{\alpha}^{2} - 2(a + \frac{1}{2})C(k)\frac{U^{2}}{b^{2}\mu}\right) = 0.$$
(81)

Rewriting these equations, by taking all terms that are a function of C(k), to the right hand side gives

$$\ddot{\xi}(1+\frac{1}{\mu}) + \xi(w_h^2) + \ddot{\alpha}(x_{\alpha} - \frac{a}{\mu}) + \dot{\alpha}(\frac{U}{b\mu}) = \\ \dot{\xi}(-\frac{2C(k)U}{b\mu}) + \dot{\alpha}\left(-2C(k)(\frac{1}{2} - a)\frac{U}{b\mu}\right) + \alpha(\frac{-2C(k)U^2}{b^2\mu}) , \qquad (82)$$

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$$\ddot{\xi}(x_{\alpha} - \frac{a}{\mu}) + \ddot{\alpha}(r_{\alpha}^{2} + \frac{\frac{1}{8} + a^{2}}{\mu}) + \dot{\alpha}\left(\frac{U}{b\mu}(\frac{1}{2} - a)\right) + \alpha(r_{\alpha}^{2}w_{\alpha}^{2}) = \\ \dot{\xi}\left(2(a + \frac{1}{2})C(k)\frac{U}{b\mu}\right) + \dot{\alpha}\left(2(a + \frac{1}{2})(\frac{1}{2} - a)C(k)\frac{U}{b\mu}\right) + \alpha\left(2(a + \frac{1}{2})C(k)\frac{U^{2}}{b^{2}\mu}\right) .$$
(83)

The right hand side of both these equations is approximated by the Nasa-Langley expression (Eq. 79) in matrix form.

The Laplace transform of Equations (82) and (83), allowing for the initial conditions

$$\xi(0) = 0 \quad \alpha(0) = 0$$

 $\dot{\xi}(0) = 0 \quad \dot{\alpha}(0) = 0$, (84)

gives

$$([M]s2 + [C]s + [K] + [D]s + [E])\tilde{z}(s) = 0.$$
(85)

where \dot{z} is the Laplace operator of z where

 $z = \left\{ \begin{array}{cc} \xi & \alpha \end{array} \right\}^T ,$

and

$$[D] = 2\left(\frac{1}{\mu}\left(\frac{U}{b}\right)\right)C(k) \begin{bmatrix} 1 & (\frac{1}{2}-a) \\ -(a+\frac{1}{2}) & -(a+\frac{1}{2})(\frac{1}{2}-a) \end{bmatrix},$$
(86)

$$[E] = 2\left(\frac{1}{\mu}\left(\frac{U}{b}\right)^2\right)C(k)\begin{bmatrix}0 & 1\\ 0 & -(a+\frac{1}{2})\end{bmatrix},$$
(87)

An approximation is required for [D] s + [E].

$$\bar{Q}(s) = [D]s + [E] \tag{88}$$

By using the following substitution

$$s = \frac{U}{b}p = \frac{U}{b}ik$$

the approximation can be written as a single matrix $[\bar{N}]$, that is a function of C(k) and ikC(k).

$$\left[\bar{N}\right] = \left[D\right] \frac{U}{b} \iota k + \left[E\right]$$

where

$$\left[\bar{N}\right] = 2\left(\frac{1}{\mu}\left(\frac{U}{b}\right)^2\right)C(k) \begin{bmatrix} ik & (\frac{1}{2}-a)ik+1\\ -(a+\frac{1}{2})ik & -(a+\frac{1}{2})(\frac{1}{2}-a)ik-(a+\frac{1}{2}) \end{bmatrix} .$$
 (89)

or

$$\left[\bar{N}\right] = x\left[N\right] \,\,, \tag{90}$$

where

$$x = 2\left(\frac{1}{\mu}\left(\frac{U}{b}\right)^2\right) \;,$$

and

$$[N] = C(k) \begin{bmatrix} ik & (\frac{1}{2} - a)ik + 1 \\ -(a + \frac{1}{2})ik & -(a + \frac{1}{2})(\frac{1}{2} - a)ik - (a + \frac{1}{2}) \end{bmatrix}.$$
 (91)

Thus, the Laplace transform of the equations of motion in the general form (Eq. 76), are given by

$$([M] s2 + [C] s + [K] + x [N])\tilde{z}(s) = 0 .$$

6.3.2 Transonic Aerodynamics

For the particular case of transonic aerodynamics, equations (33) and (34), may conveniently be written as

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + w_h^2\xi = -\frac{qc}{mb}(C_{l_h}\frac{\xi}{2} + C_{l_a}\alpha) ,$$
$$r_{\alpha}^2\ddot{\alpha} + x_{\alpha}\ddot{\xi} + w_{\alpha}^2r_{\alpha}^2\alpha = \frac{qc^2}{mb^2}(C_{m_h}\frac{\xi}{2} + C_{m_a}\alpha) .$$

Allowing for the same initial conditions as used for the incompressible aerodynamic example, the transonic equations of motion (see Table 2) can be written in the Laplace domain as

$$([M] s2 + [C] s + [K] + x [N])\tilde{z}(s) = 0$$

where

C

$$x = \frac{1}{\pi} \left(\frac{1}{\mu} \left(\frac{U}{b} \right)^2 \right) \; ,$$

and

$$[N] = \begin{bmatrix} C_{l_h}/2 & C_{l_{\alpha}} \\ -C_{m_h} & -2C_{m_{\alpha}} \end{bmatrix} .$$
(92)

Incompressible	Transonic
$[M]$ $\begin{bmatrix} 1+\frac{1}{\mu} & x_{0}-\frac{a}{\mu} \\ x_{\alpha}-\frac{a}{\mu} & r_{0}^{2}+(\frac{1}{8}+a^{2})/\mu \end{bmatrix}$	$\begin{bmatrix} M \end{bmatrix} \\ \begin{bmatrix} 1 & x_{\alpha} \\ x_{\alpha} & r_{\alpha}^{2} \end{bmatrix}$
$\begin{bmatrix} C \end{bmatrix}_{a \in rodynamic} \\ \begin{bmatrix} 0 & \frac{1}{\mu} \frac{U}{b} \\ 0 & (\frac{1}{2} - a) \frac{1}{\mu} \frac{U}{b} \end{bmatrix}$	$\begin{bmatrix} C \end{bmatrix}_{aerodynamic} \\ \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$
$\begin{bmatrix} C \end{bmatrix}_{structural} \\ \begin{bmatrix} 2\zeta_h w_h & 0 \\ 0 & 2r_\alpha^2 \zeta_\alpha w_\alpha \end{bmatrix} \\ \zeta_h, \zeta_\alpha \approx 0.01 \end{bmatrix}$	$\begin{bmatrix} C \end{bmatrix}_{structural} \\ \begin{bmatrix} 2\zeta_h w_h & 0 \\ 0 & 2r_\alpha^2 \zeta_\alpha w_\alpha \end{bmatrix} \\ \zeta_h, \zeta_\alpha \approx 0.01 \end{bmatrix}$

 $([M] s^{2} + [C] s + [K] + x [N])\dot{z}(s) = 0$

note: $[C] = [C]_{aerodynamic} + [C]_{structural}$

Table 2 (continued)

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Incompressible	Transonic					
$\begin{bmatrix} K \end{bmatrix} \\ \begin{bmatrix} w_h^2 & 0 \\ 0 & r_{\alpha}^2 w_{\alpha}^2 \end{bmatrix}$	$\begin{bmatrix} K \end{bmatrix} \\ \begin{bmatrix} w_h^2 & 0 \\ 0 & r_o^2 w_o^2 \end{bmatrix}$					
$\frac{x}{2\left(\frac{1}{\mu}\left(\frac{l^{*}}{b}\right)^{2}\right)}$	$\frac{x}{\frac{1}{\pi}}\left(\frac{1}{\mu}\left(\frac{U}{b}\right)^2\right)$					
$\begin{bmatrix} N \end{bmatrix}$ $C'(k) \begin{bmatrix} ik & (\frac{1}{2} - a)ik + 1 \\ -(a + \frac{1}{2})ik & (a^2 - \frac{1}{4})ik - (a + \frac{1}{2}) \end{bmatrix}$	$\begin{bmatrix} N \end{bmatrix}$ $\begin{bmatrix} C_{l_h}/2 & C_{l_{\alpha}} \\ -C_{m_h} & -2C_{m_n} \end{bmatrix}$					
approximation of $C(k), \ C'(k)ik$	approximation of $C_{l_h}, C_{l_{\alpha}}, C_{m_h}, C_{m_{\alpha}}$					

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Table 2Laplace Solution : Matrices in Laplace Form for Incompressibleand Transonic Flow

6.4 Solution of the Aeroelastic Equations in the Laplace Domain

The solution to the aeroelastic equations of motion in the Laplace domain follows the work done previously by Abel (1979). Consider the general form of the aeroelastic equations of motion

$$([M] s2 + [C] s + [K] + x [N]) \hat{z}(s) = 0.$$

where

$$[N] = [A_0] + \frac{b}{U} [A_1] s + (\frac{b}{U})^2 [A_2] s^2 + [A_3] \frac{s}{s + \frac{U}{b}b_1} + [A_4] \frac{s}{s + \frac{U}{b}b_2} + [A_5] \frac{s}{s + \frac{U}{b}b_3} + [A_6] \frac{s}{s + \frac{U}{b}b_4}.$$

This gives

$$([M] s^{2} + [C] s + [K] + x \{ [A_{0}] + \frac{b}{U} [A_{1}] s + (\frac{b}{U})^{2} [A_{2}] s^{2} + [A_{3}] \frac{s}{s + \frac{U}{b}b_{1}} + [A_{4}] \frac{s}{s + \frac{U}{b}b_{2}} + [A_{5}] \frac{s}{s + \frac{U}{b}b_{3}} + [A_{6}] \frac{s}{s + \frac{U}{b}b_{4}} \}) \dot{z}(s) = 0 .$$
(93)

Equation (93) is then eliminated of its denominator terms to obtain the following polynomial

$$\left(B_{0}\left[\left[\bar{A}\right]s^{2}+\left[\bar{B}\right]s+\left[\bar{C}\right]\right]+B_{1}\left[\bar{D}\right]+B_{2}\left[\bar{E}\right]+B_{3}\left[F\right]+B_{4}\left[\bar{G}\right]z(s)=0.$$
 (94)

where

$$B_{0} = (s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{1} = s(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{2} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{3} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{4}) ,$$

$$B_{4} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3}) ,$$

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$$\begin{bmatrix} \bar{A} \end{bmatrix} = [M] + x(\frac{b}{U})^2 [A_2]$$
$$\begin{bmatrix} \bar{B} \end{bmatrix} = [C] + x\frac{b}{U}[A_1] ,$$
$$\begin{bmatrix} \bar{C} \end{bmatrix} = [K] + x [A_0] ,$$
$$\begin{bmatrix} \bar{D} \end{bmatrix} = x [A_3] ,$$
$$\begin{bmatrix} \bar{E} \end{bmatrix} = x [A_4] ,$$
$$\begin{bmatrix} \bar{F} \end{bmatrix} = x [A_5] ,$$
$$\begin{bmatrix} \bar{G} \end{bmatrix} = x [A_6] .$$

Finally, this may be written as a 6^{th} order polynomial in s of the following form

$$[F_6]s^6 + [F_5]s^5 + [F_4]s^4 + \dots + [F_1]s + [F_0] = 0.$$
(95)

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where the coefficients $[F_6]$, $[F_5]$, $[F_4]$, \cdots , $[F_0]$ are functions of dynamic pressure (velocity).

This can be easily placed in the form of a typical eigenvalue problem

$$s\{x\} = [A]\{x\} , (96)$$

to give a series of 6n first order equations, where

$$\{x\} = \{s^5, s^4, s^3, s^2, s^1, s^0\}$$
(97)

and

$$[A] = \begin{bmatrix} -[F_6]^{-1}[F_5] & -[F_6]^{-1}[F_4] & \cdots & -[F_6]^{-1}[F_1] & -[F_6]^{-1}[F_0] \\ [I] & [0] & [0] & \cdots & [0] & [0] \\ [0] & [I] & \cdots & [0] & [0] \\ [0] & [0] & \cdots & [I] & [0] \end{bmatrix}$$

Formulating the equations of motion into a typical eigenvalue problem results in a robust efficient method of calculating the critical flutter speeds. The flutter characteristics are found by calculating the complex eigenvalues of [A] at various values of dynamic pressure (velocity) at a specific Mach number. The eigenvalues are the roots of the characteristic equation. The roots of the given problem will be a combination of complex conjugate pairs and real numbers

$$\beta \pm iw , \qquad (98)$$

where w and β are the frequency and damping, respectively. A 6th order degree polynomial will produce 12 roots. Of these, eight roots can be attributed to the dimension of the rational approximating function, and the remaining four roots are due to the aeroelastic modes of the airfoil. For stability, all the real parts of the roots, β , must be negative in sign. A positive real part denotes an instability, flutter. Flutter occurs when at a given velocity the frequency of the two complex conjugates begin to approach one another (coalesce). When one of the real parts, β , changes sign, flutter has occurred.

7 Results : Rational Function Optimization

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The validity of the Laplace solution for the aeroelastic equations depends greatly on how well the unsteady aerodynamics can be approximated. It is critical that the approximation be very close to the actual tabular data. Each approximant curve must follow, almost identically, every dip and twist of the original function, and give an adequate representation of both the real and imaginary parts of the aerodynamics - as a function of the reduced frequency, k. If insufficient agreement is obtained the end solution will be erroneous. If an accurate approximation of the unsteady aerodynamics can be found, the solution of the equations in the Laplace domain will eliminate the need for costly iteration, that is characteristic of the traditional U - g and p - k methods.

The unsteady aerodynamics for both the incompressible and transonic regime will be approximated by rational functions. The incompressible aerodynamics, found as a function of the Theodorsen Function, C(k) and C(k)ik, should not lead to much difficulty in funding accurate approximating functions, since the aerodynamics vary in a fairly smooth manner with reduced frequency. However, problems may be encountered in the transonic regime, where the unsteady transonic aerodynamic coefficients C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} are very irregular, demonstrating much twisting as reduced frequency is changed. At a transonic Mach number of M = 0.85, the aerodynamics are very irregular and should prove the most difficult to approximate. Thus, in order for the Laplace method to be viable, the approximations of these tabular unsteady aerodynamics cannot be considered lightly.

In order for the overall solution of the aeroelastic equations of motion to be ro-
bust, the unsteady aerodynamics should cover an adequate reduced frequency range. The governing criteria for such an analysis is that the largest frequency generated from the eigenvalue analysis should fall within the given range of frequency for the unsteady aerodynamic data. The tabular data for the unsteady aerodynamics has a lower value reduced frequency of k = 0.0, and the upper value can be checked by

$$k_{upper} = \frac{wb}{U} . \tag{99}$$

where b is the semi-chord and w is the largest frequency generated from the eigenvalue analysis at the specific velocity U.

For incompressible aerodynamics, approximations of both the Theodorsen function. C(k), and the Theodorsen function times the reduced frequency, C(k)/k are required. Initially, the optimization procedure, which used a combined least square methodology and a simplex method, considered only the Theodorsen function, C(k). This was undertaken in order to answer three fundamental questions. Firstly, the validity of the approximating function needed to be ascertained. Could the Nasa-Langley rational function lead to a valid approximation of the unsteady aerodynamics, and hence lead to a possible successful solution to the aeroelastic equations? Secondly, it was required to determine how many lag terms were needed to ensure an acceptable amount of error. Lastly, a study of the reduced frequency range was required to determine what range could be approximated successfully.

A valid approximation of the Theodorsen function, C(k), over the reduced frequency range, 0.0 < k < 10.0, was achieved by utilizing a Nasa-Langley rational function, consisting of four lag terms. The calculated optimized linear terms were given by $A_0 = 0.99858, A_1 = -0.000078, A_2 = 0.000012, A_3 = -0.040125,$

$$A_4 = -0.152297, A_5 = -0.22708, A_6 = -0.078005,$$

and the lag terms as

 $b_1 = 0.014919, b_2 = 0.080715, b_3 = 0.238540, b_4 = 0.687273,$

with a calculated error of 8.207×10^{-6} .

It was determined that the approximation error to C(k) decreased quite significantly with the increase in the number of lag terms, as shown in the table below.

Error	# of lag terms
0.02025	I
0.00093 60	2
0.00008485	3
0.000008207	-1

The Nasa-Langley rational function, incorporating four lag terms, gave an excellent approximation to the Theodorsen function over the entire reduced frequency range. This same methodology was then used to approximate C(k) and C(k)ik with great success. It should be noted that in order to simplify the subsequent analysis of the aeroelastic equations, C(k) and C(k)ik were forced to have the same lag terms.

Again, considering the same reduced frequency range, 0.0 < k < 10.0, the calculated optimized linear terms for C(k) and C(k)ik, respectively, were

 $A_0 = 0.99828, A_1 = -0.000035, A_2 = 0.0000047, A_3 = -0.041204,$

$$A_1 = -0.162893, A_5 = -0.22909, A_6 = -0.0645937,$$

and

 $A_0 = -0.00001355, A_1 = 0.50037475, A_2 = -0.00003629, A_3 = 0.00061076,$

$$A_4 = 0.01380489, A_5 = 0.05848218, A_6 = 0.04963449,$$

with the lag terms,

 $b_1 = 0.015441, b_2 = 0.084286, b_3 = 0.255951, b_4 = 0.765114,$

with a calculated error of 1.178×10^{-5} .

The Nasa-Langley rational function approximation with multi-level constraints and nonlinear term optimization proved to be quite useful. In the incompressible regime, the approximation error of the unsteady aerodynamics for C(k) and C(k)ik was found to be extremely small. Excellent agreement can be seen by referring to either the optimization table comparison (see Appendix D) or the graphs of C(k) and C(k)ik (see Figures 4a-1 to 4a-4 and Figures 4b-1 to 4b-4). Each comparison of the aerodynamic functions is given as two graphs, the real component and the imaginary component versus the reduced frequency. The choice of identical lag terms for the approximation of C(k) and C(k)ik, allows for a more efficient and robust solution of the aeroelastic equations in the Laplace domain. The given approximation to the unsteady aerodynamics gave useful results to the solution of the flutter problem, as will be seen in the next section.

In the transonic regime, approximations of four aerodynamic coefficients, namely,

 C_{l_h} , C_{l_a} , C_{m_h} and C_{m_a} , were required. This inevitably resulted in an increased approximation error. To make matters worse, the transonic aerodynamics are very irregular. This resulted in a slight modification to the way in which the unsteady aerodynamics were approximated. A decreased reduced frequency range was considered, and an analysis was undertaken to study the effect of imposed constraints at the reduced frequency, k = 0.0.

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Initially, an approximation of the unsteady aerodynamics for the transonic regime at M = 0.85 was considered and a reduced frequency range, $0.025 \le k \le 1.0$ was studied. The optimization of the aerodynamics appeared to be quite good, as can be seen from Figures 5a-1 to 5d-1; however, these approximations resulted in erroneous solutions for the flutter velocity when employed in the Laplace solution. The values of the unsteady aerodynamics at low values of reduced frequency, particularly in the interval of $0.0 \le k \le 0.025$, seemed to be of utmost importance in the analysis of the aeroelastic equations. If no consideration is given to the aerodynamics at the low reduced frequency range, the approximation procedure may give values of the aerodynamics at a value of k, that have no physical significance; i.e., a negative value of C_{lh} or C_{m_h} at k = 0.0 when they should equal 0.0, or extraordinary large values for the aerodynamics at k = 0.0. These problems are related to the value chosen for the constant A_0 in the rational function approximation. Note that the value of each aerodynamic derivative at the k = 0.0 condition is equal the value of A_0 (see Appendix E).

To solve these problems the original reduced frequency range of $0.025 \le k \le 1.0$, had to be changed to include the reduced frequency of k = 0.0. Hence, a reduced frequency range of $0.0 \le k \le 1.0$ was considered, as well as a decreased range of $0.0 \le k \le 0.5$. Imposed constraints were considered at k = 0.0 for all four aerodynamic coefficients, and as well as just for those that were zero at k = 0.0, namely C_{l_h} and C_{m_h} . All approximations of the unsteady aerodynamics resulted in functions which gave valid solutions to the aeroelastic equations in the Laplace method, as will be seen in the next section. The optimization results for the different reduced frequency ranges, and the imposed constraints are presented in Table 3. These results, for the transonic regime, were not as good as those for the incompressible regime, but were still found to be excellent overall. Note, that the requirements for the approximation of the unsteady transonic aerodynamics were much more demanding than for the incompressible aerodynamics, as can be seen from the calculated error.

Reduced Frequency Range, k	Constraints at $k = 0.0$	Error
$0.0 \le k \le 1.0$	$C_{l_h}, C_{l_\alpha}, C_{m_h}, C_{m_\alpha}$	0.009355
$0.0 \le k \le 1.0$	C_{l_h}, C_{m_h}	0 .003 602
$0.0 \le k \le 0.5$	$C_{l_h}, C_{l_\alpha}, C_{m_h}, C_{m_\alpha}$	0.004183
$0.0 \le k \le 0.5$	C_{l_h}, C_{m_h}	0.00 1396

Table 3Transonic Regime : Comparison of Rational Function Optimization Error for
Various Reduced Frequency Ranges and Imposed Constraints, M = 0.85.

A reduced frequency range of $0.0 \le k \le 0.5$, with imposed constraints at k = 0.0 for the unsteady aerodynamics, C_{l_h} and C_{m_h} , at M = 0.85, gave the minimum approximation error, and the best solution to the flutter problem, as will be shown in the next section. Excellent agreement can be seen by referring to either the optimization table comparison (see Appendix F) or the graphs of the transonic aerodynamic coefficients (see Figures 5a-2 to 5d-2). Each comparison is presented as two graphs, the real and imaginary components as functions of the reduced frequency. The decreased reduced frequency range resulted in a less robust total test range; thus, in the solution of the aeroelastic equations, some results obtained at low values of velocity were somewhat erratic. Thus a check was required to determine if the instability was totally within the accepted confines of the given unsteady tabular data range.

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Applying the same reduced frequency range, $0.0 \le k \le 0.5$, with forced constraints at k = 0.0 for C_{l_h} and C_{m_h} , the approximations to the unsteady aerodynamics for M =0.80 and M = 0.875 were also found to be very good. These results are graphically presented in Figures 6a to 6d and Figures 7a to 7d, respectively.

It should be noted that when approximations were made for the unsteady aerodynamics, a complete global error was considered. Hence, for the incompressible regime, the coefficients for the rational function were calculated simultaneously considering both C(k) and C(k)ik. The approximations were found such that they resulted in a minimum total calculated error for the same nonlinear coefficients. The transonic regime is the most interesting, since a global minimization of the four aerodynamic coefficients, C_{l_n} , C_{l_n} , C_{m_h} and C_{m_n} was performed. It should be noted, once more, that this global approximation error approach was considered at M = 0.85, where the nonlinearities are the most pronounced, and still gave adequate results.

The Nasa-Langley solution for approximating unsteady aerodynamics gave excellent results for both the incompressible and transonic regime. This method demonstrates comparably better results to previously used methods with less overall computational time and complexity (Poirel 1988).

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8 Results : Laplace Method of Solution for the Aeroelastic Equations

The rational functions obtained for incompressible flow were substituted into the twodegree-of-freedom airfoil equations. The corresponding Laplace eigenvalue problem was solved, producing twelve eigenvalues (see Appendix G). Of these, eight were completely real, attributed to the dimension of the rational approximating function, and the remaining four eigenvalues produced two complex conjugate pairs, due to the aeroelastic modes of the airfoil. The p - k solution generated four roots, which formed two complex conjugate pairs. This allowed for the direct comparison of the p - k solution with the Laplace method, when the real roots were omitted. The U - g method produced only two roots, which resembled those of the positive frequency conjugate roots. Various case studies were considered, incorporating the various methods of flutter solution.

For example, each method was used to solve the equations for the following set of parameters; $\mu = 50.0$, $r_{\alpha} = 0.5$, a = -0.5 and $w_h/w_{\alpha} = 0.2$ (Case Study 1). The results are presented as graphs of frequency ratio, w/w_{α} , and damping ratio, versus nondimensional velocity, \bar{U} (see Figures 8a-1 to 8a-3). All the methods resulted in exactly the same value for the critical flutter velocity, $\bar{U} = 4.53$ (see table 4); also, all the methods gave similar values of frequency ratio, w/w_{α} , over the non-dimensional velocity range, \bar{U} . However, the U - g method showed some variance in the damping ratio compared to either the p - k or Laplace methods. A second example was considered, which led to similar results. In this example, ($\mu = 50.0$, $r_{\alpha} = 0.6$, a = -0.6 and $w_h/w_{\alpha} = 0.4$ (Case Study 2)), all methods produced a flutter velocity, \bar{U} , of approximately 5.10. This example, showed similar response behaviour in frequency and damping ratio versus nondimensional velocity as in the previous one (see Figures 8b-1 to 8b-3). Other examples, showing comparisons of critical flutter velocity, obtained from the various methods of solution are shown in Table 4. It should be noted that the U - g method encountered some numerical difficulty in some of the cases studied, namely case studies 4 and 5. The p - k and the Laplace methods had no such problems.

Case				Flutter Velocity $ar{U}$			
Study	μ	r_{α}	a	w_h / w_{α}	U - g	p-k	Laplace
1	50.0	0.5	-0.5	0.2	4.53	4.53	4.53
2	50.0	0.6	-0.6	0.4	5.10	5.10	5.11
3	100.0	0.5	-0.5	0.2	6.24	6.26	6.26
4	75.0	0.4	-0.4	0.3	-?-	3.68	3.68
5	100.0	0.4	-0.4	0.3	-?-	4.16	4.16

Table 4 Incompressible Regime : Comparison of critical flutter velocity, \bar{U} , using the U - g, p - k and Laplace methods. The symbol + ? - indicates that it was not possible to determine a flutter velocity for these conditions due to numerical instabilities.

For transonic flow, the aeroelastic equations of motion were also solved utilizing the rational function approximation of the unsteady aerodynamics. The given Laplace eigenvalue solution produced a total of twelve eigenvalues (see Appendix H). Of these, two real roots and three complex conjugate pairs can be attributed to the rational approximating function, and the other two complex conjugate pairs are from the aeroelastic modes of



the autoil When the aeroelastic eigenvalues were isolated from the other eigenvalues, not always an easy task, comparisons can be made with the U - g and p - k methods.

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The traditional methods of solution for the aeroelastic equations were not formulated for transonic flow, but comparisons were made with U - g and p - k data from Lee (1984). The Laplace solution was formulated for transonic flow at a M = 0.85, for the parameters: $\mu = 50.0, r_{\alpha} = 0.5, a = -0.5, w_h/w_{\alpha} = 0.2$ and $x_{\alpha} = 0.25$. This allowed for direct comparison with the U - g and p - k generated results given by Lee (1984).

An initial reduced frequency range of $0.0 \le k \le 1.0$, with forced constraints at k = 0.0 for all of the aerodynamic derivatives, $C_{l_{\alpha}}$, $C_{m_{\alpha}}$ and $C_{m_{\alpha}}$ was considered. This yielded the following results, $\bar{U} = U/bw_{\alpha} = 3.36$, $w/w_{\alpha} = 0.302$ and k = wb/U = 0.090.

These results did not compare as well as had been hoped with those of Lee (1984), which are $U = U/bw_{\alpha} = 3.46$, $w/w_{\alpha} = 0.313$ and k = wb/U = 0.090.

It was concluded that better results could be obtained if the reduced frequency range was decreased to $0.0 \le k \le 0.5$ this allowed for a better approximation of the low-range reduced frequency aerodynamic. This region, say from $0.0 \le k \le 0.1$ seemed to be the most critical in the overall flutter analysis scheme, since the expected critical flutter frequency. k = 0.090, fell within this range. As well as considering this decreased reduced frequency range, constraints were placed only on the aerodynamic derivatives, whose magnitude at k = 0.0 was 0.0; i.e., C_{l_h} and C_{m_h} . The following results were obtained, $\bar{U} = U/bw_{\alpha} = 3.45$, $w/w_{\alpha} = 0.309$ and k = wb/U = 0.090.

These results agreed favourably with those given by Lee (1984). They are presented graphically in Figures 9a-1 and 9a-2. Figure 9a-1 shows excellent agreement for both the

damping and frequency ratio as a function of non-dimensional velocity, U/bw_{γ} . However, in Figure 9a-2, there is some variance observed in the magnitude of the damping ratio, particularly for values of low non-dimensional velocity. This was due to the imaginary part of the eigenvalue falling out of the considered reduced frequency range, namely, $0.0 \le k \le 0.5$, and thus out of the scope of the approximating aerodynamic function This however, is not so critical; as the root that does go unstable is well behaved, and the other root shows variation from the expected behaviour only at low values of non dimensional velocity.

For completeness, other results incorporating varying reduced frequency ranges and constraints are included below, in Table 5.

Reduced Frequency Range	Constraints applied at $k = 0.0$	Ľ	w/w_{ϕ}	k
$0.0 \le k \le 1.0$	$C_{l_h}, C_{l_a}, C_{m_h}, C_{m_a}$	3.36	0.302	0 090
$0.0 \le k \le 1.0$	C_{l_h}, C_{m_h}	3.19	0 297	0.093
$0.0 \le k \le 0.5$	$C_{l_h}, C_{l_a}, C_{m_h}, C_{m_a}$	3.31	0.302	0.091
$0.0 \le k \le 0.5$	C_{l_h}, C_{m_h}	3.45	0.309	0.090

Table 5Transonic Regime : Comparison of Laplace Solution for Various ReducedFrequency Ranges and Imposed Constraints for M = 0.85.

The Laplace Solution yielded the best results, when the approximations of the unsteady aerodynamics considered a reduced frequency range of $0.0 \le k \le 0.5$, with forced constraints for C_{l_h} and C_{m_h} at k = 0.0. Using these approximations of the unsteady aerodynamics, the Laplace solution was formulated for varying values of non-dimensional autfoil mass, μ , at M = 0.85. This allowed for the direct comparison with U - g and p - k results given by Lee (1984) (see Table 6). Excellent agreement was obtained over the entire range of airfoil non-dimensional mass, μ , such that the percentage error between the non-dimensional velocities, \bar{U} , obtained using the Laplace method and those given by Lee, was less than 1%.

Further analyses considered the transonic regime for Mach numbers of M = 0.80and M = 0.875 at varying values of airfoil non-dimensional mass, μ , (see Table 6). This showed that, although only a small interval range in Mach number was considered, 0.80 < M < 0.875, the overall behaviour was quite irregular. For example, an analysis of the aeroelastic equations of motion for the Mach numbers of M = 0.80, M = 0.85 and M = 0.875, at an airfoil non-dimensional mass, $\mu = 50.0$, gave critical flutter velocities of $\hat{U} = 3.70, 3.45$ and 6.17, respectively. If the flutter velocities obtained at M = 0.80 and M = 0.875, had been used to interpolate for the flutter velocity at M = 0.85, the actual value of $\hat{U} = 3.45$ would not have been acquired.

μ	$l^{\cdot} - g / p - k^{\bullet}$	Laplace	Laplace	Laplace
	Ī	Ū	ī.	Ū
	M = 0.85	M = 0.85	M = 0.80	M = 0.875
50.00	3.46	3.45	3.70	6.17
75.00	3.94	3.95	4.33	7.08
100.00	1.33	4.36	1.86	7.54
150.00	-4.99	5.01	5.76	7.41
200.00	5.55	5.53	6.53	7.37
250.00	5.99	5.96	7.21	7.35

* reference Lee 1984

Table 6 Transonic Regime : Comparison of critical flutter velocity, U, for M = 0.80, M = 0.85 and M = 0.875 using the U - g, p - k and Laplace methods, with the given constants : $r_{\alpha} = 0.5$, a = -0.5, $w_h/w_{\alpha} = 0.2$, $x_{\alpha} = 0.25$ and b = 1.0.

The optimized coefficients of the rational function approximations for M = 0.80, M = 0.85 and M = 0.875 are given in Appendix I. Again, from a comparison of the respective linear and nonlinear coefficient terms, it can be seen how quickly and irregularly the aerodynamics change from one Mach number to another.

The excellent comparison obtained for the flutter condition between the traditional methods and the Laplace method demonstrates the validity of the Laplace method. It also shows the excellent nature of the approximation for the unsteady aerodynamics obtained by utilizing a Nasa-Langley rational function with multi-level constraints and nonlinear term optimization. The main strength of the Laplace method is that the unsteady aerodynamics have only to be approximated once and then can be used in the analysis of many sets of parameters: thus, an iterative type solution is not required. The benefits will soon be realized if many conditions have to be studied. However, with the introduction of this approximating rational function the dimension of the problem is increased considerably. from four to twelve, thus increasing the size of the eigenvalue problem.

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9 Active Control for a Two-Dimensional Airfoil in Incompressible Flow

9.1 Equations of Motion Incorporating a Flap Deflection

Consider a three-degree-of-freedom rigid autfoil flexibly mounted undergoing bending and twisting (torsion) motion (see Figure 10). Bending (plunge) is denoted by h, positive downward, the twisting (pitch) about the elastic axis is denoted by α , positive in the clockwise direction, and the flap angle by β , positive in the clockwise direction. The accoclastic equations of motion (Fung 1955; Ashley and Bisplinghoff 1962), can be derived by summing forces and moments about the elastic axis and moments about the flap hinge line to give

$$m\ddot{h} + m(x_{\alpha}b)\ddot{\alpha} + m(x_{\beta}b)\ddot{\beta} = -K_{h}h - L , \qquad (100)$$

$$(I + m(x_{\alpha}b)^{2})\ddot{\alpha} + m(x_{\alpha}b)\ddot{h} + (I_{\beta} + (c_{\beta} - a)mb^{2}x_{\beta})\ddot{\beta} = -K_{\alpha}\alpha + M_{\alpha} , \qquad (101)$$

$$(I_{\beta} + (c_{\beta} - a)mb^{2}x_{\beta})\ddot{\alpha} + m(x_{\beta}b)\ddot{h} + (I_{\beta})\ddot{\beta} = -K_{\beta}\beta + M_{\beta}, \qquad (102)$$

where I_{β} is the flap mass moment of inertia about the flap hinge line, $x_{\beta}b$ is the distance from the flap hinge line to the flap centre of mass, $c_{\beta}b$ is the distance from the midchord to the flap hinge line, ab is the distance from the midchord to the elastic axis (measured positive towards trailing edge), K_{β} is the torsional spring constant about the flap hinge line, and M_{β} is the moment acting about the flap hinge line.

Introducing a non-dimensional generalized displacement coordinate, as before,

$$\xi = \frac{h}{b} \tag{103}$$

and making the following substitutions

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$$K_h = w_h^2 m . (104)$$

$$K_{\alpha} = w_{\alpha}^2 I_{\alpha} , \qquad (105)$$

$$K_{\beta} = w_{\beta}^2 I_{\beta} , \qquad (106)$$

$$I_{\alpha} = r_{\alpha}^2 m b^2 \,\,. \tag{107}$$

$$I_{\beta} = r_{\beta}^2 m b^2 . \tag{108}$$

where r_{α} and r_{β} are the radius of gyration about the elastic axis and the flap hinge line. respectively, the complete 3DOF aeroelastic equations of motion become

$$\ddot{\xi} + x_{\alpha}\ddot{\alpha} + x_{\beta}\ddot{\beta} + w_h^2\xi = \frac{-L}{mb} , \qquad (109)$$

$$r_{\alpha}^{2}\ddot{\alpha} + x_{\alpha}\ddot{\xi} + (r_{\beta}^{2} + (c_{\beta} - a)x_{\beta})\ddot{\beta} + w_{\alpha}^{2}r_{\alpha}^{2}\alpha = \frac{M_{\alpha}}{mb^{2}}.$$
(110)

$$(r_{\beta}^{2} + (c_{\beta} - a)x_{\beta})\ddot{\alpha} + x_{\beta}\ddot{\xi} + r_{\beta}^{2}\ddot{\beta} + w_{\beta}^{2}r_{\beta}^{2}\beta = \frac{M_{\beta}}{mb^{2}}.$$
(111)

For incompressible flow, the expressions incorporating a flap deflection for the lift force, pitching moment and the hinge moment are (Theodorsen 1935; Garrick and Theodorsen 1941)

$$L = -\pi\rho b^3 w^2 \left\{ P_w \xi + \left[P_\phi - (\frac{1}{2} + a) P_w \right] \alpha + P_\beta \beta \right\}$$
(112)

$$M_{\alpha} = \pi \rho b^{4} w^{2} \left\{ \left[M_{w} - (\frac{1}{2} + a) P_{w} \right] \xi + \left[M_{\phi} - (\frac{1}{2} + a) (P_{\phi} + M_{w}) + (\frac{1}{2} + a)^{2} P_{w} \right] \alpha + \left[M_{\beta} - (\frac{1}{2} + a) P_{\beta} \right] \beta \right\} , \qquad (113)$$

and

$$M_{\beta} = \pi \rho b^{4} w^{2} \left\{ T_{u} \xi + \left[T_{\phi} - (\frac{1}{2} + a) T_{w} \right] \alpha + T_{\beta} \beta \right\}$$
(114)

The expressions P_w , P_ϕ , P_β , M_w , M_ϕ , M_∂ , T_w , T_ϕ and T_β are included for reference in Appendix J. They are a function of the Theodorsen function, C(k), velocity, frequency and geometry.

9.2 Flutter Analysis With Active Controls

Flutter-suppression control systems have received increased attention in the last decade (Karpel 1981; Kass and Thompson 1971). They have been found to be an effective way of increasing the critical flutter speed with substantial weight savings. There are two types of flutter-suppression systems that have proved to be quite useful: passive and active flutter control.

Passive Flutter Control

Passive flutter control considers increasing the structural stiffness and/or mass balancing. Proper wing mass distribution can improve the flutter characteristics tremendously. This is so because careful placing of mass has a direct effect on the inertial coupling of the system. Inertial coupling is very often the predominant factor in flutter. Increasing the in structural stiffness (particularly the torsional stiffness) will also result in an increased critical flutter speed. It has been shown that a stiffness increase of factor n will result in an increase in flutter speed of approximately \sqrt{n} . However, the addition of stiffness unfortunately means an increase in the total overall aircraft weight.

Active Flutter Control

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The type of control that is employed in this study is active flutter-suppression. Active control uses a control surface (flap) which is deflected in response to the wing motion, resulting in a change in the aerodynamic forces and moments on the wing. Such active control has to be very reliable and durable, because the loss of flutter control would probably result in an immediate airfoil failure. Another problem encountered in the use of active control is that of obtaining accurate measurements from the airfoil. Direct measurement is required for the controller to control the control surface via some specified control law. It is also important for the controller to be robust and allow for control at various flight conditions. The benefit of active control is that its implementation will have no significant increase in the airfoil weight.

Writing the equations of motion in the Laplace domain allows for the easy implementation of active control. Consider the equations of motion of a three-degree-of-freedom rigid airfoil incorporating a flap control, β . Neglecting the hinge moment equation, the equations of motion (100) and (101) can be written in matrix form

$$([M] s2 + [C] s + [K] + x [N])\tilde{z}(s) = 0 , \qquad (115)$$

where [M], [C], [K] and [N] are 2×3 matrices, and the term \tilde{z} is the laplace operator of z where

$$z = \{ \xi \ \alpha \ \beta \}^T .$$

Equation (115) may then be rewritten in the following form

$$([M_s \mid M_c] s^2 + [C_s \mid C_c] s + [K_s \mid K_c] + x [N_s \mid N_c]) \begin{cases} q \\ q_c \end{cases} = 0 .$$
(116)

where

$$\{q\} = \{ \xi \mid \alpha \}^T . \tag{117}$$

$$\{q_c\} = \{\beta\} . \tag{118}$$

The subscripts s and c denote a structural and control quantities, respectively, where $[M_s]$, $[C_s]$ and $[K_s]$ are 2 × 2 matrices and $[M_c]$, $[C_c]$ and $[K_c]$ are 2 × 1 matrices.

A control law is considered of the following form

$$\{ \beta \} = [T] \left\{ \begin{array}{c} \xi \\ \alpha \end{array} \right\} . \tag{119}$$

where [T] is a transfer-function matrix of size 1×2 . This allows the flap angle, β , to be written as a function of the generalized aerodynamic motions, ξ and α .

Applying the control law to (116) gives

$$([M_s] s^2 + [C_s] s + [K_s] + x [N_s]) \{q\} + ([M_c] s^2 + [C_c] s + [K_c] + x [N_c]) [T] \{q\} = 0.$$
(120)

The solution procedure is the same as that employed earlier in section 6.1. Making the following substitutions in (120),

$$[N_{s}] = [A_{0}] + \frac{b}{U} [A_{1}] s + (\frac{b}{U})^{2} [A_{2}] s^{2} +$$

$$[A_{3}] \frac{s}{s + \frac{U}{b}b_{1}} + [A_{4}] \frac{s}{s + \frac{U}{b}b_{2}} + [A_{5}] \frac{s}{s + \frac{U}{b}b_{3}} + [A_{6}] \frac{s}{s + \frac{U}{b}b_{4}} , \qquad (121)$$

and

$$[N_c] = [A_{0,c}] + \frac{b}{U} [A_{1,c}] s + (\frac{b}{U})^2 [A_{2,c}] s^2 +$$

$$[A_{3,c}] \frac{s}{s + \frac{L}{b}b_1} + [A_{4,c}] \frac{s}{s + \frac{L}{b}b_2} + [A_{5,c}] \frac{s}{s + \frac{L}{b}b_3} + [A_{5,c}] \frac{s}{s + \frac{L}{b}b_4} .$$
(122)

gives

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$$([M_{4}]s^{2} + [C_{5}]s + [K_{5}] + x\{[A_{0}] + \frac{b}{U}[A_{1}]s + (\frac{b}{U})^{2}[A_{2}]s^{2} + [A_{3}]\frac{s}{s + \frac{U}{b}b_{1}} + [A_{1}]\frac{s}{s + \frac{U}{b}b_{2}} + [A_{5}]\frac{s}{s + \frac{U}{b}b_{3}} + [A_{6}]\frac{s}{s + \frac{U}{b}b_{4}}\})\{q\} + ([M_{c}]s^{2} + [C_{c}]s + [K_{c}] + x\{[A_{0,c}] + \frac{b}{U}[A_{1,c}]s + (\frac{b}{U})^{2}[A_{2,c}]s^{2} + [A_{3,c}]\frac{s}{s + \frac{U}{b}b_{1}} + [A_{1,c}]\frac{s}{s + \frac{U}{b}b_{2}} + [A_{5,c}]\frac{s}{s + \frac{U}{b}b_{3}} + [A_{6,c}]\frac{s}{s + \frac{U}{b}b_{4}}\})[T]\{q\} = 0.$$
 (123)

The denominators of the above equations are removed to obtain the following polynomial

$$(B_0\left[\left[\bar{A}\right]s^2 + \left[\bar{B}\right]s + \left[\bar{C}\right]\right] + B_1\left[\bar{D}\right] + B_2\left[\bar{E}\right] + B_3\left[\bar{F}\right] + B_4\left[\bar{G}\right])\{q\} + \\(B_0\left[\left[\bar{A}_c\right]s^2 + \left[\bar{B}_c\right]s + \left[\bar{C}_c\right]\right] + B_1\left[\bar{D}_c\right] + B_2\left[\bar{E}_c\right] + B_3\left[\bar{F}_c\right] + B_4\left[\bar{G}_c\right])[T]\{q\} = 0,$$

where

$$B_{0} = (s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{1} = s(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{2} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{3})(s + \frac{U}{b}b_{4}) ,$$

$$B_{3} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{4}) ,$$

$$B_{4} = s(s + \frac{U}{b}b_{1})(s + \frac{U}{b}b_{2})(s + \frac{U}{b}b_{3}) .$$

$$\left[\bar{A}\right] = [M] + x\left(\frac{b}{U}\right)^2 [A_2] ,$$

$$\begin{bmatrix} \bar{B} \end{bmatrix} = \begin{bmatrix} C \end{bmatrix} + x \frac{b}{U} \begin{bmatrix} A_1 \end{bmatrix}$$
$$\begin{bmatrix} \bar{C} \end{bmatrix} = \begin{bmatrix} K \end{bmatrix} + x \begin{bmatrix} A_0 \end{bmatrix} ,$$
$$\begin{bmatrix} \bar{D} \end{bmatrix} = x \begin{bmatrix} A_3 \end{bmatrix} .$$
$$\begin{bmatrix} \bar{E} \end{bmatrix} = x \begin{bmatrix} A_4 \end{bmatrix} ,$$
$$\begin{bmatrix} \bar{F} \end{bmatrix} = x \begin{bmatrix} A_5 \end{bmatrix} ,$$
$$\begin{bmatrix} \bar{G} \end{bmatrix} = x \begin{bmatrix} A_6 \end{bmatrix} ,$$

and

$$\begin{bmatrix} \bar{A}_{c} \end{bmatrix} = [M_{c}] + x(\frac{b}{U})^{2} [A_{2,c}]$$
$$\begin{bmatrix} \bar{B}_{c} \end{bmatrix} = [C_{c}] + x\frac{b}{U} [A_{1,c}] .$$
$$\begin{bmatrix} \bar{C}_{c} \end{bmatrix} = [K_{c}] + x [A_{0,c}] ,$$
$$\begin{bmatrix} \bar{D}_{c} \end{bmatrix} = x [A_{3,c}] ,$$
$$\begin{bmatrix} \bar{E}_{c} \end{bmatrix} = x [A_{4,c}] ,$$
$$\begin{bmatrix} \bar{F}_{c} \end{bmatrix} = x [A_{5,c}] ,$$
$$\begin{bmatrix} \bar{G}_{c} \end{bmatrix} = x [A_{6,c}] .$$

This is then expanded to obtain a m^{th} order polynomial in s of the following form

$$[F_m] s^m + [F_{m-1}] s^{m-1} + [F_{m-2}] s^{m-2} + \dots + [F_1] s + [F_0] = 0 , \qquad (121)$$

where the coefficients $[F_m]$, $[F_{m-1}]$, $[F_{m-2}]$, \cdots , $[F_0]$ are functions of dynamic pressure (velocity).

This can be easily placed in a typical eigenvalue problem of the given form

$$s\{x\} = [A]\{x\} , \qquad (125)$$

to give a series of $m \times n$ first order equations, where

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$$\{x\} = \{ s^{m-1} \ s^{m-2} \ \cdots \ s^1 \ s^0 \} . \tag{126}$$

The method of solution continues as shown earlier (see Section 6.4). Thus, for a given transfer function matrix [T], the problem can be analyzed to determine if such a given control law has achieved an increase in the flutter speed.

10 Results : Active Control for Incompressible Flow

The problem associated with applying an active control device to an airfoil, hinges on developing a set of aeroelastic equations, where the form of the unsteady aerodynamics and the control law are compatible. Writing the equations in the Laplace domain is ideally suited to this task. Active control was considered only for the case of incompressible flow. Active control for transonic flow could be considered; however, it would first be necessary to modify the transonic aerodynamics to incorporate a flap.

A control law was considered of the following form

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$$\{\beta\} = [T] \left\{ \begin{array}{c} \xi\\ \alpha \end{array} \right\} \,. \tag{127}$$

where [T] is the transfer-function matrix of size 1×2 . This allows the flap angle, β , to be written as a function of the generalized aerodynamic motions, ξ and α .

The transfer function used was a function of ξ and α and their first derivatives.

$$[T] = [t_{11}, t_{12}] + \frac{a_T}{w_R} s[t_{11}, t_{12}] .$$
(128)

This is referred to as a *damping type transfer function* by Nissim and Abel (1978), where a_T is a control of the amount of damping introduced by the control surface (flap), and w_R is a reference frequency, normally taken as the no-control flutter frequency. This transfer function can be considered as a type of proportional-derivative (PD) control.

A proportional-integral-derivative (PID) control was also implemented by adding an integrator term to give a transfer function of the following form

$$[T] = [t_{11}, t_{12}] + \left(\frac{a_T}{w_R}\right)_1 s[t_{11}^*, t_{12}^*] + \left(\frac{a_T}{w_R}\right)_2 \frac{1}{s}[t_{11}^{**}, t_{12}^{**}] .$$
(129)

For better numerical stability in the eigenvalue solution, modified control transfer functions [T] were employed for both PD and PID control (Nissim and Abel 1978).

PD control

$$[T] = [t_{11}, t_{12}] + \frac{a_I}{w_R} s \frac{50000}{s + 50000} [t_{11}, t_{12}]$$
(130)

PID control

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$$[T] = [t_{11}, t_{12}] + \left(\frac{a_T}{w_R}\right)_1 s \frac{50000}{s + 50000} [t_{11}^*, t_{12}^*] + \left(\frac{a_T}{w_R}\right)_2 \frac{1}{s} [t_{11}^*, t_{12}^*]$$
(131)

The added term $\frac{50000}{s+50000}$ forces both the structure and the control matrix terms to be a part of the highest order coefficient in *s* of the polynomial. Without the addition of such a term, the highest order coefficient of the polynomial would be a function of the control matrix terms alone. This results in both very large and very small solutions for the eigenvalues and thus, places undue strain on the numerical accuracy.

Initially the *damping type transfer function* (PD control) was considered for anal ysis. The values suggested by Abel and Nissim (1978), who had used a theory based on the aerodynamic energy concept, were used as a starting point to determine realistic parameters in the transfer function. A simplex method was then used to determine if these were indeed the optimal values.

The aerodynamic energy concept considers the work done by the aerodynamic forces on the airfoil per cycle of oscillation, referred to as P. Hence, the transfer function [T]was determined to give a stable system, or a negative value of work P. For an aufoil of specific geometric configuration, Nissim and Abel determined the optimal parameters for a PD transfer function control. For an airfoil flap having a length equal to 20% of the chord ($e_3 = 0.6$, see figure 10), and having its displacement measured at 30% chord from the leading edge, (a = -0.1), they found the following values for the transfer function

$$[t] = [t_{11}, t_{12}] = [0.0, -1.86],$$

and

$$[l^{+}] = [l_{11}^{-} , l_{12}^{-}] = [l 0 , 3.20] .$$

In the present study, an airfoil flap having a length equal to 20% of the chord $(c_{\beta} = 0.6, \text{see figure 10})$ and having its elastic axis at 25% of the chord from the leading edge (corresponding value of a = -0.5) was considered. A transformation was completed (see Appendix K) to account for the difference in position at which the measurements were made, this gave

$$\begin{bmatrix} t \end{bmatrix} = \begin{bmatrix} t_{11} & t_{12} \end{bmatrix} = \begin{bmatrix} 0.0 & -1.86 \end{bmatrix}$$

and

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$$[t^*] = [t_{11}^*, t_{12}^*] = [4.0, 3.60]$$

The optimization of the transfer function parameters performed by the simplex sequential method considered the net affect the active control had on the flutter velocity and not on the overall system damping. It thus determined parameter values that resulted in an increase in the critical flutter velocity. In fact, the overall magnitude of the system damping seemed to decrease over the majority of the velocity range, but still gave a tayourable increase in the flutter speed. The optimized transfer function for case studies 1, 2 and 3, utilizing a PD control, was found to be

$$[T] = [0.0, -0.3] + \frac{0.0012}{0.548} s \frac{50000}{s + 50000} [4.0, 3.60]$$
(1.32)

For case studies 4 and 5, the optimized transfer function was calculated as

$$[T] = [0.0, -0.09] + \frac{0.0012}{0.548} s \frac{50000}{s + 50000} [4.0, 3.60]$$
(1.33)

Note, the value of w_6 could have been changed to reflect the no control frequency of the case study in question. However, any change in the value of w_R would have resulted in a proportional increase in the value of a_T . What was important was the relative ratio of t_{11}^2 to t_{12}^2 .

The calculated values for the transfer function agree quite well with the first derivative generalized coordinates given by the aerodynamic energy concept, but some variance was encountered with respect to the coefficient for the coordinate α . It should be noted that the plunging coordinate, ξ , gave very little useful control information for the system. The method given by Nissim and Abel was quite helpful, in that it gave an indication of the expected magnitude and sign of the required control parameters. It should also be noted that the optimized control parameters are not unique, other values could be used which may achieve the same increase in the flutter velocity.

In order for a system to be completely stable, all the eigenvalues calculated from the aeroelastic equations of motion must possess positive damping characteristics. This can be simplified by considering only the eigenvalues with positive frequency

The addition of active control to the airfoil had varying affects on each of the cases studied. In some cases, the roots registered only minor changes, while in other cases the changes were quite significant compared to the eigenvalue solution with no control. The results obtained are given as a graph of damping ratio versus non-dimensional velocity. U (see Figure 11 and Figure 12). Both the no-control and PD-control eigenvalues are plotted versus the non-dimensional velocity, \bar{U} . The no-control roots are defined as NO CONTROL ROOT1 and NO CONTROL ROOT2. The root defined as NO CONTROL ROOT1, shows increasing damping with velocity, until at a certain non-dimensional velocity, it begins to show signs of a decrease in damping. The point where the damping is zero is referred to as the margin of instability, and any increase in velocity will result in flutter. The root denoted as NO CONTROL ROOT2, shows a proportional increase in damping with an increase in velocity. The introduction of active control changes the nature of the two roots, ROOT1 and ROOT2. The PD-control roots are denoted by PD CONTROL ROOT1 and PD CONTROL ROOT2.

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Two scenarios were investigated. For example, active control was implemented for the parameters; $\mu = 50.0$, $r_{\alpha} = 0.5$, a = -0.5 and $w_h/w_{\alpha} = 0.2$ (Case Study 1). The results are given as a graph of damping ratio versus non-dimensional velocity, \bar{U} (see Figure 11). The control root, denoted as PD CONTROL ROOT1, showed similar behaviour to that of the no-control root, NO CONTROL ROOT1, except that the control root began to show a decay in damping at a larger velocity than that of the no-control root. This resulted in an increase in the critical flutter velocity, from $\bar{U} = 4.53$ to $\bar{U} = 4.8$, an effective increase of approximately 6%. However, the implementation of the control law had quite an affect on the other root, defined as PD CONTROL ROOT2. The result was that this root showed unstable behaviour, negative damping, at low velocity. Thus, such an active control device would have to be turned off initially, and then later turned on, once the aircraft has reached a stable velocity. It may be desired to implement the flutter suppression controller at about 80% of the no control flutter speed

A second example, with the parameters: $\mu = 50.0$, $r_a = 0.6$, a = -0.6 and w_h/w_s , 0.4 (Case Study 2), was also considered (see Figure 12). In this example, the root defined as PD CONTROL ROOT2 showed stable damping characteristics over the entrie velocity range. However, the root defined as PD CONTROL ROOT1, showed negative damping at low values of velocity. This root goes through three distinct stages of development. The first stage is defined when the root has non-positive damping. The second stage is defined by the root possessing positive damping characteristics, but still has less damping in comparison to its no-control root. NO CONTROL ROOT1. The final stage begins when the PD-control root shows improved damping in comparison to the no-control root. This allowed for an increase in the flutter velocity from U = 5.1 to U = 5.6. This resulted in a net increase in flutter velocity of approximately 10% (see Figure 12). Additional comparisons of no-control and PD-control for various sets of parameters are included in Fable 7.

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Case					Flutter Velocity Û			
Stud	ין ע	r_{α}	(1	$w_h \neq w_0$	No ('ontrol	PD Control	% Increase	
I	50.0	0.5	-0.5	0.2	4.53	4.8	6 0	
2	50.0	0.6	-0.6	0.4	5.10	5.6	10.0	
.}	100.0	0.5	-0.5	0.2	6.24	6.6	5.5	
1	75-0	0.1	-0.1	0.3	3.68	3.9	5 5	
5	100/0	0.1	-0.4	0.3	4.16	4.4	5.5	

Table 7 Incompressible Regime : Comparison of critical flutter velocity, \bar{U} , using no control and PD control solved by the *Laplace* method, with the constants : $r_{\beta} = 0.5$, $x_{\beta} = 0.25$, $x_{\alpha} = 0.25$, $c_{\beta} = 0.6$ and b = 1.0.

Although the increase in flutter speed in these case studies can be considered only as fair, it has been shown that control analysis can be implemented quite easily by working in the Laplace domain. In past studies, increases in flutter speeds of 20 to 25% have been achieved

The use of a PID control law resulted in no additional increase (beyond that achieved via PD control) in the critical flutter velocity, \bar{U} . Maybe, this was to be expected. The introduction of an integrator term is usually done in the hope of cancelling the steady-state error, in comparison to a known desired control velocity. However, in this context, the object was not to control the velocity to some known required velocity, but to find the greatest possible maximum critical velocity.

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11 Conclusions

A flutter analysis was considered for a rigid autoil, flexibly mounted in the bending and torsional directions. A simple and effective way of determining flutter velocities was developed by solving the equations in the Laplace domain. In this procedure the iteration process encountered in the traditional U = q and p - k methods is eliminated, by using rational functions to approximate the unsteady aerodynamic forces as a function of reduced frequency.

The validity of the Laplace solution depends greatly on how well the unsteady acrodynamics can be approximated. It is critical that the approximation be very close to the actual data. The unsteady acrodynamics for both incompressible and transourcillow were approximated by rational functions of the Nasa-Langley form with great success

$$[N] = [A_0] + \frac{b}{U} [A_1] s + (\frac{b}{U})^2 [A_2] s^2 +$$
$$[A_3] \frac{s}{s + \frac{l}{b} b_1} + [A_1] \frac{s}{s} \frac{b_2}{b_2} + [A_5] \frac{s}{s + \frac{l}{b} b_3} + [A_0] \frac{s}{s + \frac{l}{b} b_4}$$

The optimization method employed a least squares method in conjunction with a simplex method. The use of the simplex method to optimize the lag terms showed marked improvement in the approximation in comparison with previous work, where the lag terms were chosen randomly. For incompressible flow, it was necessary to approximate Theodorsen's function, C(k) and C(k)ik; the approximation error was found to be insignificant, even for a large reduced frequency range of 0.9 < k < 10.0. This allowed for a robust evaluation of any aeroelastic problem in the incompressible flow. The Laplace method was found to give values of flutter velocity and flutter frequency identical to those given by the p - k and the U - g methods. For values of subcritical velocity, the Laplace method showed good comparison with frequency and damping obtained via the p - k method

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For transmic flow, it was necessary to approximate the four unsteady aerodynamic coefficients C_{l_h} , C_{l_h} , C_{m_h} and C_{m_h} . A much decreased range of reduced frequency was considered, $0.0 \le k \le 0.5$. Also, constraints were imposed on the aerodynamic derivatives whose magnitudes were zero at k = 0.0, namely C_{l_h} and C_{m_h} . The results were not as good as those obtained in incompressible flow, but were still found to be excellent overall. The reduced frequency range had little effect on the calculation of the flutter velocities, and fail representations of frequency and damping ratio over most of the range of nondimensional velocity were obtained. However, some discrepancies were apparent at the lower values of velocity for one of the complex conjugate roots. Here, the frequencies obtained by the eigenvalue solution fell outside of the applicable reduced frequency range, and thus, resulted in some difference in the damping and frequency response to that obtained with the p - k and U - g methods.

Solving the equations in the Laplace domain allows for the efficient implementation of active control to alleviate the flutter instability. Such control was considered with some success for incompressible flow. Although, the increase in the flutter velocity can only be considered as fair (6 to 10%), it was shown that the Laplace method was a viable way of incorporating active control into the aeroelastic equations.

Future work could consider the implementation of other control laws, in the hope of obtaining a more significant increase in flutter velocity. In addition, the control theory could be extended to account for transonic flow. This would require the determination of the unsteady aerodynamics incorporating a flap for transonic flow.

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Figure 1 Schematic representation of airfoil with freedom to move in torsional and bending directions (2DOF). (Lee 1984)

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Figure 2a Derivative of Lift Coefficient, C_l , with respect to Plunge, h, versus Reduced Frequency, k, at M = 0.85 for a NACA 64,4006 Aufoil



Figure 2b Derivative of Lift Coefficient, C_l , with respect to Pitch, α , versus Reduced Frequency, k, at M = 0.85 for a NACA 64A006 Airford

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Figure 2c Derivative of Moment Coefficient, C_m , with respect to Plunge, h, versus Reduced Frequency, k, at M = 0.85 for a NACA 64A006 Airfoil



Figure 2d Derivative of Moment Coefficient, C_m , with respect to Pitch, α , versus Reduced Frequency, k, at M = 0.85 for a NACA 64A006 Airfoil



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Figure 3 Eigenvalue Solution : Damping and Frequency versus Dynamic Pressure (Abel 1979)



Figure 1a-1 Comparison between the Rational Function Approximation of C(k)and the True C(k), Real Part, as a Function of Reduced



Figure 1a-2 Comparison between the Rational Function Approximation of C(k)and the True C(k), Real Part, as a Function of Reduced Frequency, 0.0 < k < 2.0 - a magnification of figure 4a-1



Figure 4a-3 Comparison between the Rational Function Approximation of C(k)and the True C(k), Imaginary Part, as a Function of Reduced



Figure 4a-4 Comparison between the Rational Function Approximation of C(k)and the True C(k), Imaginary Part, as a Function of Reduced Frequency, 0.0 < k < 2.0 - a magnification of figure 1a-3



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Figure 4b-2 Comparison between the Rational Function Approximation of C(k)ikand the True C(k)ik, Real Part, as a Function of Reduced Frequency, 0.0 < k < 2.0 - a magnification of figure 4b-1



Figure 1b 3 Comparison between the Rational Function Approximation of C(k)ikand the True C(k)ik. Imaginary Part, as a Function of Reduced



Figure 4b-4 Comparison between the Rational Function Approximation of C(k)ikand the True C(k)ik. Imaginary Part, as a Function of Reduced Frequency, 0.0 < k < 2.0 - a magnification of figure 4b-3

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Figure 5a-1 Comparison between the Rational Function Approximation of C_{l_h} and the Frue C_{l_h} as a Function of Reduced Frequency, $0.025 \le k \le 1.0$,

at M = 0.85 for a NACA 64A006 Airfoil.



Figure 5b-1 Comparison between the Rational Function Approximation of $C_{l_{\alpha}}$ and the True $C_{l_{\alpha}}$ as a Function of Reduced Frequency, $0.025 \le k \le 1.0$, at M = 0.85 for a NACA 64A006 Airfoil.



Figure 5c-1. Comparison between the Rational Function Approximation of C_{m_h} and the True C_{m_h} as a Function of Reduced Frequency, $0.025 \le k_{\perp} = 1.0$,

at M = 0.85 for a NACA 64,1006 Au foil.



Figure 5d-1 Comparison between the Rational Function Approximation of C_{m_0} and the True C_{m_0} as a Function of Reduced Frequency, $0.025 \le k \le 1.0$, at M = 0.85 for a NACA 64A006 Airfoil.



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Figure 5a-2 Comparison between the Rational Function Approximation of C_{lh} and the True C_{lh} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$.

at M = 0.85 for a NACA 64A006 Airfoil.



Figure 5b-2 Comparison between the Rational Function Approximation of $C_{l_{\alpha}}$ and the True $C_{l_{\alpha}}$ as a Function of Reduced Frequency, $0.0 \le k \le 0.5$, at M = 0.85 for a NACA 64 4006 Airfoil.



Figure 5c-2 Comparison between the Rational Function Approximation of C_{m_h} and the True C_{m_h} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$.

at M = 0.85 for a NACA 64 4006 Airford



Figure 5d-2 Comparison between the Rational Function Approximation of $C_{m_{+}}$ and the True $C_{m_{0}}$ as a Function of Reduced Frequency, $0.0 \le k \le 0.5$ at M = 0.85 for a NACA 64,4006 Airfoil

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Figure 6.4 Comparison between the Rational Function Approximation of C_{l_h} and

the true C_{l_h} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$.

at M = 0.80 for a NACA 64A006 Airford.







Figure 6c Comparison between the Rational Function Approximation of C_{m_h} and the True C_{m_h} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$,

at M = 0.80 for a NACA 64A006 Airfoil.



Figure 6d Comparison between the Rational Function Approximation of $C_{m_{\alpha}}$ and the True $C_{m_{\alpha}}$ as a Function of Reduced Frequency, $0.0 \le k \le 0.5$, at M = 0.80 for a NACA 64A006 Airfoil.



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Figure 7a Comparison between the Rational Function Approximation of C_{l_h} and the True C_{l_h} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$.



Figure 7b Comparison between the Rational Function Approximation of C_{l_0} and the True C_{l_0} as a Function of Reduced Frequency. $0.0 \le k \le 0.5$. at M = 0.875 for a NACA 64A006 Airfoil.



Figure 7c Comparison between the Rational Function Approximation of C_{m_h} and the True C_{m_h} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$,

at M = 0.875 for a NACA 64.4006 Airfoil.



Figure 7d Comparison between the Rational Function Approximation of C_{m_n} and the True C_{m_n} as a Function of Reduced Frequency, $0.0 \le k \le 0.5$. at M = 0.875 for a NACA 64A006 Airfoil.



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Figure 8a-1 Comparison of U - g and p - k with Laplace generated graphs

of Damping Ratio, Root 1, versus Non-Dimensional Velocity, \bar{U} .

Case Study $\#1: \mu = 50.0, r_{\alpha} = 0.5, a = -0.5, w_h/w_{\alpha} = 0.2,$



Figure Sa-2 Comparison of U - g and p - k with Laplace generated graphs of Damping Ratio, Root 2, versus Non-Dimensional Velocity, \overline{U} . Case Study #1 : $\mu = 50.0$, $r_{\alpha} = 0.5$, a = -0.5, $w_h/w_{\alpha} = 0.2$,



Figure 8a-3 Comparison of U - g and p - k with Laplace generated graphs

of Frequency Ratio, w/w_{α} , versus Non-Dimensional Velocity, U.



Figure 8b-1 Comparison of U - g and p - k with Laplace generated graphs

of Damping Ratio, Root 1, versus Non-Dimensional Velocity, \hat{U} . Case Study #2 : $\mu = 50.0$, $r_{\alpha} = 0.6$, a = -0.6, $w_h/w_{\alpha} = 0.4$, Flutter Velocity, $\bar{U} = 5.11$



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Figure 8b-2 Comparison of U - g and p - k with Laplace generated graphs

of Damping Ratio, Root 2, versus Non-Dimensional Velocity, \overline{U} .



Figure 8b-3 Comparison of U - g and p - k with Laplace generated graphs of Frequency Ratio, w/w_{α} , versus Non-Dimensional Velocity, \bar{U} . Case Study #2 : $\mu = 50.0$, $r_{\alpha} = 0.6$, a = -0.6, $w_h/w_{\alpha} = 0.4$,





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Figure 10 Schematic Representation of Airfoil with Freedom to Move in Torsional and Bending Directions with a Flap (3DOF). (Karpel 1981)



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Figure 11 Implementation of Active Control for Incompressible Regime.

Case Study #1 : $\mu = 50.0$, $r_{\alpha} = 0.5$, a = -0.5, $w_h/w_{\alpha} = 0.2$

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Figure 12 Implementation of Active Control for Incompressible Regime.

Case Study $#2: \mu = 50.0, r_a = 0.6, a = -0.6, w_b/w_a = 0.4$

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Appendix A : Objective Error Function

Many forms of an error function can be used

Q(ik) = F + iG rational function approximation

Q(ik) = F + iG tabular data (discrete values of reduced frequency)

Error #1 (one used in study)

$$err(ik) = \sum \frac{|\hat{Q}(ik) - Q(ik)|^2}{max \{1, |Q(ik)|^2\}}$$

Enor #2

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$$err(ik) = \sum \frac{|\bar{Q}(ik) - Q(ik)|^2}{\{|Q(ik)|^2\}}$$

Error #3

$$err(ik) = \sum \left\{ \frac{(\bar{F} - F)^2}{F^2} + \frac{(\bar{G} - G)^2}{G^2} \right\}$$

Error #4

$$err(\imath k) = \sum \left\{ abs\left[\frac{(\bar{F}-F)}{F}\right] + abs\left[\frac{(\bar{G}-G)}{G}\right] \right\}$$

where $|Q(ik)|^2 = F^2 + G^2$

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other variations of the error objective function could include

Erroi #5

$$err(ik) = \sum \frac{|Q(ik) - Q(ik)|^2}{max \{1, |Q(ik)|^2\} k}$$

Enor #6

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$$err(ik) = \sum \frac{|Q(ik) - Q(ik)|^2}{max\{1, |Q(ik)|^2\}k^2}$$

where $|Q(ik)|^2 = F^2 + G^2$.

Analysis has showed that the error function, given by Error #1, resulted in the best comparison of \bar{U} , $w \neq w_{\alpha}$ and k to Lee (1984). For example, the table below considers a reduced frequency range of $0.0 \le k \le 0.5$, with forced constraints for all aerodynamic coefficients at k = 0.0.

Error Type	Reduced Frequency Range, k	ľ	$w \mid w_{\alpha}$	k.,
#1	$0.0 \le k \le 0.5$	3.45	0.309	0.090
#2	$0.0 \le k \le 0.5$	3.36	0.303	0 090
#3	$0 \ 0 \le k \le 0.5$	3.42	0.305	0.089
#4	$0 \ 0 \le k \le 0.5$	3.38	0 302	0.089

Note: Depending on the original nature of the tabular data, one of the other presented error functions may have resulted in better results. It was decided to implement the same error function as used by Adams and Tiffany (1988), namely Error Type #1

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Appendix B : Linear Coefficient Optimization - Least Squares Methodology

The linear coefficients, A_0 , A_1 , ..., A_6 , are optimized by utilizing a least square methodology. In order to obtain a minimum error function, the derivative of the error function is taken with respect to each linear coefficient, and written as

$$\frac{\partial crr(\imath k)}{\partial (A_j)} = 0$$

The error function is defined as

$$err(ik) = \sum \frac{|\tilde{Q}(ik) - Q(ik)|^2}{M(ik)}$$

where

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$$M(ik) = max \{1, |Q(ik)|^2\}$$
.

The quantities $\hat{Q}(ik)$ and Q(ik) represent the rational function approximation and the exact value, respectively, at a given reduced frequency. The term M(ik) is used to normalize the aerodynamic tabular data, such that certain points do not receive larger than normal weighting. Note, the error function is defined as the total normalized sum for all the given values of the unsteady aerodynamic tabular data.

The following analysis considers an approximating function of only three lag terms, hence linear coefficients $A_0, A_1, ..., A_5$. The extension to additional lag terms is straight forward.

Making the following substitutions for $\bar{Q}(ik)$ and Q(ik),

$$Q(ik) = A_0 + A_1ik + A_2(ik)^2 + A_3\frac{ik}{ik+b_1} + A_4\frac{ik}{ik+b_2} + A_5\frac{ik}{ik+b_3}$$

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and

$$Q(ik) = F + iG .$$

into the above equations gives

$$\frac{\partial \left[\sum_{i=1}^{i} \frac{\left|\left(1_{0}+1_{1}ik+1_{2}(ik)^{2}+4_{1}\frac{ik}{ik+b_{1}}+1_{1}\frac{ik}{ik+b_{2}}+1_{5}\frac{ik}{ik+b_{1}}\right)-(F+iG)\right|^{2}}{M(ik)}\right]}{\partial(A_{J})}{=0}$$

Using $i^2 = -1$, and cleaning up the denominator of terms of the form ik + b, where j = 1 to 3, gives

$$\frac{\imath k}{\imath k + b_j} = \frac{\imath k}{\imath k + b_j} \left(\frac{\imath k - b_j}{\imath k - b_j} \right) = \frac{k^2 + \imath k b_j}{k^2 + b_j^2}$$

The numerator term, given by

$$\left[\sum \frac{\left|\left(A_{0}+A_{1}ik+A_{2}(ik)^{2}+A_{3}\frac{ik}{ik+b_{1}}+A_{4}\frac{ik}{ik+b_{2}}+A_{5}\frac{ik}{ik+b_{2}}\right)-(F+iG)\right|^{2}}{M(ik)}\right]$$

can then be rewritten to give

$$\left[\sum \frac{\left|\left(A_0 + A_1ik - A_2(k)^2 + A_3\frac{k^2 + ikb_1}{k^2 + b_1^2} + A_4\frac{k^2 + ikb_2}{k^2 + b_2^2} + A_5\frac{k^2 + ikb_3}{k^2 + b_1^2}\right) - (F + i(i))^2}{M(ik)}\right]$$

This can be further simplified to

$$\sum \frac{|(a+ib)|^2}{M(ik)} = \sum \frac{(a^2+b^2)}{M(ik)} \; .$$

where

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$$a = A_0 - A_2 k^2 + A_3 \frac{k^2}{k^2 + b_1^2} + A_4 \frac{k^2}{k^2 + b_2^2} + A_5 \frac{k^2}{k^2 + b_3^2} - F .$$

and

$$b = A_1k + A_3\frac{kb_1}{k^2 + b_1^2} + A_4\frac{kb_2}{k^2 + b_2^2} + A_5\frac{kb_3}{k^2 + b_3^2} - G$$

This can be further simplified to give

$$a^{2} = A_{0}^{2} + A_{2}^{2}k^{4} + A_{3}^{2}B_{1}^{2} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}A_{5}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{2}k^{2} - 2A_{0}F + 2A_{0}A_{3}B_{1} + A_{4}^{2}B_{2}^{2} + A_{5}^{2}B_{3}^{2} - 2A_{0}A_{3}B_{1} + A_{6}^{2}B_{2}^{2} + A_{6}^{2}B_{3}^{2} - 2A_{0}A_{2}B_{1} + A_{6}^{2}B_{2}^{2} + A_{6}^{2}B_{3}^{2} - 2A_{0}A_{2}B_{1} + A_{6}^{2}B_{2}^{2} + A_{6}^{2}B$$

$$2A_0A_4B_2 + 2A_0A_5B_3 + 2A_2k^2F - 2FA_3B_1 - 2FA_4B_2 - 2FA_5B_3 + 2A_3B_1A_4B_2 + 2A_3B_1A_5B_3 + 2A_4B_2A_5B_3 + F^2 ,$$

and

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$$b^{2} = A_{1}^{2}k^{2} + A_{3}^{2}\frac{B_{1}^{2}b_{1}^{2}}{k^{2}} + A_{4}^{2}\frac{B_{2}^{2}b_{2}^{2}}{k^{2}} + A_{5}^{2}\frac{B_{3}^{2}b_{3}^{2}}{k^{2}} - 2A_{1}kG + 2A_{1}A_{3}B_{1}b_{1} + 2A_{1}A_{4}B_{2}b_{2} + 2A_{1}A_{5}\frac{B_{1}b_{1}B_{3}b_{3}}{k} + 2A_{2}A_{4}\frac{B_{1}b_{1}B_{2}b_{2}}{k^{2}} + 2A_{3}A_{5}\frac{B_{1}b_{1}B_{3}b_{3}}{k^{2}} + 2A_{4}A_{5}\frac{B_{2}b_{2}B_{3}b_{3}}{k^{2}} ,$$

where

$$B_j = \frac{k^2}{k^2 + b_j^2}$$
, $j = 1 \text{ to } 3$.

Hence, taking the derivative of the error function with respect to the linear coefficient $I_0\ {\rm gives}$

$$\left[\sum \frac{2A_0 - 2A_2k^2 - 2F + 2A_3B_1 + 2A_4B_2 + 2A_5B_3}{M(\imath k)}\right] = 0 ,$$

or

$$\sum \frac{A_0 - A_2 k^2 + A_3 B_1 + A_4 B_2 + A_5 B_3}{M(ik)} = \sum \frac{F}{M(ik)} \,.$$

This can also be written as

$$[A'] \{x\} = [B'] ,$$

where

$$[A'] = \begin{bmatrix} 1 & 0 & -k^2 & B_1 & B_2 & B_3 \end{bmatrix},$$
$$[B'] = \begin{bmatrix} F \end{bmatrix}$$

and

\$

$$\{x\} = \{ A_0, A_1, A_2, \cdots, A_5 \} .$$

A-5

Similarly, taking the derivative with respect to the coefficients A_i , j = 1 to 5 yields a set of linear algebraic equations, that can be written as

$$[A] \{x\} = [B]$$

where

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$$[A] = \sum \frac{1}{M(ik)} \begin{bmatrix} 1 & 0 & -k^2 & B_1 & \cdots & B_3 \\ 0 & k^2 & 0 & B_1b_1 & \cdots & B_3b_3 \\ -k^2 & 0 & k^3 & -k^2B_1 & \cdots & -k^2B_3 \\ B_1 & B_1b_1 & -k^2B_1 & B_1B_1(1 + \frac{b_1b_1}{k^2}) & \cdots & B_1B_3(1 + \frac{b_1b_1}{k^2}) \\ B_2 & B_2b_2 & -k^2B_2 & B_2B_1(1 + \frac{b_2b_1}{k^2}) & \cdots & B_2B_3(1 + \frac{b_1b_1}{k^2}) \\ B_3 & B_3b_3 & -k^2B_3 & B_3B_1(1 + \frac{b_3b_1}{k^2}) & \cdots & B_3B_3(1 + \frac{b_1b_1}{k^2}) \end{bmatrix}$$

$$[B] = \sum \frac{1}{M(ik)} \left[F \cdot kG - k^2 F \cdot FB_1 + \frac{GB_1b_1}{k} \cdot \cdots \cdot FB_3 + \frac{GB_3b_3}{k} \right]^T$$

and

$$\{x\} = \{A_0, A_1, A_2, \cdots, A_5\}$$
.

Appendix C : Nonlinear Coefficient Optimization - Simplex Method

The best way to explain the simplex method is through an example. Consider a function f consisting of two independent variables (n = 2), say X_1 and X_2 . This problem requires the function evaluation $f(X_1, X_2)$ at three (n + 1) different sets of values to produce three vertices. Let these initial simplex vertices be denoted as A, B and C. which would form a triangular simplex (see figure). In general, the simplex will not be regular in shape. The simplex method then considers the function evaluations at these three points, and determines the largest, say for argument, it is f(A), and that f(C) < f(B) < f(A). It now takes a series of steps, to move the point of the simplex, which has the largest value, A, through the opposite face of the simplex, to produce a new point that will replace 4. The initial step taken, is to reflect the point A through the geometric centroid of the two vertices, B and C (the midpoint), resulting in a reflection point E. In general, the centroid is defined as the centroid of all the points $(X_1)_i$, $(X_2)_i$. \dots ; not including the point of the highest function value. Note, that the lengths AD and DE are equal. Depending on the magnitude of the function value at E, the simplex will do an extension to F, contraction to H, a lesser contraction to G, or maintain the initial reflection to form *BCE*. The process continues until convergence occurs.

Criteria:

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1) if f(E) < f(C) locate F (extension point) along AD of length DE

2) if f(E) > f(A) locate G (interior contraction point) midpoint of AD

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3) if f(B) < f(E) < f(A) locate H (exterior contraction point) midpoint of DT
4) if f(C) < f(E) < f(B) keep E (reflection point) new simplex is BCE
Simplex Method : f(X₁, X₂)



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Nasa-Langlev approximate in Laplace form

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$$[N] = [A_0] + \frac{b}{U} [A_1] s + (\frac{b}{U})^2 [A_2] s^2 +$$
$$[A_3] \frac{s}{s + \frac{U}{b}b_1} + [A_4] \frac{s}{s + \frac{U}{b}b_2} + [A_5] \frac{s}{s + \frac{U}{b}b_3} + [A_6] \frac{s}{s + \frac{U}{b}b_4}$$

Approximation of Theodorsen Function C(k)

$$[N] = [0.99828] + \frac{b}{U} [-0.00004] s + (\frac{b}{U})^2 [0.00000] s^2 + [-0.04120] \frac{s}{s + \frac{U}{b}(0.01544)} + [-0.16289] \frac{s}{s + \frac{U}{b}(0.08429)} + [-0.22909] \frac{s}{s + \frac{U}{b}(0.25595)} + [-0.06459] \frac{s}{s + \frac{U}{b}(0.76511)} .$$

Approximation of C(k) * ik

$$[.V] = [-0.00001] + \frac{b}{U} [0.50037] s + (\frac{b}{U})^2 [-0.00004] s^2 + [0.00061] \frac{s}{s + \frac{U}{b} (0.01544)} + [0.01380] \frac{s}{s + \frac{U}{b} (0.08429)} + [0.05848] \frac{s}{s + \frac{U}{b} (0.25595)} + [0.04963] \frac{s}{s + \frac{U}{b} (0.76511)} .$$

note : lag terms are equivalent

SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION INCOMPRESSIBLE REGIME GLOBAL OPTIMIZATION RANGE 0.00 < k < 10.00

N(1, 1) = C(k)

k	EXACT APPROXIMAT		TE	
	REAL	IMAG	REAL	IMAG
10.00000.	50060	- 01240	50057	
6.000000	50170	- 02060	.50057,	01257,
4.000000	50370	- 03050	· 50162,	02047,
3.000000.	.50630	-04000	50572,	03025,
2.000000	.51290	- 05770	· 50019,	03976,
1.500000	52100	- 07360	· 512/1,	05/64,
1.200000	.52100,	- 08770	.52065,	0/3/3,
1.000000	.53940	- 10030	• 52 5 54 ,	08/98,
.990000.	.54000	-10100	· 53959,	10051,
.880000	54740	-10950	· J4019,	10122,
.800000	55410	- 11650	• 54707,	10963,
.770000	55700	-11030,	• 55445,	11055,
.660000	56990	- 13080	· 55733,	-11934,
.600000	57880	- 13780	57012,	13063,
560000	58570	- 14290	· 5/695,	13/5/,
550000	58760	-14200,	· 505/9,	14253,
500000,	59790	-15070	· 28/64,	14381,
440000,	61300,	- 15070,	· 59/8/,	15049,
400000,	62500,	- 16500	·012/0,	15904,
340000	.02500,	-17390	• 02477,	- 10499,
330000	.04090,	- 17520,	• 04003,	1/39/,
325000	.05120,	- 17520,	.65106,	1/54 <i>3</i> ,
320000,	.05550,	- 17660	.00328,	1/614,
315000	.05580,	- 17720	.60002,	1/686,
.310000,	.05010,	-17700,	• 65/82,	1//56,
.310000,	.00040,	- 17960	.00017,	1/825,
.305000,	.00270,	17020	.66257,	1/894,
.300000,	.00500,	-1960,	. 66502,	1/961,
.240000,	.09090,	-10020,	.0992/,	18633,
.220000,	.71250,	-18960	./129/,	18//0,
.200000,	.72700,	-1000,	. 72802,	18841,
120000,	.70200,	-18010	. / 6288,	18/03,
.120000,	.00030,	18010,	.80384,	1/9//,
.110000,	.81880,	17660,	.81822,	1/652,
.100000,	.83200,	1/230,	.83138,	1/244,
.060000,	.00040,	10040,	.86027,	16092,
.000000,	.89200,	14260,	.89253,	14285,
.050000,	.90900,	13050,	.90957,	13041,
.040000,	.92670,	11600,	.92684,	11525,
.025000,	.95450,	08720,	.95305,	08710,
.010000,	.98240,	04820,	.98349,	04764,

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N(1, 2) = C(k) * ik

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k	EXACT		APPROXIMATE		
	REAL	IMAG	REAL	IMAG	
10 00000	12400	5 00600	12582	5 00914	
6 000000,	12360	3 01020	12292	3 01116	
4 000000,	12200	2.01480	.12110	2.01468	
3 000000,	.12000,	1.51890.	.11938.	1.51835.	
2 000000	11540	1.02580.	.11536.	1.02526.	
1 500000	11040	. 78150.	.11065.	.78113.	
1.200000	.10524.	.63600.	.10561.	.63585.	
1.000000	.10030.	.53940.	.10053.	.53954.	
.990000	.09999.	.53460,	.10023,	.53474	
.880000.	.09636.	.48171,	.09649,	.48191,	
.800000.	.09320.	. 44328	.09325,	.44352	
.770000.	.09186,	. 42889,	.09190,	.42912,	
.660000,	.08633,	. 37613,	.08621,	.37626,	
.600000.	.08268,	.34728,	.08253,	.34735,	
.560000,	.07997,	. 32799,	.07980,	.32802,	
.550000,	.07926,	. 32318,	.07908,	.32318,	
.500000	.07535,	.29895,	.07522,	.29892,	
.440000,	.07005,	.26972,	.06995,	.26961,	
.400000,	.06600,	.25000,	.06597,	.24989,	
.340000,	.05909,	.21995,	.05912,	.21990,	
.330000,	.05782,	.21490,	.05786,	.21483,	
.325000,	.05717,	.21239,	.05722,	.21229,	
.320000,	.05651,	.20986,	.05656,	.20974,	
.315000,	.05585,	.20730,	.05590,	.20719,	
.310000,	.05515,	.20472,	.05523,	.20463,	
.305000,	.05447,	.20212,	.05455,	.20206,	
.300000,	.05379,	.19950,	.05385,	.19948,	
.240000,	.04469,	.16774,	.04469,	.16780,	
.220000,	.04129,	.15675,	.04127,	.15683,	
.200000,	.03772,	.14552,	.03766,	.14558,	
.160000,	.03002,	.12205,	.02991,	.12204,	
.120000,	.02161,	.09676,	.02156,	.09669,	
.110000,	.01943,	.09007,	.01941,	.09000,	
.100000,	.01723,	.08320,	.01723,	.08313,	
.080000,	.01283,	.06883,	.01286,	.06882,	
.060000,	.00856,	.05352,	.00855,	.05356,	
.050000,	.00653,	.04545,	.00650,	.04548,	
.040000,	.00464,	.03707,	.00459,	.03708,	
.025000,	.00218,	.02386,	.00215,	.02382,	
.010000,	.00048,	.00982,	.00046,	.00983,	

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Appendix E : Rational Function Optimization - Transonic Regime

Comparison between the rational function approximation of the aerodynamic derivatives and the true aerodynamics as a function of reduced frequency, -0.025 < k < 1.000, at M = 0.85.

 $C_{l_h}, C_{l_o}, C_{m_h}$ and C_{m_o}

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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION

TRANSONIC REGIME

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 C_{l_h}



k	EXACT		APPROXIMATE		k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
.025000,	.23000,	.56000,	.21658,	.56611,	.025000,	.01000,	03000,	01235,	01705,
.050000,	.40000,	.87000,	.41419,	.84979,	.050000,	.02000,	06000,	.03307.	05976
.075000,	.58000,	1.07000,	.59756,	1.06853,	.075000,	.03000.	10000.	.03738.	12445
.100000,	.74000,	1.22000,	.75517,	1.24000,	.100000,	.03000,	14000,	.02280,	17765,
.150000,	1.00000,	1.49000,	.98456,	1.50109,	.150000,	.01000,	23000,	01110,	24884,
.200000,	1.17000,	1.69000,	1.11518,	1.72130,	.200000,	01000,	33000,	02820,	29719
.250000,	1.23000,	1.92000,	1.17742,	1.94017,	.250000,	04000,	40000,	02752,	33989.
.300000,	1.22000,	2.12000,	1.19509,	2.17025,	.300000,	06000,	41000,	01311,	38366,
.350000,	1.19000,	2.48000,	1.18273,	2.41433,	.350000,	04000,	43000,	.01179,	43100,
.400000,	1.13000,	2.63000,	1.14896,	2.67229,	.400000,	.00000,	45000,	.04502,	48292,
.450000,	1.06000,	2.93000,	1.09904,	2.94340,	.450000,	.07000,	49000,	.08512,	53991,
.500000,	.99000,	3.25000,	1.03633,	3.22697,	.500000,	.14000,	55000,	.13108,	60228,
.550000,	.92000,	3.54000,	.96317,	3.52249,	.550000,	.21000,	64000,	. 18211,	67029,
.600000,	.88000,	3.89000,	.88125,	3.82966,	.600000,	.27000,	76000,	.23757,	74417,
.650000,	.79000,	4.20000,	.79193,	4.14829,	.650000,	.34000,	86000,	.29690,	82415,
.700000,	.68000,	4.52000,	.69631,	4.47832,	.700000,	.39000,	97000,	.35958,	91043,
.750000,	.57000,	4.88000,	.59538,	4.81976,	.750000,	.46000,	-1.06000,	.42512,	-1.00321,
.800000,	.47000,	5.22000,	.49004,	5.17262,	.800000,	.53000,	-1.13000,	.49304,	-1.10268,
. 850000,	.39000,	5.53000,	.38110,	5.53699,	.850000,	.57000,	-1.19000,	.56289,	-1.20898,
.900000	.29000,	5.87000,	.26936,	5.91290,	.900000,	.61000,	-1.28000,	.63420,	-1.32226,
.950000,	.20000,	6.25000,	.15556,	6.30041,	.950000,	.65000,	-1.40000,	.70653,	-1.44262,
1.000000,	.10000,	6.58000,	.04045,	6.69956,	1.000000,	.68000,	-1.53000,	.77943,	-1.57016,

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1-	EVI CO		A DDDAYTMATF		k	EXACT		APPROXIMATE	
ĸ	DEAT.	TMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
	READ	21110							
	10 20000	-4 01000	10 80926	-4.16547	.025000,	43700,	04100,	43992,	01892,
.025000,	10.79000,	-4.01000,	0 22629	-4.10547,	.050000,	45100,	12100,	45702,	16849,
.050000,	9.10000,	-2.21000,	9.23520,	-3 92265	.075000.	48200,	26200,	50534,	27825,
.075000,	8.00000,	-3.82000,	2 20208	-3.55967	. 100000	52300.	34800.	56402	35625,
.100000,	7.21000,	-3.49000,	7.20300,	-3,55667,	150000	68100.	51100.	66848.	45194
.150000,	6.09000,	-2.96000,	6.03060,	-2.93090,	200000	- 81700	58000.	74100.	51230.
.200000,	5.41000,	-2.29000,	5.34470,	-2.29399,	.200000,	- 90800	- 59800	- 78859	- 56393.
.250000,	4.93000,	-1.79000,	4.93936,	-1.71875,	.250000,	- 00600,	- 54900	- 92047	- 61610
.300000,	4.62000,	-1.28000,	4.67882,	-1.21716,	.300000,	90000,	54000,	- 04085	- 67166
.350000,	4.43000,	81000,	4.52562,	78120,	.350000,	86900,	5/400,	84285,	0/190,
400000	4.36000,	41000,	4.43443,	39950,	.400000,	82300,	64000,	8594/,	/3091,
450000	4.35000.	03000,	4.38410,	06209,	.450000,	78600,	74800,	87262,	79411,
500000	4.40000.	. 27000.	4.36177	.23873,	.500000,	77800,	90000,	88376,	86094,
550000,	A A3000.	. 55000.	4.35925.	.50872.	.550000,	81800,	99200,	89383,	93120,
	4 50000	74000	4.37104	.75210.	.600000.	85200,	-1.08000,	90351,	-1.00472,
.600000,	4.50000,	1 00000	4. 39332	.97197.	.650000.	87800	-1.14300,	91328,	-1.08137,
. 650000,	4.30000,	1 19000	A 40306	1 17064	700000	89200.	-1.23000	92354,	-1.1610
.700000,	4.48000,	1.10000,	A 45067	1 24093	760000	- 93000.	-1.30700.	93459,	-1.24377,
.750000,	4.48000,	1.31000,	4.4300/,	1,34303,	. / 50000,	- 97800	-1.38100.	94671.	-1.32940.
.800000,	4.48000,	1.47000,	4.49/81,	1,51085,	.800000,	-1 00800	-1.43900.	- 96013.	-1.41796.
.850000,	4.49000,	1.63000,	4.53923,	1.03470,	.850000,	-1.02100	-1 49500	- 97508.	-1.50942.
.900000,	4.56000,	1.80000,	4.58165,	1./8215,	.900000,	-1.03100,	-1.49300,	- 99176	-1.60377
.950000,	4.60000,	1.91000,	4.62398,	1,89385,	.950000,	-1.04300,	-1.53400,	-1 01035	-1.70098
1.000000,	4.62000,	2.02000,	4.66521,	1.99031,	1.000000,	-1.02800'	-1.204001	-1.010331	2

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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION

TRANSONIC REGIME

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APPROXIMANT FUNCTION N (1, 1) = C_{l_h} AO A1 A2 A3 A4 A5 A6 -7907.23428 13.851605 -1.182972 7907.335646 .299359 1.150099 -22.698018 lag1 lag2 lag3 lag4 error .000000273 .042626307 .132115292 2.737040606 .065486271 .065486271 APPROXIMANT FUNCTION N (1, 2) = $C_{l_{\alpha}}$ A0 A1 λ2 λ3 λ4 Α5 Α6 170329.1714 -4.694754 1.921693 -170317.2824 -3.613661 -4.501366 23.17 lag1 lag2 lag3 lag4 error .000000273 .042626307 .132115292 2.737040606 .181049841 .306154896 APPROXIMANT FUNCTION N (2, 1) = C_{m_h} λο λ1 λ2 λ3 λ4 λ5 λ6 3001.540919 -6.292480 .797560 -3001.600689 .227531 -.334033 14.757 lag2 lag3 lag4 error lagl .000000273 .042626307 .132115292 2.737040606 .059618785 .125105055 APPROXIMANT FUNCTION N (2, 2) = $C_{m_{\alpha}}$ λ1 **A**0 እ2 እ3 A4 A5 A6 -6706.7085 -3.347323 .739087 6706.26103 .105898 -.565209 5.315633 lag1 lag2 lag3 lag4 error .000000273 .042626307 .132115292 2.737040606 .159675096 .465829992

Appendix F : Function Optimization for Transonic Regime

Nasa-Langley approximate in Laplace form

$$[N] = [A_0] + \frac{b}{U} [A_1] s + (\frac{b}{U})^2 [A_2] s^2 + [A_3] \frac{s}{s + \frac{U}{b}b_1} + [A_4] \frac{s}{s + \frac{U}{b}b_2} + [A_5] \frac{s}{s + \frac{U}{b}b_3} + [A_6] \frac{s}{s + \frac{U}{b}b_4}$$

Approximation of Transonic Coefficients of Lift and Moment with respect to plunge and pitch, for a reduced frequency range of $0.0 \le k \le 0.5$ with forced constraints for C_{l_h} and C_{m_h} at k = 0.0, at M = 0.85.

$$\begin{bmatrix} C_{l_h} & C_{l_a} \\ C_{m_h} & C_{m_a} \end{bmatrix} = \begin{bmatrix} 0.0000 & 14.236 \\ 0.0000 & -0.5168 \end{bmatrix} + \frac{b}{U} \begin{bmatrix} 7.3520 & 0.9578 \\ -1.9059 & -4.3942 \end{bmatrix} s + (\frac{b}{U})^2 \begin{bmatrix} -0.2652 & 1.25818 \\ -0.2936 & 2.7440 \end{bmatrix} s^2 + \begin{bmatrix} -8561.67 & 404976 \\ 158360 & 217699 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27893)} + \begin{bmatrix} 0.5930 & -8.14577 \\ -0.05984 & 0.16511 \end{bmatrix} \frac{s}{s + \frac{L}{b}(0.04305)} + \begin{bmatrix} 8707.6 & -408964 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & 3987.4 \\ 1543.9 & 2184.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28587 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & 3987.4 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & -2184.1 \\ -1543.9 & 2184.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28587 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.23 & -2184.1 \\ -1543.9 & 2184.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.27899)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2184.1 \\ -159903 & -219881 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2184.1 \\ -15980 & -2184.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -15980 & -2184.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -1584.1 & -2884.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -1584.1 & -2884.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -1584.1 & -2884.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -1584.1 & -2884.1 \end{bmatrix} \frac{s}{s + \frac{U}{b}(0.28587)} + \begin{bmatrix} -146.28 & -2884.1 \\ -158$$

note : lag terms are equivalent

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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION TRANSONIC REGIME M=0.85 GLOBAL OPTIMIZATION FORCED CONSTRAINTS N(r,c)=0.0 AT k=0.0 RANGE 0.00 < k < 0.50 ERROR #1

 C_{l_h}

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k	EXACT		APPROXIMATE		k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
.000000,	.00000,	.00000,	.00000,	.00000,	.000000.	14.70000.	.00000.	14,23567.	.00000.
.025000,	.17000,	.56000,	.16792,	.54877,	.025000.	12.20000,	-4.05000,	12.21288.	-3.64541.
.050000,	.39000,	.87000,	.41847,	.86344,	.050000,	9.50000,	-4.15000,	9.56427,	-4.30539,
.087500,	.66000,	1.17000,	.69680,	1.16498,	.087500.	7.63000.	-3.70000.	7,53586	-3.72869
.100000,	.76000,	1.23000,	.77287,	1.24726,	.100000.	7.22000,	-3.47000,	7.15123.	-3.52318.
.125000,	.89000,	1.39000,	.90794,	1.39206,	.125000,	6.61000,	-3.18000,	6.57096,	-3.17275,
.150000,	1.01000,	1.49000,	1.02060,	1.51454,	.150000,	6.09000,	-2.94000,	6.12655,	-2.88115,
.175000,	1.12000,	1.61000,	1.10957,	1.62079,	.175000,	5.72000,	-2.59000,	5.75592,	-2.61606,
.200000,	1.17000,	1.69000,	1.17459,	1.71761,	.200000,	5.41000,	-2.29000,	5.43806,	-2.35648,
.250000,	1.25000,	1.93000,	1.23994,	1.90689,	.250000,	4.92000,	-1.79000,	4.94237,	-1.82723,
.300000,	1.23000,	2.12000,	1.24034,	2.11637,	.300000,	4.63000,	-1.30000,	4.62099,	-1.29981,
.350000,	1.19000,	2.39000,	1.20249,	2.35917,	.350000,	4.43000,	~.83000 ,	4.44320,	81140,
.400000	1.15000,	2.64000,	1.14690,	2.63547	.400000,	4.35000,	41000,	4.36678,	38607,
.450000,	1.07000,	2.93000,	1.08659,	2.94005,	.450000,	4.35000,	03000,	4.35258,	03010,
.500000,	.98000,	3.25000,	1.02883,	3.26655,	.500000,	4.40000,	.28000,	4.37100,	.26139,

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ĸ	EXACT		APPROXIMATE		k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
.000000,	.00000,	.00000,	.00000,	.00000,	.000000,	44500,	.00000,	51684,	.00000,
.025000,	.00100,	03500,	00729,	03935.	.025000,	45000,	02700,	48373,	04150,
.050000,	.01500,	06500,	00727.	06574.	.050000,	45500,	12200,	45477,	14213,
.087500,	.02500,	12500,	.01171,	11930,	.087500,	49500,	31100,	48600,	31323,
.100000,	.03000,	14500,	.01725,	14295,	.100000,	52000,	34800,	51147,	36418,
.125000,	.02000,	19800,	.02145,	19500,	.125000,	58100,	42200,	57665,	45028,
.150000,	.01500,	23000,	.01512,	24731,	.150000,	67200,	50800,	65200,	51226,
.175000,	.01000,	28000,	.00030,	29389,	.175000,	73200,	54800,	72737,	55013,
.200000,	00500,	32500,	01874,	33146,	.200000,	81000,	58200,	79435,	56754,
.250000,	03500,	39700,	05219,	37922,	.250000,	90500,	61000,	88350,	56578,
.300000,	06000,	40300,	06191,	40363,	.300000,	89800,	- .55000,	90772,	55768,
.350000,	04500,	42500,	04310,	42338,	.350000,	86800,	57500,	88429,	57878,
.400000,	00500,	44800,	00139,	45102,	.400000,	82400,	64800,	83843,	64393,
.450000,	.06500,	49000,	.05504,	49191,	.450000,	79000,	75000,	79150,	75384,
.500000,	.14000,	55000,	.11927,	54679,	.500000,	77000,	90300,	75798,	90251,

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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION TRANSONIC REGIME M=0.85 GLOBAL OPTIMIZATION FORCED CONSTRAINTS RANGE 0.00 < k < 0.50 ERROR #1



k EXACT APPROXIMATE k EXACT APPROXIMATE REAL IMAG REAL IMAG REAL IMAG REAL IMAG .000000, .00000, .00000, .00000, .00000, .000000, 14.70000, .00000, 14.70000, .00000, .55795, .56000, .19383, .025000, .17000. .025000, 12.20000, -4.05000, 12.16128, -3.97370, .87000, .43656, .050000, .39000, .85319, .050000, 9.50000, -4.15000,9.41145, -4.23129, .087500, .66000, 1.17000, .70846, 1.17109. .087500, 7.63000, -3.70000, 7.67286, -3.66816, 1.23000, .79040, 1.25555, .100000, .76000, .100000, 7.22000, -3.47000, 7.29733, -3.52885, 1.39000, .93541, 1.39534, .125000, .89000, .125000, 6.61000, -3.18000, 6.66079, -3.25681, 1.49000, 1.04994, 1.50834, .150000, 1.01000, 6.09000, .150000, -2.94000, 6.14647, -2.96663, 1.12000, 1.61000, 1.13440, 1.60795, 5.72000, -2.59000, 5.73753, -2.66267, .175000, .175000, 1.69000, 1.19295, 1.70330, 5.41000, -2.29000, 5.41665, -2.35786, .200000, 1.17000, .200000, 1.93000, 1.25110, 1.90029, .250000, 4.92000, -1.79000, 4.97158, -1.77935, .250000, 1.25000, 2.12000, 1.25543, 2.11892, .300000, 4.63000, -1.30000, 4.70202, -1.26197, .300000, 1.23000, 1.19000, 2.39000, 1.22593, 2.36403, .350000, 4.43000, -.83000, 4.54025, -.80489, .350000, -.39879, 1.17512, 2.63603, .400000, 4.35000, -.41000, 4.44697, .400000, 1.15000, 2.64000, 4.35000, -.03000, 4.39915, -.03395, .450000, .450000, 1.07000, 2.93000, 1.11128, 2.93413, .29772, 3.25724, .500000, 4.40000, .28000, 4.38285, .500000, .98000, 3.25000, 1.04035,

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k	EXACT		APPROXIMATE		k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
. 000000 .	.00000.	.00000.	.00000.	.00000,	.000000,	44500,	.00000,	44500,	.00000,
.025000.	.00100.	03500,	04127,	05079,	.025000,	45000,	02700,	47202,	06448,
.050000.	.01500,	06500,	02884,	04927,	.050000,	45500,	12200,	42613,	12063,
.087500.	.02500.	12500.	.01021	12957.	.087500,	49500,	31100,	43443,	33035,
100000	.03000.	14500.	.00937	16165.	.100000,	52000,	34800,	47130,	39057,
125000	.02000.	19800.	00401.	~.21692.	.125000,	58100,	42200,	56157,	47416,
150000	01500	- 23000	02225	25906.	.150000,	67200,	50800,	64879,	51749,
175000	01000	- 28000	03854	29126.	.175000,	73200,	54800,	72016,	53602,
200000	- 00500	- 32500	- 05051	- 31708	.200000	81000,	58200,	77337,	54202,
.200000,	- 03500,	- 39700	- 06049	- 35882	.250000.	90500,	61000,	83305,	54527,
.250000,	- 06000	- 40300	- 05448	- 39514	.300000.	89800.	55000,	84809,	56012,
.300000,	- 04500	- 40500,	- 03590	- 42997	350000	86800.	57500.	83581,	59821,
.350000,	- 04500,	42500,	- 00707	- 16/18	400000	82400.	64800.	80830.	66360,
.400000,	00500,	44800,	00707,	40410,	.400000,	~ 79000	- 75000	77420.	75757.
.450000,	.06500,	49000,	.03062,	49834,	.450000,	- 77000,	- 90300	- 74002	- 88020.
.500000,	.14000,	55000,	.07634,	53233,	.500000,	//000,	=.90300,	.,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	

LAPLACE	METHOD	TRANSC FORCEE RANGE	ONIC M=0.85 O CONSTRAINTS 0.00 < k <	0.50	
u	r(alp)	ah	wben/wtor		Eigenvalue Analysis
50.0	. 5	5	. 2		
ND VELOCI	LTY .500				
	14470	1.281	135		
	14470	-1.281	135		
	60746	.000	000		
	67443	.000			
	00511	201	192 192		
	05524	002	255		
	05524	.002	255		
	02989	.003	22		
	02989	003	22		
	02296	.002	203		
perfomar	ice index is	002	203 75297	.00000	
ND VELOCI	TTV 3 300				
	-6.29379	.000	000		
-	4.01529	.000	00		
	19525	1.253	302		
	19525	-1.253	302		
	00033	301	22		
	00033	.301	22		
	34065	180)/4)74		
	17917	110)29		
	17917	.110	29		
	06246	060	04		
•	06246	.060	004		
periomar	ice index is	2.7	2761	.00000	
ND VELOCI	TY 3.400				
-	·6.49084	.000	000		
-	- 20562	.000	000		
	20562	-1.262	75		
	.00355	305	563	INSTABILITY	
	35100	188	800	ND VELOCITY	= 3.31
	35100	.188	300		
	.00355	.305	63		
	18377	113	162		
	103//	.113	102		
	06142	062	;/⊥)71		
perfoman	ce index is	4.3	1713	.00000	
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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION TRANSONIC REGIME M=0.85 GLOBAL OPTIMIZATION FORCED CONSTRAINTS N(r,c)=0.0 AT k=0.0 RANGE 0.00 < k < 1.00 ERROR #1

 C_{l_h}

k	ЕХАСТ		APPROXIMATE		k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
.000000.	.00000.	.00000,	.00000,	.00000,	.000000,	14.70000,	.00000,	13.9473U,	.00000,
.025000,	.17000,	.56000,	.16253,	.54899,	.025000,	12.20000,	-4.05000,	12.17360,	-3.42799,
.050000,	.39000,	.87000,	.41431,	.86878,	.050000,	9.50000,	-4.15000,	9.67237,	-4.23117,
.087500.	.66000.	1.17000.	.69195,	1.16922,	.087500,	7.63000,	-3.70000,	7.57972,	-3.78514,
.100000,	.76000,	1.23000,	.76583,	1.25158,	.100000,	7.22000,	-3.47000,	7.16425,	-3.58634,
.125000,	.89000,	1.39000,	.89576,	1.39867,	.125000,	6.61000,	-3.18000,	6.53843,	-3.22036,
.150000,	1.01000,	1.49000,	1.00408,	1.52567,	.150000,	6.09000,	-2.94000,	6.07511,	-2.89669,
.175000.	1.12000	1.61000.	1.09045.	1.63721.	.175000,	5.72000,	-2.59000,	5.70821,	-2.60052,
200000	1,17000.	1.69000.	1.15459.	1.73882,	.200000,	5.41000,	-2.29000,	5.40820,	-2.31982,
.250000,	1.25000,	1.93000,	1.22118,	1.93316,	.250000,	4.92000,	-1.79000,	4.95806,	-1.78414,
. 300000.	1.23000.	2.12000	1.22287.	2.14008.	.300000,	4.63000,	-1.30000,	4.66424,	-1.27997,
350000	1,19000	2.39000	1.18306.	2.37316.	.350000.	4.43000,	83000,	4.48829,	81771,
400000	1,15000	2.64000	1.12048.	2.63369.	.400000,	4.35000,	41000,	4.39670,	40553,
450000	1.07000.	2.93000,	1.04718,	2.91750,	.450000,	4.35000,	03000,	4.36154,	04479,
500000.	.98000.	3.25000,	.96983,	3.21913,	.500000,	4.40000,	.28000,	4.36154,	.26823,
550000	92000	3.54000	.89166.	3.53366.	.550000,	4.43000,	.55000,	4.38166,	.53973,
	88000	3.89000	.81392.	3.85723.	.600000,	4.50000,	.74000,	4.41178,	.77649,
650000,	.79000.	4.20000,	.73684,	4.18702,	.650000,	4.50000,	1.00000,	4.44540,	.98490,
.700000.	.68000,	4.52000,	.66016,	4.52100,	.700000,	4.48000,	1.18000,	4.47847,	1.17048,
750000	.57000.	4.88000	. 58345.	4.85778,	.750000,	4.48000,	1.31000,	4.50862,	1.33784,
800000	.47000.	5.22000.	.50621.	5.19641,	.800000,	4.48000,	1.47000,	4.53451,	1.49071,
.850000.	.39000,	5.53000,	.42799,	5.53621,	.850000,	4.49000,	1.63000,	4.55549,	1.63206,
.900000,	.29000,	5.87000,	.34837,	5.87675,	.900000,	4.56000,	1.80000,	4.57130,	1.76425,
.950000.	.20000.	6.25000.	.26700	6.21770,	.950000,	4.60000,	1.91000,	4.58194,	1.88914,
1.000000,	.10000,	6.58000,	.18357,	6.55886,	1.000000,	4.62000,	2.02000,	4.58755,	2.00820,

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C_{m_h}

k	EXAC	T	APPROXI	APPROXIMATE k		APPROXIMATE k EXA		EXAC	ACT APPRO		OXIMATE	
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG			
.000000,	.00000,	.00000,	.00000,	.00000,	.000000,	44500,	.00000,	57389,	.00000,			
.025000,	.00100,	03500,	.01514,	- 01722,	.025000,	45000,	02700,	50236,	00277,			
.050000,	.01500,	06500,	.03491,	05826,	.050000,	45500,	12200,	43591,	12894,			
.087500,	.02500,	12500,	.04444,	13429,	.087500,	49500,	31100,	48192,	34358,			
.100000,	.03000,	14500,	.04338,	15945,	.100000,	52000,	34800,	51936.	39784.			
.125000	.02000,	19800,	.03688,	20736,	.125000	58100.	42200,	60391.	47587,			
.150000,	.01500,	23000,	.02622,	25074,	.150000,	€7200,	50800,	68543,	51920,			
.175000,	.01000,	28000,	.01345,	28853,	.175000,	73200,	54800,	75284,	53819,			
.2000 00,	00500,	32500,	.00069,	32037,	.200000,	81000,	58200,	80187,	54342,			
.250000,	03500,	39700,	01725,	36818,	.250000,	90500,	61000,	84842,	54545,			
.300000,	06000,	40300,	01698,	40259,	.300000,	89800,	55000,	84946,	56636,			
.`50000,	04500,	42500,	.00424,	43438,	.350000,	86800,	57500,	83350,	61570,			
.400000,	00500,	44800,	.04351,	47188,	.400000,	82400,	64800,	81867,	68765,			
.450000,	.06500	49000,	.09567,	51964,	.450000,	79000,	75000,	81277,	77282,			
.500000,	.14000,	55000,	.15561,	57908,	.500000,	77000,	90300,	81723,	86341,			
.550000,	.21000,	64000,	.21928,	64973,	.550000,	81800,	99200,	83063,	95435,			
.600000,	.27000,	76000,	.28384,	73020,	.600000,	85200,	-1.08000,	85067,	-1.04290,			
.650000,	.34000,	86000,	.34752,	81879,	.650000,	- 87800,	-1.14300,	87516,	-1.12786,			
.700000,	.39000,	97000,	.40931,	91387,	.700000,	89200,	-1.23000,	90235,	-1.20894,			
.750000,	.46000,	-1.06000,	.46873,	-1.01398,	.750000,	93000,	-1.30700,	93092,	-1.28636,			
.800000,	.53000,	-1.13000,	.52566,	-1.11794,	.800000,	97800,	-1.38100,	95998,	-1.36052,			
.850000,	.57000,	-1.19000,	.58017,	-1.22475,	.850000,	-1.00800,	-1.43900,	98895,	-1.43193,			
.900000	.61000.	-1.28000	.63245,	-1.33366,	.900000	-1.03100,	-1.49500,	-1.01747,	-1.50106,			
.950000	.65000,	-1.40000,	.68275,	-1.44407,	.950000,	-1.04300,	-1.53400,	-1.04534,	-1.56834,			
1.000000,	.68000,	-1.53000,	.73133,	-1.55552,	1.000000,	-1.05800,	-1.56400,	-1.07247,	-1.63414,			

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LAPLACE N	1ETHOD	TRAN FORC RANG	TRANSONIC M=0.85 FORCED CONSTRAINTS AT N(r,c)=0.0 AT k=0.0 RANGE 0.00 < k < 1.00				
u	r(alp)	ah	wben/wto:	r	Eigenvalue Analysis		
50.0	.5	5	. 2				
ND VELOCIT	FY . 500						
-	01753	-1.1	7101				
-	01753	1.1	7101				
-	00509	.2	0181				
-	00509	2	0181				
-	19517	.0	3653				
-	17830	.0	3354				
	1951/	0	3653				
-	13698	0	3334				
-	12781	.0	0000				
-	02340	.0	0000				
-	02273	.0	0000				
perfomanc	ce index is		.75683	.00000			
ND VELOCIT	CY 3.100						
-	19085	1.2	3731				
-	19085	-1.2	3731				
-]	1.36708	.6	9722				
- 1	16692	6	9722				
L –	14692	.2	5217				
د ب	74436	2	5217				
-	00338	.0	9255				
-	.00338	2	9255				
-	.35884	.0	0000				
-	20286	.0	0000				
-	07687	. 0	0000				
perfomanc	ce index is	8	.89727	.00000			
ND VELOCIT	Y 3.200						
-	.19992	1.2	4432				
-	•.19992	-1.2	4432				
-1	41678	. 7	2501				
-1	20460	7	2501				
-1	.20460	- 2	6032				
-	.76779	- • 2 . 0	0000				
	.00058	. 2	9752	INSTABILITY			
	.00058	2	9752	ND VELOCITY =	= 3.19		
-	.36035	. 0	0000				
-	.21351	.0	0000				
-	.07676	.0	0000				
periomanc	e index 15	15	.95851	.00000			

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SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION TRANSONIC REGIME M=0.85 GLOBAL OPTIMIZATION FORCED CONSTRAINTS RANGE 0.00 < k < 1.00

k	EXAC	EXACT		APPROXIMATE		EXAC	T APPROXIMATE		
	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
.000000,	.00000,	.00000,	.00000,	.00000,	.000000,	14.70000,	.00000,	14.70000,	.00000,
.025000,	.17000,	.56000,	.20041,	.56379,	.025000,	12.20000.	-4.05000.	12.09142.	-4.07232
.050000.	. 39000.	.87000.	.43868,	.85147,	.050000.	9.50000.	-4.15000.	9.23969.	-4.20057.
.087500,	.66000,	1.17000,	.70126,	1.18815,	.087500,	7.63000.	-3.70000.	7.67001,	-3.60840,
.100000,	.76000,	1.23000,	.78739,	1.27735,	.100000,	7.22000,	-3.47000,	7.30929,	-3.51304,
.125000,	.89000,	1.39000,	.94081,	1.41745,	.125000,	6.61000	-3.18000.	6.64489,	-3.30380,
.150000,	1.01000,	1.49000,	1.05738,	1.52418,	.150000,	6.09000.	-2.94000,	6.08270,	-3.02828,
.175000,	1.12000,	1.61000,	1.13708,	1.61709,	.175000,	5.72000,	-2.59000,	5.63848,	-2.70889,
.200000,	1.17000,	1.69000,	1.18649,	1.70825,	.200000,	5.41000,	-2.29000,	5.30123,	-2.37741,
.250000,	1.25000,	1.93000,	1.22244,	1.90616,	.250000,	4.92000,	-1.79000,	4.86628,	-1.75049,
.300000,	1.23900,	2.12000,	1.20838,	2.13362,	.300000,	4.63000,	-1.30000,	4.63458,	-1.21168,
.350000.	1.19000.	2.39000.	1.16854,	2.38805,	.350000,	4.43000,	83000.	4.51518,	76053,
.400000,	1.15000,	2.64000,	1.11535,	2.66342,	.400000,	4.35000,	41000,	4.45752,	38090,
.450000,	1.07000,	2.93000,	1.05501,	2.95416,	.450000,	4.35000,	03000,	4.43406,	05620,
.500000,	.98000,	3.25000,	.99046,	3.25592,	.500000,	4.40000,	.28000.	4.42969,	.22711.
550000	92000	3.54000	.92301.	3.56554	.550000	4.43000.	.55000.	4.43604.	.47927.
.600000.	.88000.	3.89000.	.85313.	3.88073.	.600000	4.50000,	.74000,	4.44844,	.70785
.650000.	.79000,	4.20000,	.7808.,	4.19990,	.650000,	4.50000,	1.00000.	4.46424.	.91840
.700000.	.68000.	4.52000.	.70617.	4.52191	.700000	4,48000	1,18000.	4.48194.	1,11500.
750000	.57000.	4.88000.	. 62879 .	4.84597.	.750000	4.48000.	1.31000	4.50072.	1.30069.
.800000.	.47000.	5.22000.	.54853.	5.17148.	.800000,	4.48000.	1.47000,	4.52013,	1.47773,
.850000.	.39000.	5,53000.	.46518.	5.49806.	.850000.	4.49000.	1.63000.	4.53991.	1.64784.
.900000	.29000.	5.87000	.37857.	5.82540.	.900000.	4.56000	1.80000	4.55997.	1.81234.
.950000	.20000	6.25000	.28853	6.15329	.950000	4.60000	1.91000	4.58025.	1.97223.
1.000000,	.10000,	6.58000,	.19492,	6.48156,	1.000000,	4.62000,	2.02000,	4.60076,	2.12832,

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k	EXACT		APPROXIMATE		k	EXACT		APPROXIMATE	
ĸ	REAL	IMAG	REAL	IMAG		REAL	IMAG	REAL	IMAG
00000	. 00000 .	.00000.	.00000.	.00000,	.000000.	44500,	.00000,	44500,	.00000,
.025000.	.00100.	03500,	08536,	07477,	.025000,	45000,	02700,	45251,	05573,
.050000.	.01500.	06500.	07246.	03012.	.050000.	45500,	12200,	40544,	12522,
087500	02500	12500.	.00470.	13296.	.087500.	49500,	31100,	42450,	33564,
100000	.03000.	14500.	00026.	17806.	.100000.	52000,	34800,	46298,	39195,
.125000.	.02000.	19800,	03268,	24522,	.125000,	58100,	42200,	55008,	46639,
. 150000.	.01500.	23000,	06879,	28304,	.150000,	67200,	50800,	62755,	50379,
.175000.	.01000.	28000,	09475,	30307,	.175000,	73200,	54800,	68670,	52187,
.200000.	~ 00500.	32500.	10802.	31524.	. 200000.	81000,	58200,	72856,	53235,
250000	03500.	39700.	10344.	33744.	.250000,	90500,	61000,	77490,	55295,
.300000.	06000.	40300.	07271.	37041,	. 300000,	89800,	55000,	79240,	58606,
. 350000.	04500.	42500,	02907.	41710,	.350000.	86800,	57500,	79722,	63383,
400000	~.00500.	44800.	.02019.	47542.	.400000.	82400,	64800,	79794,	69385,
450000	.06500.	49000.	.07148.	54248,	.450000,	79000,	75000,	79890,	76310,
.500000.	.14000.	55000,	.12324,	61574,	.500000,	77000,	90300,	80217,	83905,
550000	21000	64000	.17495.	69328.	.550000	81800.	99200,	80866,	91976,
	27000	- 76000	22659	- 77370.	. 600000.	85200,	-1.08000,	81870,	-1.00384,
.600000,	34000	- 86000,	.27836.	85600.	.650000,	87800,	-1.14300,	83234,	-1.09027,
700000	39000	- 97000	. 33054.	93951.	.700000.	89200.	-1.23000,	84948,	-1.17836,
.700000,	. 19000,	-1.06000	38344	-1 02376	.750000	93000.	-1.30700,	86999,	-1.26760,
./50000,	.40000,	-1.00000,		-1 10842	.800000.	97800,	-1.38100,	89371,	-1.35764,
.800000,	.53000,	-1.13000,	49735,	-1.19328	.850000.	-1.00800.	-1.43900,	92048,	-1.44823,
.850000,	£1000,	-1 28000	54905	-1.27819	.900000	-1.03100.	-1,49500,	95018,	-1.53920,
.900000,	.61000,	-1 40000	60729	-1.36306.	.950000.	-1.04300,	-1.53400,	98269,	-1.63041,
.950000,	.65000,	-1 53000	.66733.	-1.44782.	1.000000,	-1.05800,	-1.56400,	-1.01789,	-1.72178,

LAPLACE I	METHOD	TRAN FORC RANG	SONIC M=0 ED CONSTRAI E 0.00 <)	.85 INTS < < 1.00	
u	r(alp)	ah	wben/wtoi	c	Eigenvalue Analysis
50.0	.5	- .5	. 2		
ND VELOCI	TY .500				
	02388	1.1	7850		
	02388	-1.1	7850		
•	00516	2	0181		
•	00516	. 2	0181		
	17990	.0	0227		
•	17990	0	0227		
•	04345	0	0962		
•	04345	.0	0962		
•	03331	.0	1123		
-	03331	0	1123		
•	02539	.0	0376		
•	02539	0	0376		
perfomanc	ce index is		.78187	.00000	
ND VELOCI					
ND VELOCI	- 21772	-1 2	1977		
	- 21773	-1.2	40/3		
- 1	28855	- 0	40/3		
-	1.28855		1047		
-	34557	2	5499		
-	00193	2	9898		
-	34557	.2	5499		
-	00193	.2	9898		
-	15908	1	3442		
•	04041	0	7268		
•	15908	.1	3442		
-	04041	.0	7268		
perfomanc	ce index is	8	.34966	.00000	
ND VELOCI					
NU VELOCIT	- 22028	-1 7	5504		
	- 22930	-1.2	5504		
_1	22930	1.2	0000		
-1	1.31940	.0	0000		
-	35652	- 2	6509		
	.00131	3	0338	TNOTARTITTY	
-	35652	J _)	6509	ND VELOCITY = 3	. 36
	.00131	.3	0338	the threatthe d	
-	16341	1	3823		
-	16341	.1	3823		
-	03761	0	7519		
-	03761	.0	7519		
perfomanc	ce index is	8	.98850	.00000	
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$\label{eq:appendix G} Appendix \ G: Incompressible \ Regime \ - \ Eigenvalue \ Solutions$

Eigenvalue solutions for the incompressible regime (Case #1) incorporating the l' = g, p = k and Laplace methods.

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u	r(alp)	ah wi	oen/wtor		
50.0	. 5	5	. 2		
k	U/b+wtor	w/wto:	r g	damping ra	tio
10.000,	.01970,	.19698	00209,	.00104,	
10.000,	. 11493,	1.14932	00569,	.00284,	
6.000.	.03283,	.19699	00348	.00174,	
6.000,	.19150,	1.14900	00948.	.00474,	
4.000.	.04925.	.19701	00524.	.00262	
4.000	.28709	1.14837	01423.	.00712	
2.000.	.09856	. 19711	01067.	.00533.	
2.000	. 57247.	1.14494	02861.	.01431.	
1.000.	. 19746 .	. 19746	02238.	.01119.	
1.000.	1.13102.	1,13102	05796.	.02898.	
. 880	. 22452	. 19758	02580	.01290.	
. 880.	1.27906	1,12558	06604	.03302.	
.800	.24711	. 19769	02871.	.01436.	
. 800.	1.40063.	1,12050	07275.	.03638.	
.600	33024	19814	- 03995	.01998	
. 600	1.82975	1 09785	- 09686	04843	
500	39712	19856	- 04954	02477	
500,	2 15096	1 07548	-11404	05747	
400	10815	10075	- 06495	03243	
400,	2 59976	1 03500	- 13957	069293,	
300	2.30370,	20062	- 00274	.00929,	
. 300,	.000/4,	.20002	-16176	.04037,	
. 300,	1 02040	.95696	, - 16760	.00230,	
.200,	1.02040,	.20410	-16100	.07660,	
.200,	3.9//29,	./9540	,10100,	.08053,	
.150,	1.39104,	.208/5	,23812,	.11906,	
.150,	4.35545,	.65332	,09675,	.04837,	
.140,	1.50237,	.21033	,26508,	.13254,	
. 140,	4.41935,	.61871	,07038,	.03519,	
.130,	1.63333,	.21233	,29866,	.14933,	
.130,	4.47879,	.58224	,03713,	.01856,	
.125,	1.70828,	.21354	,31888,	.15944,	
.125,	4.50682,	.56335	,01709,	.00854,	
.120,	1.79092,	.21491,	,34199,	.17100,	
.120,	4.53418,	.54410	, .00554,	00277,	INSTAE
.110,	1.98367,	.21820	,40100,	.20050,	U/b*wt
.110,	4.59055,	.50496	, .06404,	03202,	
.100,	2.22456,	.22246	,48395,	.24197,	
.100,	4.66272,	.46627	, .14335,	07167,	
.050,	4.97074,	.24854	, -1.75199,	.87599,	
.050,	7.19299,	.35965,	, 1.11015,	55508,	

INSTABILITY U/b*wtor = 4.53 ₹. ¥:

METHOD	CASE	#1
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l	1

PK

50.0

r(alp)

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wben/wtor

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k U/b*wtor damping damping ratio W .5000, .3966, -.00640, .19917, .03212 2.2909, .5000, -.01436, 1.14593, .01253 .2034, 1.0000, -.01487, .20335, .07295 1.1348, 1.0000, -.02958, 1.13484, .02606 .1394, 1.5000, -.02492, .20916, .11830 .7434, 1.5000, -.04621, 1.11514, .04141 .1087, 2.0000, -.03677, .21735, .16679 2.0000, .5425, -.06450, 1.08521, .05933 .0916. 2.5000. .22893, -.05169, .22026 .4169, 2.5000, -.08437, 1.04258, .08066 .3275, 3.0000, -.10547, .98336, .10664 3.0000, .0820, -.07180, .24614, .28002 .2570, 3.5000, -.12613, .89945, .13887 .0781, 3.5000, -.10247, .35074 .27359, .1923, 4.0000, .76981, -.13992, .17882 .0809, 4.0000, -.16167, .32367, .44684 .1784, 4.1000, -.13846, .73202, .18586 .0821, 4.1000, -.18312, .33747, .47693 .1633, 4.2000, -.13313,.68655, .19036 .0838, 4.2000, -.21078, .35084, .51500 .1456, 4.3000, -.11102, .62712, .17433 .0833, 4.3000, -.24824, .35793, .56990 .1306, 4.4000, -.05886, .57403, .10201 4.4000, .0805, -.29190, .35396, .63623 .1228, 4.5000, -.01093, .55238, .01977 .0749, -.33777, 4.5000, .33681, .70811 .1177, 4.6000, .02507, .54086, -.04629 INSTABILITY 4.6000, .0678, -.38361, .31138, .77642 U/b * wtor = 4.534.7000, .1133, .05465, .53225, -.10213.0599, 4.7000, -.43058, .28056, .83783 .1094, 4.8000. .08015, .52451, -.15106 .0499, 4.8000, -.48151, .23879, .89589 .1056, 4.9000, .10301, .51721, -.19532 .0390, 4.9000, **-.**53891, .19045, .94285 .1020, 5.0000, .12384, .50991, -.23600 5.0000, .0284, -.60286, .14083, .97378 .0986, 5.1000, .14317, .50263, -.27394 .0208, 5.1000, -.66517, .10505, .98776

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u	r(alp)	ah	wben/wtor		Eigenvalue Analysis
50.0	.5	5	.2		
ND VELOCIT	Y.50				
01437		1.14595			
01437		-1.14595			
00637		19921			
00637		.19921			
38147		.00000			
38256		.00000			
12656		.00000			
12798		.00000			
04200		.00000			
- 04214		.00000			
~.00772		.00000			
perfomance	index is	1	.13157	.00000	
-					
ND VELOCIT	Y 1.00				
02966		1.13486			
02900		-1.13480			
- 76511		.00000			
01482		.20360			
01482		20360			
24879		.00000			
25595		.00000			
08316		.00000			
08429		.00000			
01544		.00000			
01543	inden in	.00000			
periomance	index is	L		.00000	
ND VELOCIT	Y 1.50				
04649		-1.11520			
04649		1.11520			
-1.14241		.00000			
-1.14767		.00000			
-,02480		.20984			
02480		20984			
36693		.00000			
38393		.00000			
- 12642		.00000			
- 02316		.00000			
02313		.00000			
perfomance	index is		.29740	.00000	
-				-	
ND VELOCIT	¥ 2.00	1			
-1 500529		1.08238			
-1 52220		.00000			
-1.23053		-1 00520			
00529		-T.09238			
48020					
140020			4 - 31		

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03664	.21879	
03664	21879	
15957	.00000	
16857	.00000	
03088	.00000	
- 03080		
nerfomance index is	.00000	
perfoliance filder 13	20.13492	.00000
ND VELOCITY 2 50		
	00000	
	.00000	
- 08644	-1 04280	
- 09644	-1.04289	
- 63000	1.04289	
- 59710	.00000	
56/18	.00000	
05141	23195	
- 30325	.23195	
19235	.00000	
210/2	.00000	
03860	.00000	
03843	.00000	
perfomance index is	2.70799	.00000
ND VELOCITY 3.00		
-2.28251	.00000	
-2.29534	.00000	
11024	98348	
11024	.98348	
76786	.00000	
68609	.00000	
07108	.25240	
07108	25240	
25286	.00000	
21817	.00000	
04632	.00000	
04600	.00000	
perfomance index is	2.86213	.00000
ND VELOCITY 3.50		
-2.66251	.00000	
-2.67790	.00000	
13678	.89843	
13678	89843	
89583	.00000	
77454	.00000	
09940	.28811	
09940	28811	
29500	.00000	
23173	.00000	
05404	00000	
05348		
perfomance index is	2 21510	00000
Lananda Tudev 12	6 · 6 I J I V	
ND VELOCITY 4 00		
-3.04253	00000	
-3 06045		
- 16405	- 76005	
- 16405		
10430 -1 0000	. / 5095	
- 94967	.00000	
0480/	.00000	

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14324		36783	
14324		.36783	
33714		.00000	
22548		.00000	
06176		.00000	
06081		.00000	
perfomance	index is	2.43926	.00000
perrendice			
ND VELOCITY	4.10		
-3,11854		. 00000	
-3,13696		.00000	
- 16997		.71810	
- 16997		- 71810	
-1 04940		00000	
- 96122			
- 15509		39967	
- 15509		- 20967	
10009		39867	
34557		.00000	
2215/		.00000	
06331		.00000	
06225	1	.00000	
pertomance	index is	2.98730	.00000
ND VELOCITY	4.20		
-3.19455		.00000	
-3.21348		.00000	
-1.07500		.00000	
87286		.00000	
17319		65974	
17319		.65974	
16957		.44503	
16957		44503	
35400		.00000	
21693		.00000	
06485		.00000	
06368		.00000	
perfomance	index is	2.48161	.00000
•			
ND VELOCITY	4.40		
-3.34656		.00000	
-3.36650		.00000	
05134		.55387	
32922		.52710	
-1,12619		.00000	
89292		.00000	
05134		55387	
- 32922		- 52710	
- 37086		00000	
- 20615	•		
- 06794		.00000	
00/74 06651		.00000	
00001	indau in	.00000	
periomance	INGEX 15	0.60422	.00000
ND VELOCITY	4.50		
-3.42257		.00000	
-3.44301		.00000	
-1.15178		.00000	
01016		54976	
39043		51942	
90108		.00000 A	- 33

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01016	.54976	
39043	.51942	
- 37929	00000	
- 20025	.00000	
20035	.00000	
06948	.00000	
06/91	.00000	
perfomance index is	2.05162	.00000
ND VELOCITY 4.60		
-3.49858	.00000	
-3.51952	.00000	
-1.17738	.00000	
.02177	.54546 I	NSTABILITY
44317	.51200 U	/b*wtor = 4.53
90778	.00000	
.02177	54546	
44317	51200	
38771	.00000	
06929	.00000	
07103	.00000	
19451	.00000	
perfomance index is	5.28303	. 00000
•		
ND VELOCITY 4.80		
-3.65060	.00000	
-3.67254	.00000	
.07249	.53602	
.07249	- 53602	
-1, 22857		
- 53799	10010	
- 91590	.49017	
- 53700	- 40917	
- 40457	4901/	
- 07300	.00000	
07200	.00000	
07412	.00000	
18314	.00000	
periomance index is	2.20679	.00000
ND VELOCIEV E CO		
-2 80262		
-3.80262	.00000	
-3 - 82557	.00000	
.11382	52557	
.11382	.52557	
-1.27976	.00000	
62745	48580	
62745	.48580	
91517	.00000	
42143	.00000	
17267	.00000	
07461	.00000	
07720	.00000	
perfomance index is	1.71618	. 00000
-		

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Appendix H : Transonic Regime - Laplace Generated Eigenvalues for M=0.85

This sections gives the Laplace generated eigenvalues for M = 0.85.

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LAPLACE M	IETHOD	TRANS	SONIC M=0.85	5	
u	r(alp)	ah	wben/wtor		Eigenvalue Analysis
50.0	.5	5	. 2		
ND VELOCIT	Y .500				
-	.08315	-1.2	7662 7662		
-	.00511	2	0185		
-	00511	. 2	0185		
-	.18737	0	1797		
-	•.18737	.0	1797		
-	12884 - 12884	0	3121		
-	•.11139	0	1459		
-	.11139	.0	1459		
-	02172	.0	0019		
-	.02172	0	0019	00000	
periomano	ce index is	T	.34192	.00000	
ND VELOCIT	ry 1.000				
-	13682	-1.2	2225		
-	13682	1.2	2225		
•	01143	- 2	1140		
-	47016	2	0000		
-	41480	.0	0000		
-	24390	0	8059		
-	24390	.0	8059		
-	19773 - 19773	0	3599		
-	04279	.0	0145		
	04279	0	0145		
perfomanc	ce index is	1	.91492	.00000	
ND VELOCIT	FY 1.500				
	14304	-1.1	7204		
•	14304	1.1	7204		
•	91862	.0	0000		
•	6094/ - 01825	.0	2490		
•	01825	2	2490		
•	35435	.1	3235		
•	35435	1	.3235		
•	27054	.0	5787		
	27054	0	0460		
	06248	0	0460		
perfomance	ce index is	3	.55296	.00000	
ND VELOCT	TY 2.000				
	13090	-1.1	.6269		
-:	1.42867	. 0	0000		
•	13090	1.1	.6269		
•	81080	. 0	0000		
·		• 4	4231 4-30		

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02269 46358 46358 33331 33331 08006 08006 perfomance index is	24257 .18321 18321 .07634 07634 01006 .01006 3.54698	.00000
ND VELOCITY 2.500 -1.93877 13300 13300 -1.01285 02228 02228 57278 57278 38902 38902 09496 perfomance index is	$\begin{array}{c} .00000 \\ -1.18652 \\ 1.18652 \\ .00000 \\ .26434 \\26434 \\ .23294 \\23294 \\ .08819 \\08819 \\08819 \\01771 \\ .01771 \\ 4.61556 \end{array}$. 00000
ND VELOCITY 3.000 -2.43536 15539 15539 -1.21510 68232 68232 01448 01448 01448 44104 44104 10696 10696 perfomance index is	.00000 1.22983 -1.22983 .00000 28189 .28189 28849 .09158 09158 .02704 02704 3.86013	. 00000
ND VELOCITY 3.100 -2.53320 16210 16210 -1.25556 70427 70427 01186 01186 45129 45129 10903 10903 perfomance index is	.00000 -1.24029 1.24029 .00000 29162 .29162 29332 .29332 09119 .09119 .02906 02906 5.05791	. 0 0000
NU VELOCITY 3.200 -2.63062 16948 16948 -1.29602 72623	$\begin{array}{r} .00000\\ 1.25130\\ -1.25130\\ .00000\\30133 & 4-37 \end{array}$	

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72623 00889 00889 46152 46152 11100 11100 perfomance index is	.30133 29807 .29807 09044 .09044 .03110 03110 5.61231	.00000
ND VELOCITY 3.300 -2.72764 17749 17749 -1.33649 74821 74821 00554 00554 47177 47177 11287 11287 perfomance index is	.00000 -1.26286 1.26286 .00000 31102 .31102 30271 .30271 08933 .08933 .03319 03319 24.33285	. 00000
ND VELOCITY 3.400 -2.82428 18609 18609 -1.37696 77020 77020 00184 48205 48205 11466 11466 perfomance index is	.00000 -1.27495 1.27495 .00000 .32070 32070 .30721 30721 .08786 08786 .03530 03530 16.47175	. 00000
ND VELOCITY 3.500 -2.92058 19525 -1.41743 19525 79221 .00221 .00221 .00221 49236 49236 11637 perfomance index is	.00000 1.28757 .00000 -1.28757 33037 .31152 31152 .08603 08603 .03744 03744 5.16679	INSTABILITY ND VEL = 3.45 .00000
ND VELOCITY 4.000 -3.39765 -1.61980 24794 24794 90240	.00000 .00000 -1.35846 1.35846 .37857 A-	38

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90240	37857	
.02714	32936	
.02714	.32936	
54475	07102	
54475	.07102	
12386	.04847	
12386	04847	
perfomance index is	2.57841	. 00000

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Appendix I : Transonic Regime - Comparison of Rational Function Optimization

Comparison of rational function optimization for M = 0.80, M = 0.85 and M = 0.875.

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TRANSONIC AERODYNAMICS

	M=0.85	M=0.80	M=0.875
N(1.1)	.000000	.000000	.000000
- (- <i>i</i> - <i>i</i>	7.351972	1.908267	6.759715
$C_{l_{\mu}}$	265177	3.804198	.499484
- 11	-8561.672085	-95004.636200	3.652575
	.593074	.668869	.076634
	8707.625432	15317.270050	.928264
	-146.231038	79694.444592	-3.828033
N(1,2)	14.235672	10.400895	17.891303
<u>_</u>	.957753	6.483311	319441
$C_{l_{\alpha}}$	1.258183	637080	2.425840
	404976.770760	199406.181992	-21.350260
	-8.145772	-4.472213	11.313079
	-408964.029690	-31963.187732	-23.973478
	3987.367830	-167451.070982	22.699002
N(2,1)	.000000	.000000	.000000
~	-1,905869	-1.614667	-1.979172
C_{m_k}	293633	-0.378487	.006105
11	158360.047503	10060.571210	-3.656111
	059842	.024373	004016
	-159903.225805	-1624.362035	219576
	1543.931881	-8436.118582	4.435226
N(2,2)	516839	178982	-4.220761
	-4.394286	2.247185	-1.072839
C_{m}	2.744002	-2.311694	.007098
- ///(4	217699.240354	21872.971701	7.958546
	.165112	038852	309693
	-219881,401614	-3648.306001	3,606871
	2184.097622	-18230.087698	-8.442002
OPTIMA	L		
LAG TE	RMS .27892661	.54004054	.30284145
	.04365735	.06187302	.00461111
	.27899139	.54678439	.02192557
	.28586712	.53878120	.35119036
ER	ROR .00139602	.00106544	.00236629

Appendix J : 3DOF Incompressible Unsteady Aerodynamics

Expressions for Lift force. Pitching Moment and Hinge Moment

$$\begin{split} P_{u} &= 1 - 2i \frac{V}{bu} (F + iG) \\ P_{\sigma} &= \frac{1}{2} - i \frac{V}{bu} [1 + 2(F + iG)] - 2\left(\frac{V}{bu}\right)^{2} (F + iG) \\ P_{\sigma} &= -\frac{L_{1}}{\tau} + i \frac{V}{bu} \frac{T_{1}}{\tau} - i \frac{V}{bu} \frac{T_{11}}{\tau} (F + iG) - 2\left(\frac{V}{bu}\right)^{2} \frac{T_{10}}{\pi} (F + iG) \\ M_{w} &= \frac{1}{2} \\ M_{\sigma} &= \frac{3}{8} - i \frac{V}{bu} \\ M_{,} &= -\frac{T_{1}}{\tau} - \left(c_{,3} + \frac{1}{2}\right) \frac{T_{1}}{\tau} + i \frac{V}{bu} \frac{T_{4} - \frac{2}{3}(\sqrt{1 - c_{3}^{2}})^{3}}{\pi} - \left(\frac{V}{bw}\right)^{2} \frac{T_{4} + T_{10}}{\pi} \\ T_{w} &= -\frac{T_{1}}{\tau} - i \frac{V}{bu} \frac{T_{12}}{\pi} (F + iG) \\ T_{\sigma} &= -\frac{1}{\tau} \left[T_{\tau} + \left(c_{,3} + \frac{1}{2}\right) T_{1}\right] + i \frac{V}{bu} \frac{\frac{2}{3}(\sqrt{1 - c_{3}^{2}})^{3} + 2T_{1} + T_{4}}{2\pi} - i \frac{V}{bw} \frac{T_{12}}{\pi} (F + iG) - \left(\frac{V}{bu}\right)^{2} \frac{I_{12}}{\tau} (F + iG) \\ T_{\sigma} &= -\frac{1}{\tau} \left[T_{\tau} + \left(c_{,3} + \frac{1}{2}\right) T_{1}\right] + i \frac{V}{bu} \frac{\frac{2}{3}(\sqrt{1 - c_{3}^{2}})^{3} + 2T_{1} + T_{4}}{2\pi} - i \frac{V}{bw} \frac{T_{12}}{\pi} (F + iG) - \left(\frac{V}{bu}\right)^{2} \frac{I_{12}}{\tau} (F + iG) \\ i G_{j} \end{split}$$

$$T_{\mathcal{A}} = -\frac{T_{2}}{\tau^{2}} + i \frac{1}{h_{u}} \frac{T_{4} T_{11}}{2\pi^{2}} - i \frac{1}{h_{u}} \frac{T_{11} T_{12}}{2\pi^{2}} (F + iG) - \left(\frac{1}{h_{u}}\right)^{2} \frac{T_{2} - T_{4} T_{10}}{\tau^{2}} - \left(\frac{1}{h_{u}}\right)^{2} \frac{T_{10} T_{12}}{\tau^{2}} (F + iG)$$

$$T_{1} = -\frac{1}{3}\sqrt{1 - c_{\beta}^{2}}(2 + c_{\beta}^{2}) + c_{\beta}cos^{-1}c_{\beta}$$

$$T_{2} = c_{\beta}(1 - c_{\beta}^{2}) - \sqrt{1 - c_{\beta}^{2}}(1 + c_{\beta}^{2})cos^{-1}c_{\beta} + c_{\beta}(cos^{-1}c_{\beta})^{2}$$

$$T_{3} = -\left(\frac{1}{8} + c_{\beta}^{2}\right)(cos^{-1}c_{\beta})^{2} + \frac{1}{4}c_{\beta}\sqrt{1 - c_{\beta}^{2}}cos^{-1}c_{\beta}(7 + 2c_{\beta}^{2}) - \frac{1}{8}(1 - c_{\beta}^{2})(5c_{\beta}^{2} + 4)$$

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$$T_{4} = -cos^{-1}c_{\beta} + c_{\beta}\sqrt{1 - c_{\beta}^{2}}$$

$$T_{5} = -(1 - c_{\beta}^{2}) - (cos^{-1}c_{\beta})^{2} + 2c_{\beta}\sqrt{1 - c_{\beta}^{2}}cos^{-1}c_{\beta}$$

$$T_{6} = T_{2}$$

$$T_{7} = -\left(\frac{1}{8} + c_{\beta}^{2}\right)cos^{-1}c_{\beta} + \frac{1}{8}c_{\beta}\sqrt{1 - c_{\beta}^{2}}(7 + 2c_{\beta}^{2})$$

$$T_{8} = -\frac{1}{3}\sqrt{1 - c_{\beta}^{2}}(2c_{\beta}^{2} + 1) + c_{\beta}cos^{-1}c_{\beta}$$

$$T_{9} = \frac{1}{2}\left[\frac{1}{3}(\sqrt{1 - c_{\beta}^{2}})^{3} + aT_{4}\right]$$

$$T_{10} = \sqrt{1 - c_{\beta}^{2}} + cos^{-1}c_{\beta}$$

$$T_{11} = cos^{-1}c_{\beta}(1 - 2c_{\beta}) + \sqrt{1 - c_{\beta}^{2}}(2 - c_{\beta})$$

$$T_{12} = \sqrt{1 - c_{\beta}^{2}}(2 + c_{\beta}) - cos^{-1}c_{\beta}(2c_{\beta} + 1)$$

$$T_{14} = \frac{1}{2}[-T_{7}^{2} - (c_{\beta} - a)T_{1}]$$

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Appendix K : Control Transformation Matrix

Various Airfoil Configurations (transformation)

For an autoil of specific geometric configuration, Nissim and Abel determined the optimal parameters for a PD transfer function control. For an airfoil flap having a length equal to 20% of the chord ($c_{\beta} = 0.6$, see figure 10), and having displacement measured at 30% chord (measured from leading edge, gives a = -0.4), they found the corresponding values

$$[t] = [t_{11}, t_{12}] = [0.0, -1.86]$$

 $[t^{*}] = [t_{11}^{*}, t_{12}^{*}] = [4.0, 3.20]$.

In this study an airfoil flap having a length equal to 20% of the chord ($c_3 = 0.6$, see figure 10) and having displacement measured at 25% chord (corresponding value of a = -0.5) was considered. A transformation was required to account for the difference in position at which the measurements were made (Nissim 1971; Nissim 1977)

The transformation matrix can be written as

$$[]_{s} = []_{opt} \begin{bmatrix} 1 & -x_{s} - 0.4 \\ 0 & 1 \end{bmatrix}$$

where $[\]_{opt}$ is the original matrix values Nissim and Abel obtained for the specific geometry of an airfoil flap having a length equal to 20% of the chord ($c_{\beta} = 0.6$, and having displacement measured at 30% chord (measured from leading edge, gives a = -0.4), namely

$$[t] = [t_{11}, t_{12}] = [0.0, -1.86]$$
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$$\begin{bmatrix} t^{-1} \end{bmatrix} = \begin{bmatrix} t_{11} & t_{12} \end{bmatrix} = \begin{bmatrix} 4.0 & 3.20 \end{bmatrix}$$

The value of x, is analogous to the quantity a, that is used in this thesis. Note when a value of a = -0.4 is used in the transformation matrix, the original matrix results are correctly obtained

In this study an airfoil flap having a length equal to 20% of the chord ($c_3 = 0.6$, see figure 10) and having displacement measured at 25% chord (corresponding value of a = -0.5) was considered. A transformation was completed to account for the difference in position of measurement by

$$\begin{bmatrix} t \end{bmatrix}_{s} = \begin{bmatrix} t \end{bmatrix}_{opt} \begin{bmatrix} 1 & -x_{s} - 0.4 \\ 0 & 1 \end{bmatrix}$$
$$\begin{bmatrix} t \end{bmatrix}_{s} = \begin{bmatrix} 0 & -1.86 \end{bmatrix}_{opt} \begin{bmatrix} 1 & -(-0.5) - 0.4 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 & -1.86 \end{bmatrix}$$

and

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Appendix L : Computer Program Listings

	PROGRAM	INPUT FILES	OUTPUT FILES	COMMENTS
1	UG.FOR	UGTHEO.DAT CONST.DAT	UGDATA.OUT	eigenvalue analysis flutter calculation COMP.BAT (FORTRAN) (incompressible)
2	PK.FOR	PKTHEO.DAT CONST1.DAT	PKDATA . OUT	eigenvalue analysis flutter calculation COMP.BAT (FORTRAN) (incompressible)
3	SIMPN.FOR	N11.DAT N12.DAT	LAG.OUT LAPLACE.DAT DATA.OUT	rational function approximation (incompressible)
4	SIMPTRN.FOR	N11.DAT N12.DAT N21.DAT N22.DAT	LAG.OUT LAPLACE.DAT DATA.OUT	rational function approximation (transonic)
5	LAPINCOM.FOR	LAPLACE.DAT CONST6.DAT	EIGEN.OUT	eigenvalue analysis flutter calculation COMP.BAT (FORTRAN) (incompressible)
6	LAPTRAN.FOR	LAPLACE.DAT CONST6.DAT	EIGEN.OUT	eigenvalue analysis flutter calculation COMP.BAT (FORTRAN) (transonic)
7	LAPPDOP.FOR	LAPLACE.DAT CONST6.DAT	OPTIMAL.OUT DATA.OUT	eigenvalue analysis flutter calculation COMP.BAT (FORTRAN) (incompressible)

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```
$storage:2
$floatcalls
Sdebua
C--
С
     U-q EIGENVALUE SOLUTION
C-
С
      IMPLICIT REAL*8 (A-Z)
      INTEGER ib, n, ia, ijob, iz, ier, infer, i, j
      COMPLEX*16 eval(2),A(2,2),B(2,2),EIGA(2),EIGB(2),Z(2,2),WK(3,6)
      REAL k(40), kc(40), F(40), G(40), u
      REAL abar(2,40), bbar(2,40), ndvel(2,40), ndw(2,40)
      REAL gdamp(2, 40), dratio(2, 40)
      INTEGER pt, pts
С
C Read in constants.
      OPEN (UNIT=1, FILE='CONST.DAT', STATUS='OLD')
           READ (1,*) ia, ib, iz, n, ijob
      CLOSE (UNIT=1)
C
C open required files
      OPEN (UNIT=2, FILE='CONST.OUT', STATUS='NEW')
      OPEN (UNIT=4, FILE='UGEIGEN.OUT', STATUS='NEW')
C
C required constants
              u=50.0
              wrat=0.2
               r=0.5
               ah=-0.5
С
              x = 0.25
              semic≈1.0
              wtor=1.0
              pts=21
C
C Theodorsen function C(k) = F + iG given by k, where k = kc/2
C thus if kc=0.2, k=0.1 therefore require C(k)=2(0.1)
      OPEN (UNIT=3, FILE='UGTHEO.DAT', STATUS='( つ)
          do 50 pt=1,pts
             READ (3, *) k(pt), F(pt), G(pt)
              kc(pt) = 2.0 * k(pt)
 50
          continue
      CLOSE (UNIT=3)
C
C formation of COMPLEX matrices [A] and [B] to
C solve eigenvalue problem given by [A](x) = lamba [B](x)
          do 100 pt=1,pts
              ar11=0.25*u*kc(pt)**2+0.25*kc(pt)**2+kc(pt)*G(pt)
              aill=-kc(pt)*F(pt)
              ar12=0.25*x*u*kc(pt) **2-0.25*ah*kc(pt) **2-2*F(pt) +
     1
              (0.5-ah) *kc(pt) *G(pt)
             ai12=-0.5*kc(pt)-2*G(pt)-(0.5-ah)*kc(pt)*F(pt)
             ar21=0.25*x*u*kc(pt)**2-0.25*ah*kc(pt)**2-(ah+0.5)
     1
              *kc(pt) *G(pt)
```

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ai21=(ah+0.5)*kc(pt)*F(pt)
             ar22=0.25*r**2*u*kc(pt)**2+(0.125+ah**2)*0.25*kc(pt)**2
             +2*(ah+0.5)*F(pt)-(ah+0.5)*(0.5-ah)*kc(pt)*G(pt)
    1
             ai22=-(0.5-ah)*0.5*kc(pt)+2*(ah+0.5)*G(pt)+(ah+0.5)
    1
             *(0.5-ah)*kc(pt)*F(pt)
                a(1,1)=dcmplx(arl1,ail1)
                a(1,2) = dcmplx(ar12,ai12)
                a(2,1)=dcmplx(ar21,ai21)
                a(2,2) = dcmplx(ar22,ai22)
             brll=wrat**2
             bill=0.0
             br12=0.0
             bi12=0.0
             br21=0.0
             bi21=0.0
             br22=r**2
             bi22=0.0
                b(1,1)=dcmplx(brl1,bil1)
                b(1,2) = dcmplx(br12,bi12)
                b(2,1)=dcmplx(br21,bi21)
                b(2,2) = dcmplx(br22,bi22)
С
C subrountine for eigenvalue calculation
C abar : real part of eval()
C bbar : imaginary part of eval( )
      CALL EIGZC(A, IA, B, IB, N, IJOB, EIGA, EIGB, Z, IZ, WK, INFER, IER)
          DO 5 I=1, n
             EVAL(I) = EIGA(I) / EIGB(I)
 5
             write(2,19) k(pt),F(pt),G(pt)
 19
             format('
                        k ',f5.3,'
                                           F + iG', 2f12.7
             write(2,20) arl1,ail1,arl2,ail2,ar21,ai21
             write(2,20) ar22,ai22,br11,br22
             format(10f12.7)
 20
             write(2,21)wk(1,1)
 21
             format(' perfomance index is', 2f10.3)
             write(2,22)eval(1)
                 abar(1,pt)=dreal(eval(1))
                 bbar(1,pt)=dimag(eval(1))
             write(2,22)eval(2)
                 abar(2,pt) = dreal(eval(2))
                 bbar(2,pt)=dimag(eval(2))
             write(4,22)k(pt),eval(1)
             write(4,22)k(pt),eval(2)
          do 6 i=1,n
 6
             write(2,22) (z(i,j),j=1,n)
 22
             format(8f12.5)
 100
          continue
      CLOSE (UNIT=3)
С
C calculations and data table
C ndvel
                                                 vel/(b*wtor)
           : nd velocity
C ndw
                                                   w2/wtor
            : nd frequency
C gdamp
            : structural damping coefficient
                                                      g
C dratio
           : damping ratio
                                                    -g/2
```

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```
```
OPEN (UNIT=5, FILE='UGDATA.OUT', STATUS='NEW')
      OPEN (UNIT=6, FILE='UGDAMP.OUT', STATUS='NEW')
      OPEN (UNIT=7, FILE='UGFREQ.OUT', STATUS='NEW')
             write(5,32)
             write(5,121)
             write(5,121)
             write(5,33)
             write(5,121)
             write(5,34) u,r,ah,wrat
             write(5,121)
             write(5,120)
             write(5,121)
 32
             format(' UG METHOD ')
 33
             format('
                                     r(alp)
                                                ah
                                                       wben/wtor')
                        u
 34
             format(f7.1,4f10.1)
120
                                    U/b*wtor
             format('
                        k
                                                   w/wtor
                                                                  g
                                                                      ٠,
     1'
          damping ratio ')
121
                         1)
             format('
          do 140 pt=1,pts
          do 130 i=1,2
             ndvel(i,pt)=(u/abar(i,pt))**0.5
             ndw(i,pt) = k(pt) * (u/abar(i,pt)) **0.5
             gdamp(i,pt)=bbar(i,pt)/abar(i,pt)
             dratio(i,pt)=-gdamp(i,pt)/2
C printing data to to be graphed to file 'UGDATA.OUT'
             write(5,125) k(pt),ndvel(i,pt),
     1
             ndw(i,pt),gdamp(i,pt),dratio(i,pt)
             write(6,125) ndvel(i,pt),dratio(i,pt)
             write(7,125) ndvel(i,pt),ndw(i,pt)
             format(f6.3,',',f13.5,',',f12.5,',',f9.5,','
 125
             ,f9.5,',',f12.5,',',f12.5)
     1
 130
          continue
 140
          continue
           CLOSE (UNIT=7)
           CLOSE (UNIT=6)
           CLOSE (UNIT=5)
      END
```

INPUT FILE : UG	HEO.DA	Г
-----------------	--------	---

; 4

THEODORSEN FUNCTION C(k)

10.00	0.5006	-0.0124
6.000	0.5017	-0.0206
4.000	0.5037	-0.0305
3.000	0.5063	-0.0400
2.000	0.5129	-0.0577
1.500	0.5210	-0.0736
1.200	0.5300	-0.0877
1.000	0.5394	-0.1003
0.800	0.5541	-0.1165
0.600	0.5788	-0.1378
0.500	0.5979	-0.1507
0.400	0.6250	-0.1650
0.300	0.6550	-0.1793
0.200	0.7276	-0.1886
0.150	0.7737	-0.1857
0.130	0.7954	-0.1820
0.110	0.8188	-0.1766
0.100	0.8320	-0.1723
0.050	0.9090	-0.1305
0.025	0.9545	-0.0872
0.010	0.9824	-0.0482
k	F	iG
INPUT	FILE :	CONST.DAT

2 2 2 2 2

*

A - 50

UG METHOD

-

ŗ

u	r(alp)	ah wh	en/wtor		
50.0	.5	5	.2		
k	U/b*wtor	w/wtor	g g	damping	ratio
10.000,	.01970,	.19698,	00209,	.00104,	
10.000,	.11493,	1.14932,	00569,	.00284,	
6.000.	.03283,	. 19699,	00348,	.00174,	
6.000,	.19150,	1.14900,	00948,	.00474,	
4.000,	.04925,	.19701,	00524,	.00262,	
4.000,	.28709,	1.14837,	01423,	.00712,	
3.000,	.06568,	.19704	00703,	.00351,	
3.000,	. 38249,	1.14748,	01901,	.00950,	
2.000.	.09856,	.19711,	01067,	.00533,	
2.000,	.57247,	1.14494,	02861,	.01431,	
1.500.	.13147.	.19721.	01444,	.00722,	
1.500.	.76090	1.14135.	03831.	.01916,	
1.200.	.16444.	.19732.	01834.	.00917,	
1.200.	.94726.	1.13672.	04810,	.02405.	
1.000.	.19746.	.19746.	02238.	.01119,	
1.000.	1.13102.	1.13102	05796.	.02898,	
.800.	.24711.	. 19769 .	02871.	.01436.	
.800.	1.40063.	1,12050	07275.	.03638.	
.600	. 33024.	. 19814 .	03995.	.01998.	
.600.	1.82975.	1.09785	09686.	.04843.	
.500.	.39712.	. 19856	04954.	.02477.	
.500.	2,15096.	1.07548	11494.	.05747.	
.400	.49815.	. 19926.	06485.	.03243.	
.400.	2.58976.	1.03590	13857.	.06929.	
.300.	. 66874 .	.20062	09274.	.04637.	
.300.	3.19652.	.95896.	16476.	.08238.	
.200	1.02048	.20410	15760.	.07880.	
.200	3.97729.	.79546	16106.	.08053.	
.150.	1.39164	.20875.	23812.	.11906.	
.150.	4.35545.	.65332	09675.	.04837.	
.130	1.63333	21233	- 29866	14933	
.130	A A7879	58224	- 03713	.01856	
.110	1.98367	21820	- 40100	.20050	
. 110	4 59055	50496	06404	03202	
100	2 22456	22246	- 48305	24197	
100,	A 66272	46627	14335	- 07167	
050	A 0707A	24954	-1 75100	97500	
.050,	7 10200	144034) 266655	1 11016	- 5550g	
.030,	10 27783	.JJJ70J 2604 E	-4 24223	2 12111	
025	13 00073	· 40740		4+±4±±+, _1 32201	
.025,	13.30V/2, 37 /1033	134702, 17810	_10 09404/	-1.JEZUI,	
.010,	21.41723, 38 89513	14/919/ 14/95	-16.60464; 6 70404	-3 30040	
AUTU.					

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```
$storage:2
Sfloatcalls
$debug
C----
     p-k EIGENVALUE SOLUTION
С
C-
    _____
С
      IMPLICIT REAL*8 (A-Z)
      INTEGER ib, n, ia, ijob, iz, ier, infer, i, j
      COMPLEX + 16 eval(4), A(4,4), B(4,4), EIGA(4), EIGB(4), Z(4,4), WK(6,12)
      REAL k(40), kc(40), F(40), G(40), u, vel
      REAL abar(4), bbar(4), ndvel(4,40), ndw(4,40)
      REAL gdamp(4,40), dratio(4,40), kg, fg,gg, step
      INTEGER temp, it, pt, pts, test, soln1, isoln, check, soln(2)
С
C Read in constants.
      OPEN (UNIT=1, FILE='CONST1.DAT', STATUS='OLD')
           READ (1, \star) ia, ib, iz, n, ijob
      CLOSE (UNIT=1)
С
C open required files
      OPEN (UNIT=2, FILE='CONST.OUT', STATUS='NEW')
      OPEN (UNIT=4, FILE='PKDATA.OUT', STATUS='NEW')
      OPEN (UNIT=5, FILE='PKDAMP.OUT', STATUS='NEW')
      OPEN (UNIT=6, FILE='PKFREQ.OUT', STATUS='NEW')
С
C required constants
               u=50.0
               wrat=0.2
               r=0.5
               ah = -0.5
C
               x=0.25
               semic=1.0
               wtor=1.0
               wben=wrat*wtor
               pts=41
С
            write(4,32)
            write(4,35)
            write(4,35)
            write(4,33)
            write(4,35)
            write(4,34) u,r,ah,wrat
            write(4,35)
            write(4,36)
            write(4,35)
            format('
 32
                       PK METHOD ')
            format('
 33
                                                ah
                                                        wben/wtor')
                                    r(alp)
                        u
            format(f7.1,4f10.1)
 34
            format(' ')
 35
            format('
                                        b (damping)
 36
                        k
                             ND VEL
                                                           w/wtor ')
C
C Theodorsen function C(k) = F + iG given by k, where k = kc/2
```

A - 52

W. 1

÷.

```
C thus if kc=0.2, k=0.1 therefore require C(k)=C(0.1)
      OPEN (UNIT=3, FILE='PKTHEC.DAT', STATUS='OLD')
          do 50 pt=1,pts
             READ (3, \star) k(pt), F(pt), G(pt)
             kc(pt) = 2.0 * k(pt)
 50
          continue
      CLOSE (UNIT=3)
С
C
C iterate velocity
             vel=0.0
             step =0.5
 60
             vel=vel+step
             if(vel/(semic*wtor).gt.5.0) goto 120
          do 95 test=1,2
             check=0
             count=0
             kg=((wtor+wben)/2)*semic/vel
С
C interpolation of reduced frequency, F and G
 51
             tempw=vel*kg/semic
             it=1
 52
             it=it+1
             if(kg.lt.0.01) kg=0.01
             if(k(it).lt.kg) goto 53
             goto 52
             Fg=f(it-1)+(f(it)-f(it-1))*(kg-k(it-1))/(k(it)-k(it-1))
 53
             Gg=g(it-1)+(g(it)-g(it-1))*(kg-k(it-1))/(k(it)-k(it-1))
 54
             format(' ',8f12.5)
С
C formation of COMPLEX matrices [A] and [B] to
C solve eigenvalue problem given by [A]{x} = lambda [B]{x}
             ar11=0.0
             ai11=0.0
             ar12=0.0
             ai12=0.0
             ar13=wrat**2*wtor**2
             ai13=0.0
             ar14=(2/u)*(vel/semic)**2*Fg
             ai14=(2/u) *(vel/semic) **2*Gg
             ar21=0.0
             ai21=0.0
             ar22=0.0
             ai22=0.0
             ar23=0.0
             ai23=0.0
             ar24=r**2*wtor**2-2*(ah+0.5)*(Fg/u)*(vel/semic)**2
             ai24=-2*(ah+0.5)*(Gg/u)*(vel/semic)**2
             ar31=-1.0
             ai31=0.0
             ar32=0.0
             ai32=0.0
             ar33=0.0
             ai33=0.0
```

```
.4 – 53
```

```
----
```

```
ar34=0.0
ai34=0.0
ar41=0.0
ai41=0.0
ar42=-1.0
ai42=0.0
ar43=0.0
ai43=0.0
ar44=0.0
ai44=0.0
   a(1,1) = dcmplx(arl1,ail1)
   a(1,2) = dcmplx(ar12,ai12)
   a(1,3) = dcmplx(ar13, ai13)
   a(1,4) = dcmplx(arl4,ail4)
   a(2,1) = dcmplx(ar21,ai21)
   a(2,2) = dcmplx(ar22,ai22)
   a(2,3) = dcmplx(ar23,ai23)
   a(2,4) = dcmplx(ar24,ai24)
   a(3,1) = dcmplx(ar31,ai31)
   a(3,2) = dcmplx(ar32,ai32)
   a(3,3) = dcmplx(ar33,ai33)
   a(3,4) = dcmplx(ar34,ai34)
   a(4,1) = dcmplx(ar41,ai41)
   a(4,2) = dcmplx(ar42,ai42)
   a(4,3) = dcmplx(ar43,ai43)
   a(4,4) = dcmplx(ar44,ai44)
br11=1+1/u
bi11=0.0
br12=x-ah/u
bi12=0.0
br13=2*(Fg/u)*(vel/semic)
bi13=2*(Gg/u)*(vel/semic)
br14=((1/u)+2*(Fg/u)*(0.5-ah))*(vel/semic)
bil4=(2*(Gg/u)*(0.5-ah))*(vel/semic)
br21=x-ah/u
bi21=0.0
br22=r**2+(0.125+ah**2)/u
bi22=0.0
br23=-2*(ah+0.5)*(Fg/u)*(vel/semic)
bi23=-2*(ah+0.5)*(Gg/u)*(vel/semic)
br24=((0.5-ah)/u-2*(ah+0.5)*(Fg/u)*(0.5-ah))*(vel/semic)
bi24= (-2*(ah+0.5)*(Gg/u)*(0.5-ah))*(vel/semic)
br31=0.0
bi31=0.0
br32=0.0
bi32=0.0
br33=1.0
bi33=0.0
br34=0.0
bi34=0.0
br41=0.0
bi41=0.0
br42=0.0
bi42=0.0
```

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```
br43=0.0
             bi43=0.0
             br44=1.0
             bi44=0.0
                 b(1,1) = dcmplx(-brll,-bill)
                 b(1,2) = dcmplx(-br12,-bi12)
                 b(1,3) = dcmplx(-br13,-bi13)
                 b(1,4) = dcmplx(-br14,-bi14)
                 b(2,1) = dcmplx(-br21,-bi21)
                 b(2,2) = dcmplx(-br22,-bi22)
                 b(2,3) = dcmplx(-br23,-bi23)
                 b(2,4) = dcmplx(-br24,-bi24)
                 b(3,1) = dcmplx(-br31,-bi31)
                 b(3,2) = dcmplx(-br32,-bi32)
                 b(3,3) = dcmplx(-br33,-bi33)
                 b(3,4) = dcmplx(-br34,-bi34)
                 b(4,1) = dcmplx(-br41,-bi41)
                 b(4,2) = dcmplx(-br42,-bi42)
                 b(4,3) = dcmplx(-br43,-bi43)
                 b(4,4) = dcmplx(-br44,-bi44)
С
C subrountine for eigenvalue calculation
C abar : real part of eval()
C bbar : imaginary part of eval()
      CALL EIGZC(A, IA, B, IB, N, IJOB, EIGA, EIGB, Z, IZ, WK, INFER, IER)
С
          do 5 I=1,n
 5
              EVAL(I)=EIGA(I)/EIGB(I)
21
               format(' perfomance index is', 2f10.3)
                write(2,22)eval(1)
С
                 abar(1) = dreal(eval(1))
                 bbar(1) = dimag(eval(1))
С
                write(2,22)eval(2)
                 abar(2) = dreal(eval(2))
                 bbar(2) = dimag(eval(2))
С
                write(2,22)eval(3)
                 abar(3) = dreal(eval(3))
                 bbar(3) = dimag(eval(3))
С
                write(2,22)eval(4)
                 abar(4)=dreal(eval(4))
                 bbar(4) = dimag(eval(4))
C
C obtaining eigenvalues with positive imaginary parts for analysis
               j=0
           do 14 isoln=1,4
               if (bbar(isoln).gt.0.0) goto 12
               goto 14
 12
               j=j+1
               soln(j)=isoln
 14
           continue
               if(check.eq.0) goto 23
               if(test.eq.l.and.abs(tempw-bbar(soln(1))).lt.abs(
     1
               tempw-bbar(soln(2))) goto 23
               if(test.eq.2.and.abs(tempw-bbar(soln(2))).lt.abs(
```

```
A - 55
```

```
1
              tempw-bbar(soln(1))) goto 23
              if(test.eq.2) goto 19
              temp=soln(1)
              soln(1) = soln(2)
              soln(2) ≈temp
              goto 23
 19
              temp=soln(2)
              soln(2) = soln(1)
              soln(1)=temp
С
C convergence check
 23
              check=1
              w=dimag(eval(soln(test)))
               if (abs(w-tempw).lt.0.001) goto 94
              kg=w*semic/vel
              goto 51
С
C writing final data to file
              write(4,27)kg,vel/(semic*wtor),dreal(eval(soln(test))),
 94
     1
               dimag(eval(soln(test)))/wtor
               write(5,27) vel/(semic*wtor),-dreal(eval(soln(test)))/
               ((dreal(eval(soln(test))))**2+(dimag(eval(soln(test))))
     1
     1
               **2)**0.5
               write(6,27) vel/(semic*wtor),dimag(eval(soln(test)))/wtor
               format(f6.4,',',f7.4,',',f12.5,',<sup>1</sup>,f12.5)
 27
               format(8f12.5)
 22
 95
         continue
С
C changing velocity step size
               if(vel/(semic*wtor).ge.5.0) step=0.1
С
                if(vel/(semic*wtor).ge.4.8) step=0.05
С
                if(vel/(semic*wtor).ge.5.0) step=0.5
                if(vel/(semic*wtor).ge.6.0) goto 120
С
 100
               goto 60
       CLOSE (UNIT=4)
 120
       CLOSE (UNIT=2)
 200
            END
 300
```

INPUT FILE : PKTHEO.DAT

THEODORSEN FUNCTION C(k)

10.00 6.000 4.000	0.5006 0.5017 0.5037	-0.0124 -0.0206 -0.0305
3.000	0.5063	-0.0400
2.000	0.5129	-0.0577
1.500	0.5210	-0.0736
1.200	0.5300	-0.0877
1.000	0.5394	-0.1003
0.990	0.5400	-0.1010
0.800	0.54/4	-0.1095 -0.1165
0.770	0.5570	-0.1193
0.660	0.5699	-0.1308
0.600	0.5788	-0.1378
0.560	0.5857	-0.1428
0.550	0.5876	-0.1441
0.500	0.5979	-0.1507
0.440	0.6130	-0.1592
0.400	0.6250	-0.1650
0.340	0.6469	-0.1738
0.330	0.6512	-0.1752
0.325	0.6535	-0.1759
0.320	0.6558	-0.1766
0.315	0.6581	-0.1773
0.310	0.6604	-0.179
0.305	0.002/	-0.1788
0.240	0.6989	-0.1862
0.220	0.7125	-0.1877
0.200	0.7276	-0.1886
0.160	0.7628	-0.1876
0.120	0.8063	-0.1801
0.110	0.8188	-0.1766
0.100	0.8320	-0.1723
0.080	0.8604	-0.1604
0.060	0.8920	-0.1426
0.050	0.9090	-0.1305
0.040	0.9267	-0.1160
0.025	0.9545	-0.0872
0.010	0.9824	-0.0482
0.000	T.0000	-0.0000
k	F	iG

INPUT FILE : CONST1.DAT

4 4 4 4 2

ŧ

OUTPUT FILE : PKDATA.OUT

PK METHOD

I

u	r(alı	o) ah	wben/wtor
50.0	• 5	5	. 2
k	ND VEL E	(damping)	w/wtor
.3965, 2.2909, .2034,	.5000, .5000, 1.0000,	00640, 01436, 01487,	.19917 1.14593 .20335
1.1348, .1394, .7434,	1.0000, 1.5000, 1.5000,	02958, 02492, 04621,	1.13484 .20916 1.11514
.1087, .5425, .0916,	2.0000, 2.0000, 2.5000,	03677, 06450, 05169.	.21735 1.08521 22893
.4169, .3275, .0820.	2.5000, 3.0000, 3.0000	08437, 10547, 07180	1.04258 .98336
.2570, .0781,	3.5000, 3.5000, 4.0000	12613,10247,10247,12092	.24614 .89945 .27359
.0809,	4.0000 , 4.5000 , 4.5000 ,	16167, 01093,	. 76981 . 32367 . 55238
.1020,	5.0000, 5.0000,	33///, .12384, 60286,	.33681 .50991 .14083

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C-----------SEQUENTIAL SIMPLEX PROGRAM (MINIMIZATION PROBLEM SOLVING) С С RATIONAL FUNCTION APPROXIMATION - INCOMPRESSIBLE REGIME C-С С OPERATION С REQUIRE: 1) INPUT OF ICOUNT, N, REQMIN, START(I), STEP(I) С 2) USER SPECIFIED FUNCTION SUBPROGRAM (DOUBLE PRECISION С FUNCTION FN(X) - MINIMIZATION FUNCTION) С С DOUBLE PRECISION START(20), STEP(20), XMIN(20), 1XSEC(20), YNEWLO, YSEC, REQMIN С OPEN(UNIT=6, FILE='DATA.OUT') С С С * С * ICOUNT, N, REQMIN, START(N), STEP(N) С С ********** С C LET N=4 AT ALL TIMES, SIMPLY SET START(I)=0.0 AND STEP(I)=0.0 С ICOUNT=260 N=4REQMIN=0.000000000001 START(1) = 0.015START(2) = 0.084464515START(3) = 0.25635START(4) = 0.76550587STEP(1) = 0.0105STEP(2) = 0.01STEP(3) = 0.05STEP(4) = 0.1С С DO 60 I=1,N XMIN(I) = 0.D0XSEC(I) = 0.D060 CONTINUE YNEWLO=0.D0 YSEC=0.DO С С CALL NELDER-MEAD SUBROUTINE С CALL NELMIN (N, START, XMIN, XSEC, YNEWLO, YSEC, **1REQMIN**, STEP, ICOUNT) С OUTPUT FROM PROGRAM С С WRITE(6,64) WRITE(6,65) ICOUNT WRITE(6,75)

```
WRITE(6,77)
      DO 79 I=1,N
 79
      WRITE(6,80) I,XSEC(I),XMIN(I)
      WRITE(6,82)
      WRITE(6,83)
      WRITE(6,84) YSEC, YNEWLO
С
 64
      FORMAT(6X,1H //,1H ,42H SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION)
 65
      FORMAT(1H //,1H ,15,12H TRIALS USED/)
 75
      FORMAT(1H ,21X,9HESTIMATES/)
 77
      FORMAT(6X,1H,9HPARAMETER,7X,12HNEXT-TO-BEST,8X,
     14HBEST/)
      FORMAT(1H , 15, 2F20.7)
 80
      FORMAT (6X, 1H //, 1H ,6X, 15HFUNCTION VALUES/)
 82
 83
      FORMAT(6X,1H,5X,13H NEXT-TO-BEST,8X,4HBEST/)
 84
      FORMAT(6X,1H,2F15.9)
      STOP
      END
C
С
              NELDER-MEAD SUBROUTINE
С
      SUBROUTINE NELMIN (N, START, XMIN, XSEC, YNEWLO, YSEC,
     1REQMIN, STEP, ICOUNT)
      DOUBLE PRECISION START(N), STEP(N), XMIN(N),
     1XSEC(N), YNEWLO, YSEC, REQMIN, P(20, 21), PSTAR(20),
     2P2STAR(20), PBAR(20), Y(20), DN, Z, YLO, RCOEFF,
     3YSTAR, ECOEFF, Y2STAR, CCOEFF, FN, DABIT, DCHK,
     4COORD1, COORD2
      DATA RCOEFF/1.0D0/, ECOEFF/2.0D0/, CCOEFF/0.5D0/
      KCOUNT=ICOUNT
      ICOUNT=0
С
      IF (REQMIN.LE.O.DO) ICOUNT=ICOUNT-1
      IF(N.LE.0) ICOUNT=ICOUNT-10
      IF(N.GT.20) ICOUNT=ICOUNT-10
      IF(ICOUNT.LT.0) RETURN
С
      DABIT=2.04607D-35
      BIGNUM=1.0D38
      KONVGE=5
      XN=FLOAT(N)
      DN=DFLOAT(N)
      NN=N+1
С
С
      CONSTRUCTION OF INITIAL SIMPLEX
С
 1001 DO 1 I=1,N
    1 P(I,NN) = START(I)
      Y(NN) = FN(START)
      ICOUNT=ICOUNT+1
      DO 2 J=1,N
      DCHK=START(J)
      START(J) = DCHK + STEP(J)
      DO 3 I=1,N
```

```
A - 60
```

```
3 P(I,J) = START(I)
      Y(J) = FN(START)
      ICOUNT=ICOUNT+1
    2 START(J)=DCHK
С
С
       SIMPLEX CONSTRUCTION COMPLETE
С
       FIND HIGHEST AND LOWEST Y VALUES
С
С
       YNEWLO (Y(IHI)) INDICATES THE VERTEX OF
С
       THE SIMPLEX TO BE REPLACED
С
 1000 \text{ YLO}=Y(1)
      YNEWLO=YLO
      ILO=1
      IHI=1
      DO 5 I=2,NN
      IF(Y(I).GE.YLO) GOTO 4
      YLO=Y(I)
      ILO=I
    4 IF(Y(I).LE.YNEWLO) GOTO 5
      YNEWLO=Y(I)
      IHI=I
    5 CONTINUE
С
С
        PERFORM CONVERGENCE CHECKS ON FUNCTION
С
      DCHK=(YNEWLO+DABIT)/(YLO+DABIT)-1.DO
      IF (DABS (DCHK). LT. REQMIN) GOTO 900
С
      KONVGE=KONVGE-1
      IF(KONVGE.NE.0) GOTO 2020
      KONVGE=5
С
С
        CHECK CONVERGENCE OF COORDINATES ONLY
С
        EVERY 5 SIMPLEXES
С
      DO 2015 I=1,N
      COORD1 = P(I, 1)
      COORD2=COORD1
      DO 2010 J=2,NN
      IF(P(I,J).GE.COORD1) GOTO 2005
      COORD1=P(I,J)
 2005 IF(P(I,J).LE.COORD2) GOTO 2010
      COORD2=P(I,J)
 2010 CONTINUE
      DCHK=(COORD2+DABIT)/(COORD1+DABIT)-1.D0
      IF (DABS (DCHK).GT.REQMIN) GOTO 2020
 2015 CONTINUE
      GOTO 900
 2020 IF(ICOUNT.GE.KCOUNT) GOTO 900
С
С
         CALCULATE PBAR, THE CENTROID OF THE
С
         SIMPLEX VERTICES EXCEPTING THAT WITH
С
         Y VALUE YNEWLO
```

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С
      DO 7 I=1,N
      Z=0.0D0
      DO 6 J=1, NN
    6 Z=Z+P(I,J)
      Z=Z-P(I,IHI)
    7 PBAR(I) = Z / DN
С
С
         REFLECTION THROUGH THE CENTROID
С
      DO 8 I=1,N
    8 PSTAR(I) = (1.0D0+RC0EFF) *PBAR(I)-RC0EFF*P(I,IHI)
      YSTAR=FN (PSTAR)
      ICOUNT=ICOUNT+1
      IF (YSTAR.GE.YLO) GOTO 12
      IF (ICOUNT.GE.KCOUNT) GOTO 19
С
С
         SUCCESSFUL REFLECTION, SO EXTENSION
С
      DO 9 I=1,N
    9 P2STAR(I) = ECOEFF*PSTAR(I)+(1.0D0-ECOEFF)*PBAR(I)
      Y2STAR=FN (P2STAR)
      1COUNT=ICOUNT+1
С
С
         RETAIN EXTENSION OR CONTRACTION
С
      IF(Y2STAR.GE.YSTAR) GOTO 19
   10 DO 11 I=1,N
   11 P(I, IHI) = P2STAR(I)
      Y(IHI)=Y2STAR
      GOTO 1000
С
С
         NO EXTENSION
С
   12 L=0
      DO 13 I=1,NN
      IF(Y(I).GT.YSTAR) L=L+1
   13 CONTINUE
      IF(L.GT.1) GOTO 19
      IF(L.EQ.O) GOTO 15
С
С
          CONTRACTION ON THE REFLECTION SIDE OF THE
С
         CENTROID
С
      DO 14 I=1,N
   14 P(I, IHI) = PSTAR(I)
      Y(IHI)=YSTAR
С
С
          CONTRACTION ON THE Y(IHI) SIDE OF THE CENTROID
С
   15 IF (ICOUNT.GE.KCOUNT) GOTO 900
      DO 16 I=1,N
   16 P2STAR(I) = CCOEFF*P(I, IHI) + (1.0D0 - CCUEFF)*PBAR(I)
      Y2STAR=FN (P2STAR)
```

```
ICOUNT=ICOUNT+1
      IF(Y2STAR.LT.Y(IHI)) GOTO 10
С
C
         CONTRACT THE WHOLE SIMPLEX
C
      DO 18 J=1, NN
      DO 17 I=1,N
      P(I,J) = (P(I,J) + P(I,ILO)) * 0.5D0
   17 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   18 CONTINUE
      ICOUNT=ICOUNT+NN
      IF(ICOUNT.LT.KCOUNT) GOTO 1000
      GOTO 900
С
С
         RETAIN REFLECTION
С
С
    19 CONTINUE
   19 DO 20 I=1,N
   20 P(I, IHI) = PSTAR(I)
      Y(IHI)=YSTAR
      GOTO 1000
С
С
         SELECT THE TWO BEST FUNCTION VALUES (YNEWLO
С
         AND YSEC) AND THEIR COORDS. (XMIN AND XSEC)
С
  900 DO 23 J=1,NN
      DO 22 I=1,N
   22 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   23 CONTINUE
      YNEWLO=BIGNUM
      DO 24 J=1,NN
      IF(Y(J).GE.YNEWLO) GOTO 24
      YNEWLO=Y(J)
      IBEST=J
   24 CONTINUE
      Y(IBEST)=BIGNUM
      YSEC=BIGNUM
      DO 25 J=1,NN
      IF(Y(J).GE.YSEC) GOTO 25
      YSEC=Y(J)
      ISEC=J
   25 CONTINUE
      DO 26 I=1,N
      XMIN(I) = P(I, IBEST)
      XSEC(I) = P(I, ISEC)
   26 CONTINUE
      RETURN
      END
С
С
     APPROXIMATING FUNCTION : PADE APPROXIMATE WITH LAG TERMS
С
       DOUBLE PRECISION FUNCTION FN(lag)
```

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```
IMPLICIT REAL*8 (A-Z)
       COMPLEX*16 ckquad, cklag, ckbar(40,2,2), p
       DOUBLE PRECISION LAG(4)
       REAL*8 w(40), k(40, 2, 2), F(40, 2, 2), G(40, 2, 2), a(9, 18), b(9, 1), x(9, 1)
       REAL*8 mn(40)
       INTEGER i,m, in, nst, nf, j, prt, totlag, c, r, count, icount
       INTEGER D, COL, ROW
С
С
         OPEN (UNIT=1, FILE='LAG.OUT')
         OPEN (UNIT=10, FILE='LAPLACE.DAT')
С
С
С
       С
       *
С
       *
           OPEN REQUIRED INPUT TABULAR DATA FILES
                                                       *
С
       *
           MINIMUM ERROR DESIRED FOR TABLE
С
           TOTAL # OF ITERATIONS
       *
                                  ICOUNT
С
       *
           INPUT NUMBER OF LAG TERMS REQUIRED
С
       *
           # OF POINTS (m) IN DATA FILES
С
       *
С
       *****
С
          icount=260+5
          minerr=0.0000167
          totlag=4
          m=40
С
          count=count+1
          if(count.gt.1) goto 88
С
С
     FILES CONTAIN
                     k(i,r,c), F(i,r,c), G(i,r,c)
С
     k(i) REDUCED FREQUENCY
С
     F(i) REAL
С
     G(i) IMAGINARY
С
C APPROXIMATING C(k) AND C(k)*ik
         OPEN (UNIT=2, FILE='N11.DAT', STATUS='OLD')
         OPEN (UNIT=4, FILE='N12.DAT', STATUS='OLD')
         do 150 i=1,m
             READ(2,*) k(i,1,1), F(i,1,1), G(i,1,1)
             READ(4,*) k(i,1,2),F(i,1,2),G(i,1,2)
 150
         continue
         CLOSE (UNIT=4)
         CLOSE(UNIT=2)
С
С
    REQUIRED CONSTANTS
С
             matrix [A] ( size
                                in x in )
 88
             nst=4
             in=nst+(totlag-1)
             nf=2*in
С
C NOTE: lag(1) \iff lag(2) \iff lag(3) \iff lag(4)
С
        lag(1), lag(2), lag(3), lag(4) must be all non-negative
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C BOUNDARY (CONSTRAINT) CONDITIONS
           do 156 i=1,totlag
              if(lag(i).lt.0.0) lag(i)=i*0.0001
           do 157 j=1, totlag-1
              if(i.eq.j) goto 157
              if(lag(i).eq.lag(j)) goto 163
              goto 157
              print 164, j
 163
               format (15, ' EQUAL LAG TERMS REQUIRES IMMEDIATE STOP')
 164
              stop
           continue
 157
 156
           continue
С
              FN=0.0
С
           do 750 c=1,2
               r=1
               err=0.0
С
     RESETTING MATRIX ELEMENT VALUES TO ZERO
 390
           do 410 i=1, in
           do 400 j=1,nf
               a(i,j) = 0.0
 400
           continue
               b(i,1) = 0.0
               x(i,1) = 0.0
 410
           continue
С
С
      CREATING MATRICES [A] AND [B] : SOLVE
                                                     [A] \{x\} = [B]
           do 175 i=1,m
               bk1=k(i,r,c)**2/(k(i,r,c)**2+lag(1)**2)
               if(totlag.eq.1) goto 415
               bk2=k(i,r,c)**2/(k(i,r,c)**2+lag(2)**2)
               if(totlag.eq.2) goto 415
               bk_3=k(i,r,c)**2/(k(i,r,c)**2+lag(3)**2)
               if(totlag.eq.3) goto 415
               bk4=k(i,r,c)**2/(k(i,r,c)**2+lag(4)**2)
 415
               mn(i) = F(i, r, c) * * 2 + G(i, r, c) * * 2
               if(mn(i).lt.1.0) mn(i)=1.0
               w(i) = 1.0
               w(i) = w(i) / mn(i)
               a(1,1) = a(1,1) + w(i) * (1.0)
               a(1,2) = a(1,2) + w(i) * (0.0)
               a(1,3) = a(1,3) + w(i) * (-k(i,r,c) * * 2)
               a(1,4) = a(1,4) + w(1) * (bk1)
               a(1,5) = a(1,5) + w(i) * (bk2)
               a(1,6) = a(1,6) + w(i) * (bk3)
               a(1,7) = a(1,7) + w(i) * (bk4)
               a(2,1) = a(2,1) + w(i) * (0.0)
               a(2,2) = a(2,2) + w(i) * (k(i,r,c) * * 2)
               a(2,3) = a(2,3) + w(i) * (0.0)
               a(2,4) = a(2,4) + w(i) * (bk1 * lag(1))
               a(2,5) = a(2,5) + w(i) * (bk2 * lag(2))
               a(2,6) = a(2,6) + w(i) * (bk3 * lag(3))
               a(2,7) = a(2,7) + w(i) * (bk4 * lag(4))
```

a(3,1)=a(3,1)+w(i)*(-k(i,r,c)**2)a(3,2)=a(3,2)+w(i)*(0.0)a(3,3)=a(3,3)+w(i)*(k(i,r,c)**4)a(3,4)=a(3,4)+w(i)*(-k(i,r,c)**2*bk1)a(3,5)=a(3,5)+w(i)*(-k(i,r,c)**2*bk2)a(3,6)=a(3,6)+w(i)*(-k(i,r,c)**2*bk3)a(3,7)=a(3,7)+w(i)*(-k(i,r,c)**2*bk4)a(4,1)=a(4,1)+w(i)*(bk1)a(4,2)=a(4,2)+w(i)*(bkl*lag(1))a(4,3)=a(4,3)+w(i)*(-k(i,r,c)**2*bk1)a(4,4) = a(4,4) + w(i) * (bk1 * 2 * (1 + (lag(1) * 2/k(i,r,c) * 2)))a(4,5)=a(4,5)+w(i)*(bk1*bk2*(1+(lag(1)*lag(2)/1 k(i,r,c) * * 2))a(4,6) = a(4,6) + w(i) * (bk1 * bk3 * (1+(lag(1) * lag(3)))1 k(i,r,c) **2))) a(4,7)=a(4,7)+w(i)*(bk1*bk4*(1+(lag(1)*lag(4)/k(i,r,c)**2))) 1 a(5,1) = a(5,1) + w(i) * (bk2)a(5,2)=a(5,2)+w(i)*(bk2*lag(2))a(5,3)=a(5,3)+w(i)*(-k(i,r,c)**2*bk2)a(5,4)=a(5,4)+w(i)*(bk1*bk2*(1+(lag(1)*lag(2)/k(i,r,c)**2))) 1 a(5,5) = a(5,5) + w(i) * (bk2*bk2*(1+(lag(2)*lag(2)/1 k(i,r,c) * * 2)))a(5,6)=a(5,6)+w(i)*(bk3*bk2*(1+(lag(3)*lag(2)/1 k(i,r,c) * * 2)))a(5,7)=a(5,7)+w(i)*(bk4*bk2*(1+(lag(4)*lag(2)/ 1 k(i,r,c) * * 2)))a(6,1)=a(6,1)+w(i)*(bk3)a(6,2)=a(6,2)+w(i)*(bk3*lag(3))a(6,3)=a(6,3)+w(i)*(-k(i,r,c)**2*bk3)a(6,4)=a(6,4)+w(i)*(bk1*bk3*(1+(lag(1)*lag(3)/ 1 k(i,r,c) * * 2)))a(6,5)=a(6,5)+w(i)*(bk2*bk3*(1+(lag(2)*lag(3)/ 1 k(i,r,c) * * 2)))a(6,6)=a(6,6)+w(i)*(bk3*bk3*(1+(lag(3)*lag(3)/ 1 k(i,r,c) * * 2))a(6,7)=a(6,7)+w(i)*(bk4*bk3*(1+(lag(4)*lag(3)/1 k(i,r,c) * * 2))a(7,1)=a(7,1)+w(1)*(bk4)a(7,2)=a(7,2)+w(i)*(bk4*lag(4))a(7,3)=a(7,3)+w(i)*(-k(i,r,c)**2*bk4)a(7,4)=a(7,4)+w(i)*(bk1*bk4*(1+(lag(1)*lag(4)/ 1 k(i,r,c) * * 2)))a(7,5)=a(7,5)+w(i)*(bk2*bk4*(1+(lag(2)*lag(4)/1 k(i,r,c) * * 2)))a(7,6)=a(7,6)+w(i)*(bk3*bk4*(1+(lag(3)*lag(4)/1 k(i,r,c) * * 2)))a(7,7)=a(7,7)+w(i)*(bk4*bk4*(1+(lag(4)*lag(4)/))1 k(i,r,c) * * 2)))b(1,1)=b(1,1)+w(i)*(f(i,r,c))b(2,1)=b(2,1)+w(i)*(g(i,r,c)*k(i,r,c))b(3,1)=b(3,1)+w(i)*(-f(i,r,c)*k(i,r,c)**2)b(4,1)=b(4,1)+w(i)*(bk1*(f(i,r,c)+g(i,r,c)*

```
1
              lag(1)/k(i,r,c))
              b(5,1)=b(5,1)+w(i)*(bk2*(f(i,r,c)+g(i,r,c)*)
     1
              lag(2)/k(i,r,c))
              b(6,1)=b(6,1)+w(i)*(bk3*(f(i,r,c)+q(i,r,c)*)
     1
              lag(3)/k(i,r,c))
              b(7,1)=b(7,1)+w(i)*(bk4*(f(i,r,c)+g(i,r,c)*)
     1
              lag(4)/k(i,r,c))
  175
          continue
 205
              format(8f10.4)
 206
              format(' ')
С
      PROGRAM INVMAT
С
С
      A(N,NF) -- CHANGE BOTH THE DECLARATION I.E REAL A(N,NF), AS WELL
С
      AS THE DECLARED VALUES OF N & NF. ALSO CHANGE THE NF VALUE IN THE
С
      FORMAT STATEMENT AT THE END. I.E (NFF6.2)
С
      DO 4 ROW=1, IN
     DO 3 COL=(IN+1), NF
       A(ROW, COL) = 0.0
       IF((ROW+IN).EQ.COL) A(ROW,COL)=1.0
3
     CONTINUE
4
      CONTINUE
      DO 50 D=1, IN
     TEMP=A(D, D)
     DO 5 COL=D,NF
       A(D, COL) = A(D, COL) / TEMP
5
     CONTINUE
     DO 15 ROW=D+1, IN
       TEMP = A(ROW, D)
       DO 10 COL=D,NF
          A(ROW, COL) = A(ROW, C) - TEMP*A(D, COL)
10
       CONTINUE
15
     CONTINUE
50
      CONTINUE
      DO 100 D=IN, 2, -1
     DO 80 ROW=D-1,1,-1
       TEMP=A(ROW, D)
       DO 70 COL=D,NF
          A(ROW, COL) = A(ROW, COL) - TEMP*A(D, COL)
70
       CONTINUE
80
     CONTINUE
100
      CONTINUE
      DO 120 ROW=1, IN
      DO 125 COL=IN+1,NF
      X(ROW, 1) = X(ROW, 1) + A(ROW, COL) + B(COL-IN, 1)
125
       CONTINUE
       CONTINUE
120
      FORMAT(' ',9F15.5)
130
C CALCULATION OF APPROXIMATE FUNCTION
С
           do 300 i≕1,m
              p=dcmplx(0.0,k(i,r,c))
              ckquad=x(1,1)+x(2,1)*p+x(3,1)*p**2
              cklag=x(4,1)*p/(p+lag(1))+x(5,1)*p/(p+lag(2))+x(6,1)*
```

	1	<pre>p/(p+lag(3 ckbar(i,r, mn(i)=F(i, if(mn(i).1</pre>))+x(7,1)*p/ c)=ckquad+ck r,c)**2+G(i, t.1.0) mn(i)	(p+lag(4)) lag r,c)**2 =1.0		
		err=err+((dreal(ckbar(i,r,c))-f(i	,r,c))**2+	
	1	(dimag(ckb	ar(i,r,c))-g	(i,r,c))**2))/mn(i)	
300		continue				
301		FN=FN+err				
305		format(6X,	f12.6,',',f1	0.5 ,',',f 10	.5,',',f10.5	, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1,
	1	f10.5,',',	2f10.5)			
		print 634,	count			
500		continue				
600		if(count.lt.i	count-1) got	o 750		
		write (1,602)	r,c			
		write (1,604)				
		write (1,604)				
		write (1,603)				
		write (1,604)				
		write (1,605)	x(1,1),x(2,	1), x(3, 1), x	(4,1)	
		write (1,604)				
		write (1,601)				
		write (1,604)				
		write (1,605)	x(5,1), x(6,	1), x(7, 1)		
		write (1,604)				
		write (1,604)				
		write (1,606)				
		write (1,604)				
		write (1,635)	<pre>lag(1),lag(</pre>	2),lag(3),la	ag(4),err,FN	
		write (1,604)				
		write (1,604)				
		write (1,610)				
		write (1,611)				
		write (1,604)				
		write (1,604)				
602		format(6X, ' A	PPROXIMANT F	UNCTION	N (',i2,'	,',i2,')')
603		format(5x, '	AO	A1	A2	A3')
601		format(5x, '	A4	A5	A6')	•
604		format(' ')				
605		format(5x,8f1	6.8)			
634		format(215,7f	12.9)			
635		format(5x,7f1	2.9)			
606		format(5x, '	lag1	lag2	lag3	lag4',
	1	' err	or')			
610		format(6X, '	k	EXACT		APPROXIMATE')
611		format(6X, '		REAL	IMAG	REAL '
	1	' IMAG')				
		do $650 \text{ prt}=1.$	m			
		write(1,305)	k(prt.r.c).f	(prt.r.c).a	(prt.r.c).ck	bar(prt.r.c)
650		continue		(=========	(==-/=/-//	
c						
C CRI	EATE	FILE 'LAPLACE.	DAT' CONTAIN	S b1, b2.	AO. A1.	
		write (10.605) $x(1,1).x(2)$	(1), x(3.1)	x(4,1).x(5.1)	.),
	1	x(6,1).x(7,1)	· · · · · · · · · · · · · · · · · · ·	, , , , , , - , ,	· · - / · · · · · · · · · · · · · · · ·	* *
	_	if(r.eg.1.and	.c.eg.2) got	o 584		

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584 C	goto 750 write (10,635)	lag(1),lag(2),lag(3),lag(4),FN
750	continue RETURN	
800	END	

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3. #

INPUT	FILE :	N11.DAT
10.000	0.5006	-0.0124
6.000	0.5017	-0.0206
4.000	0.5037	-0.0305
3.000	0.5063	-0.0400
2.000	0.5129	-0.0577
1.500	0.5210	-0.0736
1.200	0.5300	-0.0877
1.000	0.5394	-0.1003
0.990	0.5400	-0.1010
0.880	0.5474	-0.1095
0.800	0.5541	-0.1165
0.770	0.5570	-0.1193
0.660	0.5699	-0.1308
0.600	0.5788	-0.1378
0.560	0.5857	-C.1428
0.550	0.5876	-0.1441
0.500	0.5979	-0.1507
0.440	0.6130	-0.1592
0.400	0.6250	-0.1650
0.340	0.6469	-0.1738
0.330	0.6512	-0.1752
0.325	0.6535	-0.1759
0.320	0.6558	-0.1766
0.315	0.6581	-0.1773
0.310	0.6604	-0.1779
0.305	0.6627	-0.1786
0.300	0.6650	-0.1793
0.240	0.6989	-0.1862
0.220	0.7125	-0.1877
0.200	0.7276	-0.1886
0.160	0.7628	-0.1876
0.120	0.8063	-0.1801
0.110	0.8188	-0.1766
0.100	0.8320	-0.1723
0.080	0.8604	-0.1604
0.060	0.8920	-0.1426
0.050	0.9090	-0.1305
0.040	0.9267	-0.1160
0.025	0.9545	-0.0872
0.010	0.9824	-0.0482
DATA	TABLE	C(k)
k	F	iG

INPUT FILE : N12.DAT

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10.000000	.1240000	5,0060000
6.000000	.1236000	3.0102000
4.0000000	.1220000	2.0148000
3.000000	.1200000	1.5189000
2.0000000	.1154000	1.0258000
1.5000000	.1104000	.7815000
1.2000000	.1052400	.6360000
1.0000000	.1003000	.5394000
.9900000	.0999900	.5346000
.8800000	.0963600	.4817120
.8000000	.0932000	.4432800
.7700000	.0918610	.4288900
.6600000	.0863280	.3761340
.6000000	.0826800	.3472800
.5600000	.0799680	.3279920
.5500000	.0792550	.3231800
.5000000	.0753500	.2989500
- 4400000	.0700480	.2697200
- 4000000	.0660000	.2500000
.3400000	.0590920	.2199460
.3300000	.0578160	.2148960
.3250000	.0571675	.2123875
.3200000	.0565120	.2098560
.3150000	.0558495	.2073015
.3100000	.0551490	.2047240
.3050000	.0544730	.2021235
.3000000	.0537900	.1995000
.2400000	.0446880	.1677360
.2200000	.0412940	.1567500
.2000000	.0377200	.1455200
.1600000	.0300160	.1220480
.1200000	.0216120	.0967560
.1100000	.0194260	.0900680
.1000000	.0172300	.0832000
.0800000	.0128320	.0688320
.0600000	.0085560	.0535200
.0500000	.0065250	.0454500
.0400000	.0046400	.0370680
.0250000	.0021800	.0238625
.0100000	.0004820	.0098240

DATA TABLE C(k) *ik

k

٩. 2

-G*k iF*k

OUTPUT FILE : DATA.OUT

N.X

SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION

260 TRIALS USED

ESTIMATES

	PARAMETER	NEXT-TO-BEST	BEST
1		.0154409	.0154409
2		.0842865	.0842865
3		.2559507	.2559507
4		.7651135	.7651136

FUNCTION VALUES

NEXT-TO-BEST	BEST
.000011782	.000011782

OUTPUT FILE : LAG.OUT

APPROXIMA	NT FUNCTIO	N N	(1, 1)		
AO		A1	A2	А3	
.99	828226	000035	62	.00000474	04120375
A4		A5	A 6		
16	289333	229087	45 -	06459371	
lagl	. 1	lag2	lag3	lag4	error
.01544085	2 .0842	286467 .	255950673	.765113621	.000009348

k EXACT APPROXIMATE REAL IMAG REAL IMAG .50060, -.01240, .50057, -.01257, 10.000000, -.02060, -.02047, 6.000000, .50182, .50170, -.03025, .50372, 4.000000, .50370, -.03050, 3.000000, .50630, -.04000, .50619, -.03976, 2.000000, .51271, -.05764, .51290, -.05770, 1.500000, .52100, -.07360, .52083, -.07373, -.08798, .52994, 1.200000, .53000, -.08770, -.10030, -.10051, .53940, .53959, 1.000000, .990000, .54000, -.10100, .54019, -.10122,

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.015440852	.084286467	.25595067	.76511	3621 .000002434
lagi	lag2	lag3	lag4	error
.013804	89 .058	48218	.04963449	
A4	Α5	A 6		
000013	55 .500	37475	00003629	.00061076
AO	A1	λ2	A3	
APPROXIMANT	FUNCTION	N (1, 2)	
.010000,	.98240,	04820,	.98349,	04764,
.025000,	.95450,	08720,	.95305,	08710,
.040000,	.92670,	11600,	.92684,	11525,
.050000,	.90900,	13050,	.90957	13041,
.060000.	.89200.	14260.	.89253	14285.
. 080000	.86040	16040	.03130, 86027	-16092
.110000,	.81880,	17660,	.81822,	17652,
.120000,	.80630,	18010,	.80584,	17977,
.160000,	.76280,	18760,	.76288,	18703,
.200000,	.72760,	18860,	.72802,	18841,
.220000,	.71250,	18770,	.71297,	18770,
.240000,	.69890,	18620,	.69927,	18633,
.300000,	.66500,	17930,	.66502,	17961,
.305000,	.66270,	17860,	.66257,	17894,
.310000,	.66040,	17790,	.66017,	17825,
.315000,	.65810,	17730,	.65782.	17756.
.320000,	.65580,	17660.	.65552.	17686
.325000,	.65350.	17590.	.65128	-17614
.340000,	.04090,	17380,	• 04083, 65109	17542
.400000,	.62500,	16500,	.62477,	16499,
.440000,	.61300,	15920,	.61278,	15904,
.500000,	.59790,	15070,	.59787,	15049,
.550000,	.58760,	14410,	.58764,	14381,
.560000,	.58570,	14280,	.58579,	14253,
.600000,	.57880,	13780,	.57895,	13757,
.660000,	.56990,	13080,	.57012,	13063,
.770000,	.55700,	11930,	.55733.	11934 .
.800000,	.55410.	11650.	.55443.	11655
.880000.	.54740.	10950.	54767	- 10963

.000011782

1

k	EXACT		APPROXIMATE	
	REAL	IMAG	REAL	IMAG
10.000000,	.12400,	5.00600,	.12582,	5.00914,
6.000000,	.12360,	3.01020,	.12292,	3.01116,

		0 01 1 0 0		0 00 4 6 0
4.000000,	.12200,	2.01480,	.12110,	2.01468,
3.000000,	.12000,	<u>1</u> .51890,	.11938,	1.51835,
2.000000,	.11540,	1.02580,	.11536,	1.02526,
1.500000,	.11040,	.78150,	.11065,	.78113,
1.200000,	.10524,	.63600,	.10561,	.63585,
1.000000,	.10030,	.53940,	.10053,	.53954,
.990000,	. 09999	.53460,	.10023,	.53474,
.880000,	.09636,	.48171,	.09649,	.48191,
.800000.	.09320	.44328,	.09325,	.44352,
.770000	.09186,	.42889,	.09190,	.42912,
.660000,	.08633,	.37613,	.08621,	.37626,
.600000,	.08268,	.34728,	.08253,	.34735,
.560000,	.07997,	.32799,	.07980,	.32802,
.550000,	.07926,	.32318,	.07908,	.32318,
.500000,	.07535,	. 29895,	.07522,	.29892,
.440000,	.07005,	.26972,	.06995,	.26961,
.400000,	.06600,	.25000,	.06597,	.24989,
.340000,	.05909,	.21995,	.05912,	.21990,
.330000,	.05782,	.21490,	.05786,	.21483,
.325000,	.05717,	.21239,	.05722,	.21229,
.320000,	.05651,	.20986,	.05656,	.20974,
.315000,	.05585,	.20730,	.05590,	.20719,
.310000,	.05515,	.20472,	.05523,	.20463,
.305000	.05447	.20212.	.05455,	.20206,
.300000.	.05379	.19950,	.05385,	.19948,
.240000.	.04469.	.16774.	.04469.	.16780,
.220000.	.04129.	.15675.	.04127.	.15683,
.200000.	.03772.	.14552.	.03766.	.14558,
.160000.	.03002.	.12205.	.02991.	.12204.
.120000.	.02161.	.09676.	.02156.	.09669.
.110000.	.01943.	.09007.	.01941.	.09000,
.100000.	.01723.	.08320.	.01723.	.08313.
.080000	.01283.	.06883.	.01286.	.06882.
.060000	.00856	.05352	.00855	.05356
.050000	.00653	.04545	.00650	.04548.
.040000	.00464	.03707	.00459	.03708
.025000	.00218	.02386	.00215.	.02382
.010000	. 00048	.00982	.00046	.00983
	1000407			

OUTPUT FILE : LAPLACE.DAT

.9 16289332	9828226 22908	00003562 74506	0000047. 459372	0 4120375
0 .01380489	0001355	.50037475 217 .04	0000362 963449	.00061076
.015440	851 .08428	6462 .255950	657 .76511354	9.000011782

С С SEQUENTIAL SIMPLEX PROGRAM (MINIMIZATION PROBLEM SOLVING) С RATIONAL FUNCTION APPROXIMATION - TRANSONIC REGIME С С С OPERATION С ICOUNT, N, REQMIN, START(I), STEP(I) REQUIRE: 1) INPUT OF С 2) USER SPECIFIED FUNCTION SUBPROGRAM (DOUBLE PRECISION С FUNCTION FN(X) - MINIMIZATION FUNCTION) С С DOUBLE PRECISION START(20), STEP(20), XMIN(20), 1XSEC(20), YNEWLO, YSEC, REQMIN С OPEN (UNIT=6, FILE='DATA.OUT') С С С * С ICOUNT, N, REQMIN, START(N), STEP(N) * С С ***** С C LET N=4 AT ALL TIMES, SIMPLY SET START(I)=0.0 AND STEP(I)=0.0 C ICOUNT=450N=4REQMIN=0.00000000001 START(1) = 0.000592052START(2) = 0.07383689 **START(3)** =0.27305258 START(4) = 0.47381410STEP(1) = 0.1STEP(2) = 0.1051STEP(3) = 0.210412STEP(4) = 0.310323С С DO 60 I=1, NXMIN(I) = 0.D0XSEC(I) = 0.D060 CONTINUE YNEWLO=0.D0 YSEC=0.D0 С С CALL NELDER-MEAD SUBROUTINE С CALL NELMIN (N, START, XMIN, XSEC, YNEWLO, YSEC, **1REQMIN**, STEP, ICOUNT) С С OUTPUT FROM PROGRAM С WRITE(6,64) WRITE(6,65) ICOUNT WRITE(6,75)

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```
WRITE(6,77)
      DO 79 I=1,N
 79
      WRITE(6,80) I,XSEC(I),XMIN(I)
      WRITE(6,82)
      WRITE(6,83)
      WRITE(6,84) YSEC, YNEWLO
С
 64
      FORMAT(6X,1H //,6x,1H ,
     142H SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION)
      FORMAT(6X,1H //,6x,1H ,I5,12H TRIALS USED/)
 65
      FORMAT(6X,1H ,21X,9HESTIMATES/)
 75
      FORMAT(6X,1H,9HPARAMETER,7X,12HNEXT-TO-BEST,8X,
 77
     14HBEST/)
      FORMAT(6X,1H,15,2F20.7)
 80
      FORMAT(6X,1H //,1H ,6X,15HFUNCTION VALUES/)
 82
      FORMAT(6X,1H,5X,13H NEXT-TO-BEST,8X,4HBEST/)
 83
 84
      FORMAT(6X,1H,2F15.9)
      STOP
      END
С
С
             NELDER-MEAD SUBROUTINE
С
      SUBROUTINE NELMIN(N, START, XMIN, XSEC, YNEWLO, YSEC,
     1REQMIN, STEP, ICOUNT)
      DOUBLE PRECISION START(N), STEP(N), XMIN(N),
     1XSEC(N), YNEWLO, YSEC, REQMIN, P(20,21), PSTAR(20),
     2P2STAR(20), PBAR(20), Y(20), DN, Z, YLO, RCOEFF,
     3YSTAR, ECOEFF, Y2STAR, CCOEFF, FN, DABIT, DCHK,
     4COORD1, COORD2
      DATA RCOEFF/1.0D0/, ECOEFF/2.0D0/, CCOEFF/0.5D0/
      KCOUNT=ICOUNT
      ICOUNT=0
С
      IF(REQMIN.LE.O.DO) ICOUNT=ICOUNT-1
      IF(N.LE.0) ICOUNT=ICOUNT-10
      IF(N.GT.20) ICOUNT=ICOUNT-10
      IF(ICOUNT.LT.0) RETURN
С
      DABIT=2.04607D-35
      BIGNUM=1.0D38
      KONVGE=5
      XN = FLOAT(N)
      DN=DFLOAT(N)
      NN=N+1
С
С
      CONSTRUCTION OF INITIAL SIMPLEX
С
 1001 DO 1 I=1,N
    1 P(I,NN) = START(I)
      Y(NN) = FN(START)
      ICOUNT=ICOUNT+1
      DO 2 J=1,N
      DCHK=START(J)
      START(J) = DCHK + STEP(J)
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DO 3 I=1,N 3 P(I,J) = START(I)Y(J) = FN(START)ICOUNT=ICOUNT+1 2 START(J)=DCHK С С SIMPLEX CONSTRUCTION COMPLETE С С FIND HIGHEST AND LOWEST Y VALUES С YNEWLO (Y(IHI)) INDICATES THE VERTEX OF С THE SIMPLEX TO BE REPLACED С 1000 YLO=Y(1)YNEWLO=YLO ILO=1IHI=1 DO 5 I=2,NNIF(Y(I).GE.YLO) GOTO 4 YLO=Y(I)ILO=I 4 IF(Y(I).LE.YNEWLO) GOTO 5 YNEWLO=Y(I)IHI=I **5 CONTINUE** С С PERFORM CONVERGENCE CHECKS ON FUNCTION С DCHK=(YNEWLO+DABIT)/(YLO+DABIT)-1.D0 IF (DABS (DCHK).LT.REOMIN) GOTO 900 С KONVGE=KONVGE-1 IF(KONVGE.NE.0) GOTO 2020 KONVGE=5С С CHECK CONVERGENCE OF COORDINATES ONLY С EVERY 5 SIMPLEXES С DO 2015 I=1,N COORD1 = P(I, 1)COORD2=COORD1 DO 2010 J=2,NN IF(P(I,J).GE.COORD1) GOTO 2005 COORD1 = P(I, J)2005 IF(P(I,J).LE.COORD2) GOTO 2010 COORD2 = P(I,J)2010 CONTINUE DCHK=(COORD2+DABIT)/(COORD1+DABIT)-1.D0 IF(DABS(DCHK).GT.REQMIN) GOTO 2020 2015 CONTINUE **GOTO 900** 2020 IF(ICOUNT.GE.KCOUNT) GOTO 900 С С CALCULATE PBAR, THE CENTROID OF THE С SIMPLEX VERTICES EXCEPTING THAT WITH

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Y VALUE YNEWLO С С DO 7 I=1,N Z=0.0D0 DO 6 J=1,NN6 Z=Z+P(I,J)Z=Z-P(I,IHI)7 PBAR(I) = Z/DNС С REFLECTION THROUGH THE CENTROID С DO 8 I=1,N 8 PSTAR(I)=(1.0D0+RCOEFF)*PBAR(I)-RCOEFF*P(I,IHI) **YSTAR=FN(PSTAR)** ICOUNT=ICOUNT+1 IF(YSTAR.GE.YLO) GOTO 12 IF (ICOUNT.GE.KCOUNT) GOTO 19 С С SUCCESSFUL REFLECTION, SO EXTENSION С DO 9 I=1,N 9 P2STAR(I)=ECOEFF*PSTAR(I)+(1.0D0-ECOEFF)*PBAR(I) Y2STAR=FN(P2STAR) ICOUNT=ICOUNT+1 С С RETAIN EXTENSION OR CONTRACTION С IF(Y2STAR.GE.YSTAR) GOTO 19 10 DO 11 I=1,N 11 P(I, IHI) = P2STAR(I)Y(IHI)=Y2STAR GOTO 1000 С С NO EXTENSION С 12 L=0 DO 13 I=1,NN IF(Y(I).GT.YSTAR) L=L+1 **13 CONTINUE** IF(L.GT.1) GOTO 19 IF(L.EQ.0) GOTO 15 С С CONTRACTION ON THE REFLECTION SIDE OF THE С CENTROID С DO 14 I=1,N 14 P(I, IHI) = PSTAR(I)Y(IHI)=YSTAR С С CONTRACTION ON THE Y(IHI) SIDE OF THE CENTROID С 15 IF(ICOUNT.GE.KCOUNT) GOTO 900 DO 16 I=1,N 16 P2STAR(I) = CCOEFF*P(I,IHI)+(1.0D0-CCOEFF)*PBAR(I)

```
Y2STAR=FN (P2STAR)
      ICOUNT=ICOUNT+1
      IF(Y2STAR.LT.Y(IHI)) GOTO 10
С
С
         CONTRACT THE WHOLE SIMPLEX
С
      DO 18 J=1,NN
      DO 17 I=1,N
      P(I,J) = (P(I,J) + P(I,ILO)) *0.5D0
   17 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   18 CONTINUE
      ICOUNT=ICOUNT+NN
      IF(ICOUNT.LT.KCOUNT) GOTO 1000
      GOTO 900
С
С
         RETAIN REFLECTION
С
С
    19 CONTINUE
   19 DO 20 I=1,N
   20 P(I, IHI) = PSTAR(I)
      Y(IHI)=YSTAR
      GOTO 1000
С
С
         SELECT THE TWO BEST FUNCTION VALUES (YNEWLO
С
         AND YSEC) AND THEIR COORDS. (XMIN AND XSEC)
С
  900 DO 23 J=1,NN
      DO 22 I=1,N
   22 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   23 CONTINUE
      YNEWLO=BIGNUM
      DO 24 J=1,NN
      IF(Y(J).GE.YNEWLO) GOTO 24
      YNEWLO=Y(J)
      IBEST=J
   24 CONTINUE
      Y(IBEST)=BIGNUM
      YSEC=BIGNUM
      DO 25 J=1,NN
      IF(Y(J).GE.YSEC) GOTO 25
      YSEC=Y(J)
      ISEC=J
   25 CONTINUE
      DO 26 I=1,N
      XMIN(I) = P(I, IBEST)
      XSEC(I) = P(I, ISEC)
   26 CONTINUE
      RETURN
      END
С
С
     APPROXIMATING FUNCTION : PADE APPROXIMATE WITH LAG TERMS
С
```

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1
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```
DOUBLE PRECISION FUNCTION FN(lag)
       IMPLICIT REAL*8 (A-Z)
       COMPLEX*16 ckquad, cklag, ckbar(40,2,2), p
       DOUBLE PRECISION LAG(4)
       REAL*8 w(40),k(40,2,2),F(40,2,2),G(40,2,2),a(9,18),b(9,1),x(9,1)
       REAL*8 mn(40)
       INTEGER i, m, in, nst, nf, j, prt, totlag, c, r, count, icount
       INTEGER D, COL, ROW
С
С
        OPEN(UNIT=1, FILE='LAG.OUT')
        OPEN(UNIT=10, FILE='LAPLACE.DAT')
С
С
С
       С
       *
С
          OPEN REQUIRED INPUT TABULAR DATA FILES
       *
                                                     *
С
       *
          TOTAL # OF ITERATIONS ICOUNT (2)
С
       *
           # OF POINTS (m) IN FILES
С
       *
          CHECK FORCED CONDITIONS
С
       *
С
       С
          icount=450+5
         minerr=0.0000167
         totlag=4
         m=15
         count=count+1
         if(count.gt.1) goto 88
С
С
     FILES CONTAIN
                    k(i,r,c), F(i,r,c), G(i,r,c)
С
    k(i) REDUCED FREQUENCY
С
     F(i) REAL
С
    G(i) IMAGINARY
С
C APPROXIMATING UNSTEADY AERODYNAMICS
        OPEN(UNIT=2, FILE='N11.DAT', STATUS='OLD')
        OPEN(UNIT=4, FILE='N12.DAT', STATUS='OLD')
        OPEN (UNIT=7, FILE='N21.DAT', STATUS='OLD')
        OPEN(UNIT=8, FILE='N22.DAT', STATUS='OLD')
        do 150 i=1,m
            READ(2,*) k(i,1,1),F(i,1,1),G(i,1,1)
            READ(4,*) k(i,1,2),F(i,1,2),G(i,1,2)
            READ(7,*) k(i,2,1),F(i,2,1),G(i,2,1)
            READ(8,*) k(i,2,2),F(i,2,2),G(i,2,2)
 150
        continue
        CLOSE(UNIT=8)
        CLOSE(UNIT=7)
        CLOSE(UNIT=4)
        CLOSE(UNIT=2)
С
С
    REQUIRED CONSTANTS
С
            matrix [A] (size in x in )
88
            nst=4
```

```
in=nst+(totlag-1)
              nf=2*in
C
C NOTE: lag(1) <> lag(2) <> lag(3) <> lag(4)
С
         lag(1), lag(2), lag(3), lag(4) must be all non-negative
C BOUNDARY (CONSTRAINT) CONDITIONS
          do 156 i=1, totlag
              if(lag(i).lt.0.0) lag(i)=i*0.000000001
           do 157 j=1,totlag-1
              if(i.eq.j) goto 157
              if(lag(i).eq.lag(j)) goto 163
              goto 157
              print 164, j
 163
              format(i5, ' EQUAL LAG TERMS REQUIRES IMMEDIATE STOP')
 164
              stop
 157
           continue
 156
           continue
С
              FN=0.0
С
           do 750 r=1,2
           do 700 c=1,2
              err=0.0
С
     RESETTING MATRIX ELEMENT VALUES TO ZERO
 390
           do 410 i=1,in
           do 400 j=1,nf
              a(i,j) = 0.0
 400
           continue
              b(i,1)=0.0
              x(i,1) = 0.0
 410
           continue
С
С
     CREATING MATRICES [A] AND [B] :
                                          SOLVE
                                                  [A] \{x\} = [B]
           do 175 i=1,m
              if(k(i,r,c).eq.0.0) qoto 175
              bk1=k(i,r,c)**2/(k(i,r,c)**2+lag(1)**2)
              if(totlag.eq.1) goto 415
              bk2=k(i,r,c)**2/(k(i,r,c)**2+lag(2)**2)
              if(totlag.eq.2) goto 415
              bk3=k(i,r,c)**2/(k(i,r,c)**2+lag(3)**2)
              if(totlag.eq.3) goto 415
              bk4=k(i,r,c)**2/(k(i,r,c)**2+lag(4)**2)
 415
              mn(i) = F(i,r,c) * * 2 + G(i,r,c) * * 2
              if(mn(i).lt.1.0) mn(i)=1.0
              w(i) = 1.0
              w(i) = w(i) / mn(i)
              a(1,1) = a(1,1) + w(i) * (1.0)
              a(1,2) = a(1,2) + w(i) * (0.0)
              a(1,3) = a(1,3) + w(i) * (-k(i,r,c) * * 2)
              a(1,4) = a(1,4) + w(i) * (bk1)
              a(1,5) = a(1,5) + w(i) * (bk2)
              a(1,6) = a(1,6) + w(i) * (bk3)
              a(1,7) = a(1,7) + w(i) * (bk4)
              a(2,1) = a(2,1) + w(i) * (0.0)
```

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```
a(2,2)=a(2,2)+w(i)*(k(i,r,c)**2)
        a(2,3) = a(2,3) + w(i) * (0.0)
        a(2,4) = a(2,4) + w(i) * (bk1 * lag(1))
        a(2,5) = a(2,5) + w(i) * (bk2 * lag(2))
        a(2,6) = a(2,6) + w(i) * (bk3 * lag(3))
        a(2,7) = a(2,7) + w(i) * (bk4 * lag(4))
        a(3,1)=a(3,1)+w(i)*(-k(i,r,c)**2)
        a(3,2)=a(3,2)+w(i)*(0.0)
        a(3,3)=a(3,3)+w(i)*(k(i,r,c)**4)
        a(3,4)=a(3,4)+w(i)*(-k(i,r,c)**2*bk1)
        a(3,5)=a(3,5)+w(i)*(-k(i,r,c)**2*bk2)
        a(3,6)=a(3,6)+w(i)*(-k(i,r,c)**2*bk3)
        a(3,7)=a(3,7)+w(i)*(-k(i,r,c)**2*bk4)
        a(4,1)=a(4,1)+w(i)*(bk1)
        a(4,2) = a(4,2) + w(i) * (bk1 * lag(1))
        a(4,3)=a(4,3)+w(i)*(-k(i,r,c)**2*bk1)
        a(4,4) = a(4,4) + w(i) * (bk1**2*(1+(lag(1)**2/k(i,r,c)**2)))
        a(4,5)=a(4,5)+w(i)*(bk1*bk2*(1+(lag(1)*lag(2)/
1
        k(i,r,c) * * 2)))
        a(4,6) = a(4,6) + w(i) * (bk1 * bk3 * (1 + (lag(1) * lag(3)))
1
        k(i,r,c)**2))
        a(4,7)=a(4,7)+w(i)*(bk1*bk4*(1+(lag(1)*lag(4)/
1
        k(i,r,c)**2)))
        a(5,1)=a(5,1)+w(i)*(bk2)
        a(5,2) = a(5,2) + w(i) * (bk2 * lag(2))
        a(5,3)=a(5,3)+w(i)*(-k(i,r,c)**2*bk2)
        a(5,4)=a(5,4)+w(i)*(bk1*bk2*(1+(lag(1)*lag(2)/
1
        k(i,r,c)**2)))
        a(5,5) = a(5,5) + w(i) * (bk2 * bk2 * (1+(lag(2) * lag(2)))
1
        k(i,r,c)**2)))
        a(5,6)=a(5,6)+w(i)*(bk3*bk2*(1+(lag(3)*lag(2)/
1
        k(i,r,c)**2)))
        a(5,7) = a(5,7) + w(i) * (bk4 * bk2 * (1 + (lag(4) * lag(2)))
1
        k(i,r,c)**2)))
        a(6,1) = a(6,1) + w(i) * (bk3)
        a(6,2)=a(6,2)+w(i)*(bk3*lag(3))
        a(6,3)=a(6,3)+w(i)*(-k(i,r,c)**2*bk3)
        a(6,4)=a(6,4)+w(i)*(bk1*bk3*(1+(lag(1)*lag(3)/
1
        k(i,r,c)**2)))
        a(6,5) = a(6,5) + w(i) * (bk2 * bk3 * (1 + (lag(2) * lag(3)))
1
        k(i,r,c) * * 2)))
        a(6,6)=a(6,6)+w(i)*(bk3*bk3*(1+(lag(3)*lag(3)/))
1
        k(i,r,c) * * 2)))
        a(6,7)=a(6,7)+w(i)*(bk4*bk3*(1+(lag(4)*lag(3)/)))
        k(i,r,c)**2)))
1
        a(7,1)=a(7,1)+w(i)*(bk4)
        a(7,2)=a(7,2)+w(i)*(bk4*lag(4))
        a(7,3)=a(7,3)+w(i)*(-k(i,r,c)**2*bk4)
        a(7,4)=a(7,4)+w(i)*(bk1*bk4*(1+(lag(1)*lag(4)/
1
        k(i,r,c)**2)))
        a(7,5)=a(7,5)+w(i)*(bk2*bk4*(1+(lag(2)*lag(4)/
1
        k(i,r,c) * * 2)))
        a(7,6) = a(7,6) + w(i) * (bk3 * bk4 * (1+(lag(3) * lag(4)))
1
        k(i,r,c) * * 2))
```

```
A - 82
```

```
a(7,7)=a(7,7)+w(i)*(bk4*bk4*(1+(lag(4)*lag(4)/
     1
             k(i,r,c)**2)))
             b(1,1)=b(1,1)+w(i)*(f(i,r,c))
             b(2,1)=b(2,1)+w(i)*(g(i,r,c)*k(i,r,c))
             b(3,1)=b(3,1)+w(i)*(-f(i,r,c)*k(i,r,c)**2)
              b(4,1)=b(4,1)+w(i)*(bk1*(f(i,r,c)+g(i,r,c)*
     1
              lag(1)/k(i,r,c))
              b(5,1)=b(5,1)+w(i)*(bk2*(f(i,r,c)+g(i,r,c)*
     1
              lag(2)/k(i,r,c))
              b(6,1)=b(6,1)+w(i)*(bk3*(f(i,r,c)+g(i,r,c)*
     1
              lag(3)/k(i,r,c))
              b(7,1)=b(7,1)+w(i)*(bk4*(f(i,r,c)+g(i,r,c)*)
     1
              lag(4)/k(i,r,c))
  175
          continue
 205
              format(8f10.4)
 206
              format('
                       • • )
С
      PROGRAM INVMAT
С
С
      A(N,NF) -- CHANGE BOTH THE DECLARATION I.E REAL A(N,NF), AS WELL
С
      AS THE DECLARED VALUES OF N & NF. ALSO CHANGE THE NF VALUE IN THE
С
      FORMAT STATEMENT AT THE END. I.E (NFF6.2)
С
С
      DO 4 ROW=1, IN
     DO 3 COL=(IN+1), NF
       A(ROW, COL) = 0.0
       IF((ROW+IN).EQ.COL) A(ROW,COL)=1.0
3
     CONTINUE
4
      CONTINUE
      DO 50 D=1, IN
     TEMP=A(D, D)
     DO 5 COL=D,NF
       A(D, COL) = A(D, COL) / TEMP
5
     CONTINUE
     DO 15 ROW=D+1, IN
       TEMP=A(ROW, D)
       DO 10 COL=D,NF
         A(ROW, COL) = A(ROW, COL) - TEMP * A(D, COL)
10
       CONTINUE
15
     CONTINUE
      CONTINUE
50
      DO 100 D=IN,2,-1
     DO 80 ROW=D-1,1,-1
       TEMP=A(ROW, D)
       DO 70 COL=D,NF
         A(ROW, COL) = A(ROW, COL) - TEMP + A(D, COL)
70
       CONTINUE
80
     CONTINUE
100
      CONTINUE
      DO 120 ROW=1, IN
      DO 125 COL=IN+1, NF
      X(ROW, 1) = X(ROW, 1) + A(ROW, COL) + B(COL-IN, 1)
125
       CONTINUE
120
       CONTINUE
```

```
FORMAT(' ',9F15.5)
130
С
C FORCED AGREEMENT WITH TABULAR DATA AT k=0.0
              if(r.eq.1.and.c.eq.1.and.k(1,1,1).eq.0.0) \times (1,1) = F(1,r,c)
              if(r.eq.2.and.c.eq.1.and.k(1,2,1).eq.0.0) x(1,1)=F(1,r,c)
              if(r.eq.1.and.c.eq.2.and.k(1,1,2).eq.0.0) x(1,1)=F(1,r,c)
С
              if(r.eq.2.and.c.eq.2.and.k(1,2,2).eq.0.0) x(1,1)=F(1,r,c)
С
C CALCULATION OF APPROXIMATE FUNCTION
C
          do 300 i=1,m
             p=dcmplx(0.0,k(i,r,c))
              ckquad=x(1,1)+x(2,1)*p+x(3,1)*p**2
              cklag=x(4,1)*p/(p+lag(1))+x(5,1)*p/(p+lag(2))+x(6,1)*
              p/(p+lag(3))+x(7,1)*p/(p+lag(4))
     1
              ckbar(i,r,c)=ckquad+cklag
c error calculation
              mn(i) = F(i, r, c) * * 2 + G(i, r, c) * * 2
              if(mn(i).lt.1.0) mn(i)=1.0
c square error (complete) #1 #2
              err=err+(((dreal(ckbar(i,r,c))-f(i,r,c))**2+
              (dimag(ckbar(i,r,c))-g(i,r,c))**2)/mn(i))/m
     1
c % square error (separate) #3
               err=err+((dreal(ckbar(i,r,c))-f(i,r,c))**2/f(i,r,c)**2+
С
               (dimag(ckbar(i,r,c))-g(i,r,c))**2/g(i,r,c)**2)/m
С
      1
c % error (separate) #4
С
               err=err+(abs((dreal(ckbar(i,r,c))-f(i,r,c))/f(i,r,c))+
С
      1
               abs((dimag(ckbar(i,r,c))-g(i,r,c))/g(i,r,c)))/m
С
 300
          continue
 301
             FN=FN+err
              print 305, x(1,1),x(2,1),x(3,1),x(4,1),x(5,1),x(6,1),x(7,1)
C
 305
              format(6X,f12.6,',',f10.5,',',f10.5,',',f10.5,',',
              f10.5,',',2f10.5)
     1
               print 634, r,c,lag(1),lag(2),lag(3),lag(4),err,FN
С
              print 634, count
 500
          continue
 600
           if(count.lt.icount-2) goto 700
          write (1,604)
          write (1,602) r,c
          write (1,604)
          write (1,604)
          write (1,603)
          write (1,604)
          write (1,605) \times (1,1), \times (2,1), \times (3,1), \times (4,1)
          write (1,604)
          write (1,601)
          write (1,604)
          write (1,605) \times (5,1), \times (6,1), \times (7,1)
          write (1,604)
          write (1,604)
          write (1,606)
          write (1,604)
          write (1,635) lag(1), lag(2), lag(3), lag(4), err, FN
          write (1,604)
```

```
A - 84
```
```
write (1,604)
           write (1,610)
           write (1,611)
           write (1,604)
           write (1,604)
С
C CREATE FILE 'LAPLACE.DAT' CONTAINS b1, b2, ..., A0, A1, ...
           write (10,605) \times (1,1), \times (2,1), \times (3,1), \times (4,1), \times (5,1),
     1
           x(6,1), x(7,1)
           if(r.eq.2.and.c.eq.2) goto 584
           goto 585
 584
           write (10,635) lag(1),lag(2),lag(3),lag(4),FN
С
 602
           format(6X, ' APPROXIMANT FUNCTION
                                                       N (',i2,',',i2,')')
 603
           format(5x,'
                                A0
                                                  A1
                                                                    A2',
     1
                            A3')
 601
           format(5x,'
                                A4
                                                  A5
                                                                    A6')
           format(' ')
 604
 605
           format(5x, 8f16.6)
 634
           format(215,7f12.9)
 635
           format(7x,7fll.8)
 606
           format(7x, '
                            lag1
                                        lag2
                                                     lag3
                                                                 lag4',
     1
                   error')
 610
           format(6X,'
                                k
                                               EXACT
                                                                     ۰,
     1
           'APPROXIMATE')
 611
           format(6X,'
                                         REAL
                                                      IMAG
                                                                     REAL
                                                                           ١,
                 IMAG')
     1
 585
           do 650 prt=1,m
           write(1,305) k(prt,r,c),f(prt,r,c),g(prt,r,c),ckbar(prt,r,c)
 650
           continue
 700
           continue
 750
           continue
           RETURN
 800
           END
```

0.00	0.00	0.00
0.025	0.17	0.56
0.050	0.39	0.87
0.0875	0.66	1.17
0.100	0.76	1.23
0.125	0.89	1.39
0.150	1.01	1.49
0.175	1.12	1.61
0.200	1.17	1.69
0.25	1.25	1.93
0.3	1.23	2.12
0.35	1 .19	2.39
0.4	1.15	2.64
0.45	1.07	2.93
0.5	0.98	3.25

DATA TABLE

Cl(h)

ų.

k	Real	Imaginary

INPUT FILE : N12.DAT

0.000	14.70	-0.00
0.025	12.20	-4.05
0.050	9.50	-4.15
0.0875	7.63	-3.70
0.100	7.22	-3.47
0.125	6.61	-3.18
0.150	6.09	-2.94
0.175	5.72	-2.59
0.200	5.41	-2.29
0.25	4.92	-1.79
0.3	4.63	-1.30
0.35	4.43	-0.83
0.4	4.35	-0.41
0.45	4.35	-0.03
0.5	4.40	0.28

DATA TABLE

Cl(alpha)

k Real	Imaginary
--------	-----------

0.000	0.000	0.000
0.025	0.001	-0.035
0.050	0.015	-0.065
0.0875	0.025	-0.125
0.100	0.030	-0.145
0.125	0.02	-0.198
0.150	0.015	-0.23
0.175	0.010	-0.280
0.200	-0.005	-0.325
0.25	-0.035	-0.397
0.3	-0.060	-0.403
0.35	-0.045	-0.425
0.4	-0.005	-0.448
0.45	0.065	-0.490
0.5	0.14	-0.55

DATA TABLE

Cm(h)

l

INPUT FILE : N22.DAT

-0.445	0.000
-0.450	-0.027
-0.455	-0.122
-0.495	-0.311
-0.520	-0.348
-0.581	-0.422
-0.672	-0.508
-0.732	-0.548
-0.810	-0.582
-0.905	-0.610
-0.898	-0.550
-0.868	-0.575
-0.824	-0.648
-0.790	-0.750
-0.770	-0.903
	-0.445 -0.450 -0.455 -0.520 -0.581 -0.672 -0.732 -0.810 -0.905 -0.898 -0.868 -0.824 -0.790 -0.770

DATA TABLE

Cm(alpha)

k	Real	Imaginary
---	------	-----------

,

SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION

451 TRIALS USED

ESTIMATES

PARAMETER	NEXT-TO-BEST	BEST
1	.2789266	.2789266
2	.0436573	.0436573
3	.2789914	.2789914
4	.2858671	.2858671

FUNCTION VALUES

NEXT-TO-BEST	BEST
.001414830	.001396017

OUTPUT FILE : LAG.OUT

APPROXIMANT	FUNCTION	1 (1, 1)	
A0	A1	A2	A3
.000000	7.351972	265177	-8561.672085
A4	A 5	A 6	

.593074	8707.625432	-146.231038

	error	lag4	lag3	lag2	lagl
.00020817	.00020817	.28586712	.27899139	.04365735	.27892661

k	EXAC	Т	APPROXIMATE	
	REAL	IMAG	REAL	IMAG
.000000,	.00000,	.00000,	.00000,	.00000,
.025000,	.17000,	.56000,	.16792,	.54877,
.050000,	.39000,	.87000,	.41847,	.86344,
.087500,	.66000,	1.17000,	.69680,	1.16498,
.100000,	.76000,	1.23000,	.77287,	1.24726,

.125000,	.89000,	1.39000,	.90794,	1.39206,
.150000,	1.01000,	1.49000,	1.02060,	1.51454.
.175000,	1.12000,	1.61000,	1.10957,	1.62079.
.200000,	1.17000,	1.69000,	1.17459,	1.71761.
.250000,	1.25000,	1.93000,	1.23994,	1.90689,
.300000,	1.23000,	2.12000,	1.24034,	2.11637,
.350000,	1.19000,	2.39000,	1.20249,	2.35917
.400000,	1.15000,	2.64000,	1.14690,	2.63547,
.450000,	1.07000,	2.93000,	1.08659,	2.94005,
.500000,	.98000,	3.25000,	1.02883,	3.26655,

*

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APPROXIMANT FUN	ICTION	N (1,	, 2)	
AO	A1		A2	A3
14.235672	. 9	957753	1.258183	404976.7 70760
A4	A5		A 6	
-8.145772	-408964.(29690	3987.367830	
lagl	lag2	lag3	lag4	error

.27892661	.04365735	.27899139	.28586712	.00020294	.00041111

k	EXACT		APPROXI	MATE
	REAL	IMAG	REAL	IMAG
.000000,	14.70000,	.00000,	14.23567,	.00000,
.025000,	12.20000,	-4.05000,	12.21288,	-3.64541,
.050000,	9.50000,	-4.15000,	9.56427,	-4.30539,
.087500,	7.63000,	-3.70000,	7.53586,	-3.72869,
.100000,	7.22000,	-3.47000,	7.15123,	-3.52318,
.125000,	6.61000,	-3.18000,	6.57096,	-3.17275,
.150000,	6.09000,	-2.94000,	6.12655,	-2.88115,
.175000,	5.72000,	-2.59000,	5.75592,	-2.61606,
.200000,	5.41000,	-2.29000,	5.43806,	-2.35648,
.250000,	4.92000,	-1.79000,	4.94237,	-1.82723,
.300000,	4.63000,	-1.30000,	4.62099,	-1.29981,
.350000,	4.43000,	83000,	4.44320,	81140,
.400000,	4.35000,	41000,	4.36678,	38607,
.450000,	4.35000,	03000,	4.35258,	03010,
.500000,	4.40000,	.28000,	4.37100,	.26139,
APPROXIMANT F	UNCTION	N (2, 1	.)	

A0	A1	λ2	A3
.000000	-1.905869	293633	158360.047503

A4

A6

-.059842 -159903.225805 1543.931881

A5

	error	lag4	lag3	lag2	lagl
.00060945	.00019833	.28586712	.27899139	.04365735	.27892661

k	EXAC	CT APPROXIMATE			T APPROXIMATE		MATE
	REAL	IMAG	REAL	IMAG			
.000000,	.00000,	.00000,	.00000,	.00000,			
.025000.	.00100.	03500,	00729,	03935,			
.050000,	.01500,	06500,	00727,	06574,			
.087500,	.02500,	12500,	.01171,	11930,			
.100000,	.03000,	14500,	.01725,	14295,			
.125000,	.02000,	19800,	.02145,	19500,			
.150000,	.01500,	23000,	.01512,	24731,			
.175000,	.01000,	28000,	.00030,	29389,			
.200000,	00500,	32500,	01874,	33146,			
.250000,	03500,	39700,	05219,	37922,			
.300000,	06000,	40300,	06191,	40363,			
.350000,	04500,	42500,	04310,	42338,			
.400000,	00500,	44800,	00139,	45102,			
.450000,	.06500,	49000,	.05504,	49191,			
.500000,	.14000,	55000,	.11927,	54679,			

APPROXIMANT FUNCTION N (2, 2)

A 0	Al	A2	A3
516839	-4.394286	2.744002	217699.240354
A4	23	A 6	
.165112	-219881.401614	2184.097622	

lagi	lag2	lag3	lag4	error	
.27892661	.04365735	.27899139	.28586712	.00078657	.00139602

k	EXAC	Г	APPROXIMATE	
	REAL	IMAG	REAL	IMAG
.000000,	44500,	.00000,	51684,	.00000,
.025000,	45000,	02700,	48373,	04150,
.050000,	··· 45500,	12200,	45477,	14213,
.087500,	49500,	31100,	48600,	31323,

.100000,	52000,	34800,	51147,	36418,
.125000,	58100,	42200,	57665,	45028,
.150000,	67200,	50800,	65200,	51226,
.175000,	73200,	54800,	72737,	55013,
.200000,	81000,	58200,	79435,	56754
.250000,	90500,	61000,	88350,	56578,
.300000,	89800,	55000,	90772,	55768,
.350000,	86800,	57500,	88429,	57878,
.400000,	82400,	64800,	83843,	64393,
.450000,	79000,	75000,	79150,	75384
.500000,	77000,	90300,	75798,	90251,
•	•	•	•	•

OUTPUT FILE : LAPLACE.DAT

I

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.593074	00000 . 870	0 7.625432	7.351972 -146.23	26517 7 1038	-8561. 672085
	14.23567	2	.957753	1.25818	404976. 770760
-8.145772	-40896	4.029690	3987.36	/830	
059842	.00000	0 ·	-1.905869 1543.93	2936 3:	3 158360.047503
			1010000		
.165112	51683	9 1.401614	-4.394286 2184.09	2.74400 7622	2 217699.240354
.27	892661	.04365735	5.27899139	.28586712	.00139602

```
$storage:2
$floatcalls
$debug
C----
   LAPLACE EIGENVALUE SOLUTION
С
С
  FLUTTER CALCULATION (INCOMPRESSIBLE REGIME)
C-----
С
      IMPLICIT REAL*8 (A-Z)
      INTEGER ib, n, ia, ijob, iz, ier, infer, i, j
      COMPLEX*16 eval(12), A(12, 12), B(12, 12), EIGA(12), EIGB(12), Z(12, 12)
      COMPLEX*16 WK(21, 42)
      REAL*8 u,vel,atemp1,atemp2,atemp3,atemp4
      REAL*8 finv(7,2,2),f(8,2,2),m(2,2),c(2,2),k(2,2)
      REAL*8 a0(2,2),a1(2,2),a2(2,2),a3(2,2),a4(2,2),a5(2,2),a6(2,2)
      real*8 n1(2,7)
      INTEGER y
С
C Read in constants.
      OPEN (UNIT=1, FILE='CONST6.DAT', STATUS='OLD')
      READ (1,*) ia, ib, iz, n, ijob
С
                 12 12 12 12 2
      CLOSE (UNIT=1)
С
C open required files
      OPEN (UNIT=4, FILE='EIGEN.OUT', STATUS='NEW')
C
C required constants
             u=50.0
             wrat=0.2
             r=0.5
             ah=-0.5
С
             xtor=0.25
             semic=1.0
             wtor=1.0
             wben=wrat*wtor
С
C required constants
С
      OPEN (UNIT=10, FILE='LAPLACE.DAT', STATUS='OLD')
          do 25 i=1,2
      READ (10,*) n1(i,1),n1(i,2),n1(i,3),n1(i,4),n1(i,5),
     1
                  nl(i,6),nl(i,7)
 25
          continue
      READ (10,*) b1,b2,b3,b4,err
      CLOSE (UNIT=10)
С
C titles and initial parameter listing to file
             write(4, 32)
             write(4,121)
             write(4,121)
             write(4, 33)
             write(4,121)
```

```
write(4,34) u,r,ah,wrat
              write(4,121)
              write(4,121)
 32
              format(6x, ' LAPLACE METHOD
                                                INCOMPRESSIBLE (THEODORSEN)')
              format(6x,'
 33
                                           r(alp)
                              u
                                                        ah
                                                                wben/wtor')
 34
              format(6x, f7.1, 4f10.1)
 121
              format('
                           1)
С
C FOUR LAG TERMS
С
C PADE APPROXIMANT
                       THEODORSEN FUNCTION
                                                 C(k)
                                                        AND
                                                              C(k) * ik
С
C term Q(1,1): C(k) * ik
С
              a0(1,1) = n1(2,1)
              a1(1,1) = n1(2,2)
              a2(1,1) = n1(2,3)
              a3(1,1) = n1(2,4)
              a4(1,1) = n1(2,5)
              a5(1,1) = n1(2,6)
              a6(1,1) = n1(2,7)
С
C term Q(1,2): C(k) * ik * (0.5-ah) + C(k)
С
              a0(1,2) = n1(2,1) * (0.5-ah) + n1(1,1)
              al(1,2) = nl(2,2) * (0.5-ah) + nl(1,2)
              a2(1,2) = n1(2,3) * (0.5-ah) + n1(1,3)
              a3(1,2) = n1(2,4) * (0.5-ah) + n1(1,4)
               a4(1,2) = n1(2,5) * (0.5-ah) + n1(1,5)
              a5(1,2) = n1(2,6) * (0.5-ah) + n1(1,6)
              a6(1,2) = n1(2,7) * (0.5-ah) + n1(1,7)
С
C term Q(2,1): C(k)*ik*(-(ah+0.5))
C
              a0(2,1) = n1(2,1) * (-0.5-ah)
              al(2,1) = nl(2,2) * (-0.5-ah)
              a2(2,1) = n1(2,3) * (-0.5-ah)
              a3(2,1) = n1(2,4) * (-0.5-ah)
              a4(2,1) = n1(2,5) * (-0.5-ah)
              a5(2,1) = n1(2,6) * (-0.5-ah)
              a6(2,1) = n1(2,7) * (-0.5-ah)
С
C term Q(2,2): C(k)*ik*(-(ah+0.5)*(0.5-ah))-(ah+0.5)*C(k)
С
              a0(2,2) = n1(2,1) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,1)
              a1(2,2) = n1(2,2) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,2)
              a2(2,2) = n1(2,3) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,3)
              a3(2,2) = n1(2,4) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,4)
              a4(2,2) = n1(2,5) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,5)
              a5(2,2) = n1(2,6) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,6)
              a6(2,2) = n1(2,7) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,7)
С
С
```

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С

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```
b1234=b1+b2+b3+b4
              b123=b1+b2+b3
              b124 = b1 + b2 + b4
              b134 = b1 + b3 + b4
              b234 = b2 + b3 + b4
              bt123=b1*b2*b3
              bt124=b1*b2*b4
              bt134=b1*b3*b4
              bt234=b2*b3*b4
              bq1234=b1*b2*b3*b4
              bt1234=b1*b2*b3+b1*b2*b4+b1*b3*b4+b2*b3*b4
              bd1234=b1*b2+b1*b3+b2*b3+b1*b4+b2*b4+b3*b4
              bd123=b1*b2+b1*b3+b2*b3
              bd124=b1*b2+b1*b4+b2*b4
              bd134=b1*b3+b1*b4+b3*b4
              bd234=b2*b3+b3*b4+b2*b4
С
C iterate velocity
              vel=0.0
              step =0.5
 60
              vel=vel+step
              if(vel/(semic*wtor).gt.5.0) goto 120
              wf=vel/semic
              xvar=2*wf**2/u
С
C formation of matrices :
                             [M], [C], [K], [D] and [E]
              m(1,1) = 1 + 1/u
              m(1,2) = xtor-ah/u
              m(2,1) = xtor-ah/u
              m(2,2) = r**2+(0.125+ah**2)/u
              c(1,1)=0.0
              c(1,2) = wf/u
              c(2,1)=0.0
              c(2,2) = (0.5-ah) * wf/u
              k(1,1)=wrat**2*wtor**2
              k(1,2)=0.0
              k(2,1)=0.0
              k(2,2)=r**2*wtor**2
С
C formation of polynomial coefficients :
                                              [F7], [F6], [F5], [F4]
С
                                              [F3], [F2], [F1]
С
С
    [F7]s6+[F6]s5+ \dots + [F2]s+[F1] = 0
           do 200 i=1,2
           do 100 j=1,2
С
              calculation of [F7]
              f(7, i, j) = m(i, j) + xvar * (1/wf) * * 2 * a 2(i, j)
C
              calculation of [F6]
              f(6,i,j)=c(i,j)+wf*b1234*m(i,j)+xvar*(1/wf)*(a1(i,j)+
     1
              b1234*a2(i,j))
С
              calculation of [F5]
              f(5,i,j)=k(i,j)+wf*b1234*c(i,j)+wf**2*bd1234*m(i,j)+
              xvar * (a0(i,j)+b1234*a1(i,j)+bd1234*a2(i,j)+a3(i,j)+
     1
```

```
1
                             a4(i,j)+a5(i,j)+a6(i,j)
С
                             calculation of [F4]
                              f(4,i,j)=wf*b1234*k(i,j)+wf**2*bd1234*c(i,j)+wf**3*bt1234*
           1
                             m(i,j) + xvar + wf + (b1234 + a0(i,j) + bd1234 + a1(i,j) + bt1234 + at(i,j) + bt1234 + 
           1
                              a2(i,j)+b234*a3(i,j)+b134*a4(i,j)+b124*a5(i,j)+
           1
                             b123*a6(i,j))
С
                              calculation of [F3]
                              f(3,i,j)=wf**2*bd1234*k(i,j)+wf**3*bt1234*c(i,j)+wf**4*
                             bg1234*m(i,j)+xvar*wf**2*(bd1234*a0(i,j)+bt1234*a1(i,j)+
           1
                              bq1234*a2(i,j)+bd234*a3(i,j)+bd134*a4(i,j)+bd124*
           1
           1
                              a5(i,j)+bd123*a6(i,j))
С
                              calculation of [F2]
                              f(2,i,j) = wf * *3 * bt 1234 * k(i,j) + wf * *4 * bg 1234 * c(i,j) +
           1
                              xvar*wf**3*(bt1234*a0(i,j)+bq1234*a1(i,j)+bt234*a3(i,j)+
                              bt134*a4(i,j)+bt124*a5(i,j)+bt123*a6(i,j))
           1
С
                              calculation of [F1]
                              f(1,i,j)=wf**4*bq1234*k(i,j)+xvar*wf**4*bq1234*a0(i,j)
  100
                       continue
                       continue
  200
С
С
                              calculation of inverse
                                                                                      (inv[F8])
                              f7det=f(7,1,1)*f(7,2,2)-f(7,1,2)*f(7,2,1)
                              if(f7det.eq.0.0) goto 210
                              goto 220
  210
                              write(4,215)
                              format(' DETERMINANT EQUALS ZERO - INVESTIGATE INPUT '.
  215
           1
                              'CONSTANT ... ah ')
                              goto 120
                              finv(7,1,1) = f(7,2,2)/f7det
  220
                              f(8,1,1) = finv(7,1,1)
                              finv(7,2,2) = f(7,1,1)/f7det
                              f(8,2,2) = finv(7,2,2)
                              finv(7,1,2) = -f(7,1,2)/f7det
                              f(8,1,2) = finv(7,1,2)
                              finv(7,2,1) = -f(7,2,1)/f7det
                              f(8,2,1)=finv(7,2,1)
С
C formation of matrices [A] and [B] to
C solve eigenvalue problem given by [A]{x} = lambda [B]{x}
                       do 400 i=1,n
                       do 300 j=1,n
                              a(i,j) = dcmplx(0.0,0.0)
                              if(i.eq.j+2) a(i,j)=dcmplx(1.0,0.0)
                              b(i,j) = dcmplx(0.0,0.0)
                              if(i.eq.j) b(i,j) = dcmplx(1.0,0.0)
  300
                       continue
  400
                       continue
                       do 500 i=1,n/2
                              y=n/2+1-i
                              atemp1=-(finv(7,1,1)*f(y,1,1)+finv(7,1,2)*f(y,2,1))
                              a(1,2*i-1)=dcmplx(atemp1,0.0)
                              atemp2 = -(finv(7,1,1) * f(y,1,2) + finv(7,1,2) * f(y,2,2))
                              a(1,2*i)=dcmplx(atemp2,0.0)
                              atemp3 = -(finv(7,2,1) * f(y,1,1) + finv(7,2,2) * f(y,2,1))
```

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A - 95
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a(2,2*i-1) = dcmplx(atemp3,0.0)
              atemp4=-(finv(7,2,1)*f(y,1,2)+finv(7,2,2)*f(y,2,2))
              a(2,2*i) = dcmplx(atemp4,0.6)
 500
          continue
С
 748
              format(6x, 2f24.6)
              format(6x, 28e8.1)
 750
 749
              format(6x, i3)
              format(' ')
 760
С
C subrountine for eigenvalue calculation
C abar : real part of eval()
C bbar : imaginary part of eval()
      CALL EIGZC(A, IA, B, IB, N, IJOB, EIGA, EIGB, Z, IZ, WK, INFER, IER)
          do 5 I=1,n
 5
             EVAL(I) = EIGA(I) / EIGB(I)
C printout of velocity and associated eigenvalues
             write(4,23) vel/(semic*wtor)
          do 6 i=1,n
             write(4,22) eval(i)
 6
          continue
              write(4,21) wk(1,1)
              write(4,121)
              format(6x,' perfomance index is', 2f15.5)
 21
              format(6x, ' ND VELOCITY ', f6.3)
 23
 22
              format(6x,8f17.5)
 28
              format(6x, f17.5, ', ', f17.5, ', ', f17.5)
С
С
C changing velocity step size
С
               if(vel/(semic*wtor).ge.3.5) step=0.1
С
               if(vel/(semic*wtor).ge.5.0) step=0.1
С
               if(vel/(semic*wtor).ge.6.0) step=0.5
               if(vel/(semic*wtor).ge.6.0) goto 120
С
 700
               goto 60
 120
      CLOSE (UNIT=4)
 800
      END
```

INPUT FILE : LAPLACE.DAT

1

 .99828226
 -.00003562
 .00000474
 -.04120375

 -.16289333
 -.22908745
 -.06459371
 -.00001355

 -.00001355
 .50037475
 -.00003629
 .00061076

 .01380489
 .05848218
 .04963449
 .000011782

INPUT FILE : CONST6.DAT

12 12 12 12 2

OUTPUT FILE : EIGEN.OUT

i

LAPLACE	METHOD	INCOMP	RESSI	BLE	(THEODORSEN)	1
u	r(al	.p)	ah	wł	pen/wtor	
50.0	.5		5		.2	
ND VELO	CITY .50	00				
	01437		1.	14595	5	
	01437		-1.	10021		
	00637			1992:		
	38147			00000		
	38256		•	00000)	
	12656		•	00000)	
	12798		•	00000)	
	04200		•	00000		
	04214		•	00000		
	00772		•	00000	<i>ו</i> ר	
perfoma	ance index	ris	•	1.252	299	.00000
ND VELOO	CITY 1.00	0				
	02966		1.	13480	5	
	02966		-1.	13486	5	
	76212		•	00000)	
	76511		•	00000	2	
	01482		•	20360		
	- 24979			20360		
	-,24879		•	00000		
	08316			00000	5	
	08429		•	00000	5	
	01544		•	00000	0	
	01543		•	00000	כ	
perfoma	ance index	(is		1.50	717	.00000
ND VELO	CITY 1.50	00	_		_	
	04649		-1.	11520	5	
	04649		1.	11520		
	-1.14241		•		5	
	-1.14707		•	2098	1	
	02480			2098	4	
	36693			0000	0	
	38393		•	0000	D	
	12267		•	0000	0	
	12643		•	0000	D	
	02316		•	0000	D	
	02313	• _	•	0000	0	
perfoma	ance index	(1S		3.53	341	.00000

ND VELOCITY 2.000		
06529	1.08538	
-1.52250	.00000	
-1.53023	.00000	
06529	-1.08538	
51190	.00000	
48021	.00000	
03664	.21879	
03664	21879	
- .15957	.00000	
16857	.00000	
03088	.00000	
03080	.00000	
perfomance index is	17.25293	.00000
ND VELOCITY 2.500		
-1.90251	.00000	
-1.91278	.00000	
08644	-1.04289	
08644	1.04289	
63988	.00000	
58718	.00000	
05141	-,23195	
05141	23195	
19235	.00000	
-21072	.00000	
03860	.00000	
- 03843	00000	
perfomance index is	1.84973	.00000
ND VELOCITY 3 000		
-2 29251	00000	
-2.20231	.00000	
-2.29004 1100A	- 98348	
- 11024	20340	
- 76795	.98348	
- 68609	.00000	
08009	.00000	
07108	- 25240	
- 25286	25240	
- 21216	.00000	
21816	.00000	
04632	.00000	
	.00000	
periomance index is	3.86263	.00000
ND VELOCITY 3.500		
-2.66251	.00000	
-2.67790	.00000	
13678	.89843	
13678	89843	
89583	.00000	
77455	.00000	
09940	.28811	

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09940	28811	
29500	.00000	
23173	.00000	
05404	.00000	
05348	.00000	
perfomance index is	2.67030	.00000
ND VELOCITY 4.000		
-3.04253	.00000	
-3.06045	.00000	
16495	76095	
16495	.76095	
-1.02380	.00000	
84867	.00000	
14323	36783	
14323	.36783	
33715	.00000	
·~·22548	.00000	
06176	.00000	
06081	.00000	
perfomance index is	3.78368	.00000
ND VELOCITY 4.500		
-3.42257	.00000	
-3.44301	.00000	
-1.15178	.00000	
01016	54976	
39043	51943	
90109	.00000	
01016	.54976	
39043	.51943	
37929	.00000	
20035	.00000	
06948	.00000	
06791	.00000	
perfomance index is	3.19514	.00000
ND VELOCITY 5.000		
-3.80262	.00000	
-3.82557	.00000	
.11382	52557	
.11382	.52557	
-1.27975	.00000	
62744	48581	
62744	.48581	
9151.8	.00000	
42143	.00000	
17266	.00000	
07461	.00000	
07720	.00000	
perfomance index is	2.91782	.00000

```
$storage:2
$floatcalls
$debug
C---
С
   LAPLACE EIGENVALUE SOLUTION
C
   FLUTTER CALCULATION (TRANSONIC REGIME)
C-----
      IMPLICIT REAL*8 (A-Z)
      INTEGER ib, n, ia, ijob, iz, ier, infer, i, j
      COMPLEX*16 eval(12), A(12,12), B(12,12), EIGA(12), EIGB(12), Z(12,12)
      COMPLEX*16 WK(21,42)
      REAL*8 u,vel,atemp1,atemp2,atemp3,atemp4
      REAL*8 finv(7,2,2),f(8,2,2),m(2,2),c(2,2),k(2,2),d(2,2),e(2,2)
      REAL*8 a0(2,2),a1(2,2),a2(2,2),a3(2,2),a4(2,2),a5(2,2),a6(2,2)
      INTEGER y
С
C Read in constants.
      OPEN (UNIT=1, FILE='CONST6.DAT', STATUS='OLD')
      READ (1,*) ia, ib, iz, n, ijob
С
                  12 12 12 12 2
      CLOSE (UNIT=1)
C
C open required files
      OPEN (UNIT=4, FILE='EIGEN.OUT', STATUS='NEW')
С
C required constants
              u=50.0
              r = 0.5
              ah = -0.5
              wrat=0.2
С
              xtor=0.25
              semic=1.0
              wtor=1.0
              wben=wrat*wt >r
              pi=4.0*atan(1.0)
С
C required constants (approximation)
      OPEN (UNIT=10, FILE='LAPLACE.DAT', STATUS='OLD')
           do 50 i=1,2
           do 25 j=1,2
      READ (10,*) a0(i,j),a1(i,j),a2(i,j),a3(i,j),a4(i,j),
      1
                   a5(i,j),a6(i,j)
 25
           continue
 50
           continue
      READ (10,*) b1,b2,b3,b4,err
      CLOSE (UNIT=10)
C titles and initial parameter listing to file
              write(4, 32)
              write(4,121)
              write(4,121)
              write(4,33)
              write(4,121)
```

write(4,34) u,r,ah,wrat write(4,121) write(4,121) 32 format(6X,' LAPLACE METHOD TRANSONIC M=0.85') format(6X,' r(alp) ah wben/wtor') 33 u format(6X, f7.1, 4f10.1) 34 **''**) format(' 121 С С C term Q(1,1): N(1,1) = Clh/2С a0(1,1) =a0(1,1)/2.0al(1,1) =a1(1,1)/2.0a2(1,1) =a2(1,1)/2.0a3(1,1) = $a_3(1,1)/2.0$ a4(1,1) =a4(1,1)/2.0a5(1,1) =a5(1,1)/2.0a6(1,1) =a6(1,1)/2.0С C term Q(1,2): N(1,2) = CltorC a0(1,2) =a0(1,2) a1(1,2) =a1(1,2)a2(1,2) =a2(1,2) $a_3(1,2) =$ a3(1,2) a4(1,2) =a4(1,2) a5(1,2) =a5(1,2)a6(1,2) =a6(1,2)С C term Q(2,1): N(2,1) = Cmh*(-1)С a0(2,1) =a0(2,1) * (-1.0)al(2,1) =a1(2,1) * (-1.0)a2(2,1) =a2(2,1) * (-1.0)a3(2,1) =a3(2,1) * (-1.0)a4(2,1) =a4(2,1) * (-1.0)a5(2,1) =a5(2,1) * (-1.0)a6(2,1) =a6(2,1) * (-1.0)С C term Q(2,2): N(2,2) = Cmtor*(-2)С a0(2,2) =a0(2,2)*(-2.0) a1(2,2) =a1(2,2)*(-2.0)a2(2,2) =a2(2,2) * (-2.0)a3(2,2) =a3(2,2)*(-2.0)a4(2,2) =a4(2,2)*(-2.0)a5(2,2) =a5(2,2)*(-2.0) a6(2,2) =a6(2,2) * (-2.0)С С С b1234=b1+b2+b3+b4 b123=b1+b2+b3 b124=b1+b2+b4 A - 102

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b134=b1+b3+b4
             b234=b2+b3+b4
             bt123=b1*b2*b3
             bt124=b1*b2*b4
             bt134=b1*b3*b4
             bt234=b2*b3*b4
             bq1234=b1*b2*b3*b4
             bt1234=b1*b2*b3+b1*b2*b4+b1*b3*b4+b2*b3*b4
             bd1234=b1*b2+b1*b3+b2*b3+b1*b4+b2*b4+b3*b4
             bd123=b1*b2+b1*b3+b2*b3
             bd124=b1*b2+b1*b4+b2*b4
             bd134=b1*b3+b1*b4+b3*b4
             bd234=b2*b3+b3*b4+b2*b4
C
C iterate velocity
             vel=0.0
             step =0.5
 60
             vel=vel+step
              if(vel/(semic*wtor).gt.4.0) goto 120
             wf=vel/semic
             xvar=(1/pi)*wf**2/u
C
C formation of matrices :
                            [M], [C], [K], [D] and [E]
             m(1,1) = 1.0
             m(1,2) = xtor
             m(2,1) = xtor
             m(2,2) = r * * 2
              c(1,1)=0.0
              C(1,2)=0.0
              c(2,1)=0.0
              c(2,2)=0.0
             k(1,1) = wben * * 2
             k(1,2)=0.0
             k(2,1)=0.0
             k(2,2)=r**2*wtor**2
С
C formation of polynomial coefficients :
                                            [F7], [F6], [F5], [F4]
С
                                            [F3], [F2], [F1]
С
С
    [F7]s6+[F6]s5+ \dots + [F2]s+[F1] = 0
          do 200 i=1,2
          do 100 j=1,2
С
              calculation of [F7]
              f(7,i,j)=m(i,j)+xvar*(1/wf)**2*a2(i,j)
С
              calculation of [F6]
              f(6,i,j)=c(i,j)+wf*bl234*m(i,j)+xvar*(1/wf)*(al(i,j)+
     1
              b1234*a2(i,j))
С
              calculation of [F5]
              f(5,i,j)=k(i,j)+wf*bl234*c(i,j)+wf**2*bdl234*m(i,j)+
     1
              xvar*(a0(i,j)+b1234*a1(i,j)+bd1234*a2(i,j)+a3(i,j)+
     1
              a4(i,j)+a5(i,j)+a6(i,j)
С
              calculation of [F4]
              f(4,i,j)=wf*b1234*k(i,j)+wf**2*bd1234*c(i,j)+wf**3*bt1234*
     1
              m(i,j)+xvar*wf*(b1234*a0(i,j)+bd1234*a1(i,j)+bt1234*
```

0	1 1	a2(i,j)+b234*a3(i,j)+b134*a4(i,j)+b124*a5(i,j)+ b123*a6(i,j))
C	1	calculation of [F3] f(3,i,j)=wf**2*bd1234*k(i,j)+wf**3*bt1234*c(i,j)+wf**4* bq1234*m(i,j)+xvar*wf**2*(bd1234*a0(i,j)+bt1234*a1(i,j)+
с	1 1	bq1234*a2(1,j)+bd234*a3(1,j)+bd134*a4(1,j)+bd124*a5(i,j)+bd123*a6(i,j)) calculation of [F?]
-	1 1	f(2,1,j)=wf**3*bt1234*k(1,j)+wf**4*bq1234*c(1,j)+ xvar*wf**3*(bt1234*a0(i,j)+bq1234*a1(i,j)+bt234*a3(i,j)+ bt134*a4(i,j)+bt124*a5(i,j)+bt123*a6(i,j))
C		f(1,i,i) = wf**4*bg1234*k(i,i)+xvar*wf**4*bg1234*a0(i,i)
100	С	ontinue
200	С	ontinue
C C		calculation of inverse (inv[F8])
0		f7det=f(7,1,1) *f(7,2,2)-f(7,1,2) *f(7,2,1) if(f7det.eq.0.0) goto 210
210		goto 220 write(4 215)
215	1	format(' DETERMINANT EQUALS ZERO - INVESTIGATE INPUT ', 'CONSTANT ah ')
220		finv(7,1,1)=f(7,2,2)/f7det
		f(8,1,1) = finv(7,1,1)
		finv(7,2,2)=f(7,1,1)/f7det
		I(8,2,2) = IINV(7,2,2) finv(7,1,2) = - f(7,1,2) / f7det
		f(8,1,2) = finv(7,1,2)
		finv(7,2,1) = -f(7,2,1) / f7det
0		f(8,2,1)=finv(7,2,1)
C for	mation	of matrices [1] and [B] to
C sol	ve eig	envalue problem given by $[A] \{x\} = lambda [B] \{x\}$
	d	o 400 i=1,n
	d	0 300 j=1, n
		a(1,j)=acmpix(0.0,0.0) if(i.eq.i+2) $a(i,i)=dcmpix(1,0,0,0)$
		b(i,j)=dcmplx(0.0,0.0)
		if(i.eq.j) b(i,j) = dcmplx(1.0,0.0)
300	C	ontinue
400	C C	ontinue o 500 j=1 p/2
	u	v=n/2+1-i
		atempl=-(finv(7,1,1)*f(y,1,1)+finv(7,1,2)*f(y,2,1))
		a(1,2*i-1) = dcmplx(atemp1,0.0)
		atemp2=-(finv(7,1,1)*f(y,1,2)+finv(7,1,2)*f(y,2,2)) (1,2*i)-domply(atomp2,0,0)
		a(1,2*1) - a(mpix(a(emp2,0.0))) atemp3 = -(finy(7,2,1) * f(y,1,1) + finy(7,2,2) * f(y,2,1))
		a(2,2*i-1)=dcmplx(atemp3,0.0)
		atemp4=-(finv(7,2,1)*f(y,1,2)+finv(7,2,2)*f(y,2,2))
500	~	a(2,2*1)=dcmplx(atemp4,0.0)
500	C	

```
С
C subrountine for eigenvalue calculation
C abar : real part of eval( )
C bbar : imaginary part of eval()
      CALL EIGZC(A, IA, B, IB, N, IJOB, EIGA, EIGB, Z, IZ, WK, INFER, IER)
          do 5 I=1,n
 5
              EVAL(I) = EIGA(I) / EIGB(I)
С
C printout of velocity and associated eigenvalues
              write(4,24)
             write(4,23) vel/(semic*wtor)
          do 6 i=1,n
              write(4,22) eval(i)
 6
          continue
              write(4,21) wk(1,1)
 21
              format(6x,' perfomance index is', 2f15.5)
              format(6x, ' ND VELOCITY ', f6.3)
 23
 22
              format(6x,8f17.5)
              format(' ')
 24
С
C changing velocity step size
               if(vel/(semic*wtor).ge.3.0) step=0.1
С
С
               if(vel/(semic*wtor).ge.3.5) step=0.5
С
               if(vel/(semic*wtor).ge.4.5) step=0.5
               if(vel/(semic*wtor).ge.5.0) step=0.5
С
 700
               goto 60
      CLOSE (UNIT=4)
 120
 800
      END
```

INPUT FILE : LAPLACE.DAT

į.

351972 -.265177 -8561.672085 -146.231038 .000000 7.351972 8707.625432 .593074 14.235672 .957753 -8.145772 -408964.029690 3987.36783 **1.258183** 404976.770760 3987.367830 .000000 -1.905869 **-.293633** 158360.047503 -.059842 -159903.225805 1543.931881 -.516839 -4.394286 **2.744002 217699.24**0354 .165112 -219881.401614 2184.097622

.27892661 .04365735 .27899139 .28586712 .00139602

INPUT FILE : CONST6.DAT

12 12 12 12 2

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1

LAPLACE	METHOD	TRANSONIC	M=0.85	
u	r(alp)	ah	wben/wtor	
50.0	• 5	 5	.2	
ND VELOC	ITY .500			
	08315	1.276	562	
	08315	-1.276	562	
	00511	201	L85	
	- 10727	.201	185	
	- 19737	017	/9/	
	12884	- 031	101	
	12884	.031	121	
	11139	014	159	
	11139	.014	159	
	02172	.000	019	
	02172	000	019	
perfoma	nce index is	1.3	34192	.00000
ND VELOC	ITY 1.000			
	13682	-1.222	225	
	13682	1.222	225	
	01143	.21	140	
	01143	21	L40	
	47016	.000	000	
	41480	.000	000	
	- 24390	080	159	
	19773	- 03	599	
	19773	039	599	
	04279	.001	145	
	04279	001	L45	
perfoma	nce index is	1.9	91492	.00000
ND VELOC	ITY 1.500			
	14304	-1.172	204	
	14304	1.172	204	
	91862	.000	000	
	60947	.000	000	
	01825	.224	90	
	01825	224	190	
	30430	- 132	(J) 25	
	27054	134	200 787	
	27054	057	787	
	06248	.004	60	
	06248	004	60	
perfoma	nce index is	3,6	55296	00000

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4 , **4**

ND VELOCITY 2.000		
13090	-1.16269	
-1.42867	.00000	
13090	1.16269	
81080	.00000	
02269	.24257	
02269	24257	
46358	.18321	
46358	18321	
3333L	.07634	
33331	- 01006	
- 08006	01006	
perfomance index is	3.54698	.00000
E 1 1 1 1		
ND VELOCITY 2.500		
-1.93877	.00000	
13300	-1.18652	
13300	1.18652	
-1.01285	.00000	
02228	.26434	
02228	26434	
5/2/8	.23294	
5/2/8	23294	
- 38902	.08819	
30902		
- 09496	01771	
perfomance index is	4.61556	.00000
-		
ND VELOCITY 3.000		
-2.43536	.00000	
15539	1.22983	
15539	-1.22983	
-1.21510	.00000	
68232	28189	
08232	.28189	
	28849	
- 44104	.20049	
- 44104	- 09158	
- 10696	02704	
- 10696	- 02704	
perfomance index is	3.86013	.00000
•		
ND VELOCITY 3.500		
-2.92058	.00000	
19525	1.28757	
-1.41743	.00000	
19525	-1.28757	
79221	33037	
79221	.33037	
.00221	.31152	

.00221 49236 49236 11637 11637 perfomance index is	31152 .08603 08603 .03744 03744 5.60378	.00000
ND VELOCITY 4.000		
-3.39765	.00000	
-1.61980	.00000	
24794	-1.35846	
24794	1.35846	
90240	.37857	
90240	37857	
.02714	32936	
.02714	.32936	
54475	07102	
54475	.07102	
12386	.04847	
12386	04847	
perfomance index is	3.19862	.00000

```
C ----
     LAPLACE EIGENVALUE SOLUTION
С
С
     IMPLEMENTATION OF ACTIVE CONTROL
С
     FLUTTER CALCULATION (INCOMPRESSIBLE REGIME)
С
         С
С
     OF'ERATION
С
     REQUIRE: 1) INPUT OF ICOUNT, N, REQMIN, START(I), STEP(I)
С
              2) USER SPECIFIED FUNCTION SUBPROGRAM (DOUBLE PRECISION
С
                 FUNCTION FN(X) - MINIMIZATION FUNCTION)
С
С
     DOUBLE PRECISION START(20), STEP(20), XMIN(20),
    1XSEC(20), YNEWLO, YSEC, REQMIN
С
     OPEN(UNIT=6, FILE='DATA.OUT', STATUS='NEW')
С
С
С
С
       С
                                                       *
С
           ICOUNT, REQMIN, START(N), STEP(N)
       *
                                                       *
С
       *
                                                       *
С
         velndnc : no control non-dimensional velocity
       *
                                                       *
С
       *
           step : velocity step size
                                                       *
С
           u : airfoil-air mass ratio
                                                       *
       *
С
       *
         wrat : frequency ratio (wben/wtor)
                                                       *
С
         r : non-dimensional distance
                                                       *
С
       *
          ah : non-dimensional distance
                                                       *
С
                                                       *
С
       * OUTPUT FILE : OPTIMAL.OUT
                                                       *
С
С
       ****
С
C
С
C LET N=4 AT ALL TIMES, SIMPLY SET START(I)=0.0 AND STEP(I)=0.0
С
С
     ICOUNT=8
     N=4
     REQMIN=0.00000001
     START(1) = 0.0
     START(2) = -0.3
     START(3) = 4.0
     START(4) = 3.6
     STEP(1) = 0.1
     STEP(2) = -0.1
     STEP(3) = 0.1
     STEP(4) = 0.1
С
С
     DO 60 I=1,N
     XMIN(I) = 0.D0
```

```
XSEC(I) = 0.D0
 60
      CONTINUE
      YNEWLO=0.D0
      YSEC=0.D0
С
С
                CALL NELDER-MEAD SUBROUTINE
С
      CALL NELMIN(N, START, XMIN, XSEC, YNEWLO, YSEC,
     1REQMIN, STEP, ICOUNT)
С
                OUTPUT FROM PROGRAM
С
С
      WRITE(6,64)
      WRITE(6,65) ICOUNT
      WRITE(6,75)
      WRITE(6,77)
      DO 79 I=1,N
 79
      WRITE(6,80) I,XSEC(I),XMIN(I)
      WRITE(6,82)
      WRITE(6,83)
      WRITE(6,84) YSEC, YNEWLO
С
 64
      FORMAT(1H //,1H ,42H SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION)
 65
      FORMAT(1H //,1H ,15,12H TRIALS USED/)
 75
      FORMAT(1H ,21X,9HESTIMATES/)
 77
      FORMAT(1H, 9HPARAMETER, 7X, 12HNEXT-TO-BEST, 8X,
     14HBEST/)
 80
      FORMAT(1H , 15, 2F20.7)
 82
      FORMAT(1H //,1H ,6X,15HFUNCTION VALUES/)
      FORMAT(1H ,5X,13H NEXT-TO-BEST,8X,4HBEST/)
 83
 84
      FORMAT(1H ,2F15.9)
      STOP
      END
С
С
              NELDER-MEAD SUBROUTINE
С
      SUBROUTINE NELMIN(N, START, XMIN, XSEC, YNEWLO, YSEC,
     1REQMIN, STEP, ICOUNT)
      DOUBLE PRECISION START(N), STEP(N), XMIN(N),
     1XSEC(N), YNEWLO, YSEC, REQMIN, P(20, 21), PSTAR(20),
     2P2STAR(20), PBAR(20), Y(20), DN, Z, YLO, RCOEFF,
     3YSTAR, ECOEFF, Y2STAR, CCOEFF, FN, DABIT, DCHK,
     4COORD1, COORD2
      DATA RCOEFF/1.0D0/, ECOEFF/2.0D0/, CCOEFF/0.5D0/
      KCOUNT=ICOUNT
      ICOUNT=0
С
      IF(REQMIN.LE.O.DO) ICOUNT=ICOUNT-1
      IF(N.LE.0) ICOUNT=ICOUNT-10
      IF(N.GT.20) ICOUNT=ICOUNT-10
      IF(ICOUNT.LT.0) RETURN
C
      DABIT=2.04607D-35
      BIGNUM=1.0D38
```

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A - 111
```

```
KONVGE=5
      XN=FLOAT(N)
      DN=DFLOAT(N)
      NN=N+1
С
      CONSTRUCTION OF INITIAL SIMPLEX
С
С
 1001 DO 1 I=1,N
    1 P(I,NN)=START(I)
      Y(NN) = FN(START)
      ICOUNT=ICOUNT+1
      DO 2 J=1,N
      DCHK=START(J)
      START(J) = DCHK + STEP(J)
      DO 3 I=1,N
    3 P(I,J)=START(I)
      Y(J) = FN(START)
      ICOUNT=ICOUNT+1
    2 START (J) = DCHK
С
С
       SIMPLEX CONSTRUCTION COMPLETE
С
С
       FIND HIGHEST AND LOWEST Y VALUES
С
       YNEWLO (Y(IHI)) INDICATES THE VERTEX OF
С
       THE SIMPLEX TO BE REPLACED
С
 1000 YLO=Y(1)
      YNEWLO=YLO
      ILO=1
      IHI=1
      DO 5 I=2,NN
      IF(Y(I).GE.YLO) GOTO 4
      YLO=Y(I)
      ILO=I
    4 IF(Y(I).LE.YNEWLO) GOTO 5
      YNEWLO=Y(I)
      IHI=I
    5 CONTINUE
С
С
        PERFORM CONVERGENCE CHECKS ON FUNCTION
С
      DCHK=(YNEWLO+DABIT)/(YLO+DABIT)-1.DO
      IF (DABS (DCHK).LT.REQMIN) GOTO 900
С
      KONVGE=KONVGE-1
      IF(KONVGE.NE.0) GOTO 2020
      KONVGE=5
С
```

CHECK CONVERGENCE OF COORDINATES ONLY EVERY 5 SIMPLEXES

```
DO 2015 I=1,N
COORD1=P(I,1)
COORD2=COORD1
```

14 14 С

С

С

```
DO 2010 J=2,NN
      IF(P(I,J).GE.COORD1) GOTO 2005
      COORD1=P(I,J)
 2005 IF(P(I,J).LE.COORD2) GOTO 2010
      COORD2 = P(I, J)
 2010 CONTINUE
      DCHK=(COORD2+DABIT)/(COORD1+DABIT)-1.D0
      IF(DABS(DCHK).GT.REQMIN) GOTO 2020
 2015 CONTINUE
      GOTO 900
 2020 IF(ICOUNT.GE.KCOUNT) GOTO 900
С
С
         CALCULATE PBAR, THE CENTROID OF THE
С
         SIMPLEX VERTICES EXCEPTING THAT WITH
С
         Y VALUE YNEWLO
С
      DO 7 I=1,N
      Z=0.0D0
      DO 6 J=1,NN
    6 Z=Z+P(I,J)
      Z=Z-P(I,IHI)
    7 PBAR(I) = Z/DN
С
С
         REFLECTION THROUGH THE CENTROID
С
      DO 8 I=1,N
    8 PSTAR(I)=(1.0D0+RCOEFF)*PBAR(I)-RCOEFF*P(I,IHI)
      YSTAR=FN(PSTAR)
      ICOUNT=ICOUNT+1
      IF(YSTAR.GE.YLO) GOTO 12
      IF(ICOUNT.GE.KCOUNT) GOTO 19
С
С
         SUCCESSFUL REFLECTION, SO EXTENSION
C
      DO 9 I=1,N
    9 P2STAR(I)=ECOEFF*PSTAR(I)+(1.0D0-ECOEFF)*PBAR(I)
      Y2STAR=FN(P2STAR)
      ICOUNT=ICOUNT+1
С
С
         RETAIN EXTENSION OR CONTRACTION
С
      IF(Y2STAR.GE.YSTAR) GOTO 19
   10 DO 11 I=1,N
   11 P(I,IHI)=P2STAR(I)
      Y(IHI)=Y2STAR
      GOTO 1000
С
С
         NO EXTENSION
С
   12 L=0
      DO 13 I=1,NN
      IF(Y(I).GT.YSTAR) L=L+1
   13 CONTINUE
      IF(L.GT.1) GOTO 19
```

```
A - 113
```

```
IF(L.EQ.0) GOTO 15
С
С
         CONTRACTION ON THE REFLECTION SIDE OF THE
С
         CENTROID
С
      DO 14 I=1,N
   14 P(I,IHI)=PSTAR(I)
      Y(IHI)=YSTAR
С
С
         CONTRACTION ON THE Y(IHI) SIDE OF THE CENTROID
С
   15 IF(ICOUNT.GE.KCOUNT) GOTO 900
      DO 16 I=1,N
   16 P2STAR(I)=CCOEFF*P(I,IHI)+(1.0D0-CCOEFF)*PBAR(I)
      Y2STAR=FN(P2STAR)
      ICOUNT=ICOUNT+1
      IF(Y2STAR.LT.Y(IHI)) GOTO 10
С
С
         CONTRACT THE WHOLE SIMPLEX
С
      DO 18 J=1,NN
      DO 17 I=1,N
      P(I,J) = (P(I,J) + P(I,ILO)) * 0.5D0
   17 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   18 CONTINUE
      ICOUNT=ICOUNT+NN
      IF(ICOUNT.LT.KCOUNT) GOTO 1000
      GOTO 900
С
С
         RETAIN REFLECTION
C
С
    19 CONTINUE
   19 DO 20 I=1,N
   20 P(I,IHI)=PSTAR(I)
      Y(IHI)=YSTAR
      GOTO 1000
С
С
         SELECT THE TWO BEST FUNCTION VALUES (YNEWLO
С
         AND YSEC) AND THEIR COORDS. (XMIN AND XSEC)
С
  900 DO 23 J=1,NN
      DO 22 I=1,N
   22 XMIN(I) = P(I,J)
      Y(J) = FN(XMIN)
   23 CONTINUE
      YNEWLO=BIGNUM
      DO 24 J=1,NN
      IF(Y(J).GE.YNEWLO) GOTO 24
      YNEWLO=Y(J)
      IBEST=J
   24 CONTINUE
      Y(IBEST)=BIGNUM
      YSEC=BIGNUM
```

```
DO 25 J=1,NN
      IF(Y(J).GE.YSEC) GOTO 25
      YSEC=Y(J)
      ISEC=J
   25 CONTINUE
      DO 26 I=1,N
      XMIN(I) = P(I, IBEST)
      XSEC(I) = P(I, ISEC)
   26 CONTINUE
      RETURN
      END
С
С
C-
С
   LAPLACE SOLUTION (EIGENVALUE SOLUTION)
С
                      THEODORSEN: 7th order in s
С
                                     approximate C(k) \& C(k) * ik
C--
С
      DOUBLE PRECISION FUNCTION FN(lag)
      IMPLICIT REAL*8 (A-Z)
      DOUBLE PRECISION lag(4),t(1,2),tst(1,2)
      INTEGER ib, n, ia, ijob, iz, ier, infer, i, j
      COMPLEX*16 eval(14), A(14, 14), B(14, 14), EIGA(14), EIGB(14), Z(14, 14)
      COMPLEX*16 WK(24,45)
      REAL*8 u,vel,velndnc,velck,atemp1,atemp2,atemp3,atemp4
      REAL*8 trm1, trm2, trm3, trm4, trm5, trm6, trm7, trm8
      REAL*8 trm9, trm10, trm11, trm12, trm13, trm14
      REAL*8 finv(8,2,2),f(9,2,2),m(2,2),c(2,2),k(2,2)
      REAL*8 mc(2,1), cc(2,1), kc(2,1)
      REAL*8 a0(2,2),a1(2,2),a2(2,2),a3(2,2),a4(2,2),a5(2,2),a6(2,2)
      REAL*8 a0c(2,1), a1c(2,1), a2c(2,1), a3c(2,1)
      REAL*8 a4c(2,1), a5c(2,1), a6c(2,1)
      REAL*8 n1(2,7)
       INTEGER y, check
C
C transfer function constants
              t(1,1) = LAG(1)
              t(1,2) = LAG(2)
              tst(1,1) = LAG(3)
              tst(1,2) = LAG(4)
С
С
              velndnc=4.53
              step=0.05
              u=50.0
              wrat=0.2
              r=0.5
              ah=-0.5
С
              ta=0.0012
              tb=0.5484
              cst=50000.0
              ds=cst*ta/tb
```

```
velndbig=2*velndnc
С
C Read in constants.
      OPEN (UNIT=1, FILE='CONST6.DAT', STATUS='OLD')
      READ (1,*) ia, ib, iz, n, ijob
                  14 14 14 14 2
С
      CLOSE (UNIT=1)
С
C open required files
      check=check+1
      if(check.gt.1) goto 66
      OPEN (UNIT=4, FILE='EIGEN.OUT', STATUS='NEW')
      OPEN (UNIT=5, FILE='OPTIMAL.OUT', STATUS='NEW')
C required constants
              pi=3.14159
С
              xtor=0.25
              semic=1.0
              wtor=1.0
              wben=wrat*wtor
С
              xbeta=0.25
              rbeta=0.5
              c1=0.6
С
C required constants
C
              t_{1}=-(1/3)*(1-c_{1}*2)**0.5*(2+c_{1}*2)+c_{1}*acos(c_{1})
              t2=c1*(1-c1**2)-(1-c1**2)**0.5*(1+c1**2)*acos(c1)+c1*
              (acos(c1))**2
     1
              t3=-((1/8)+c1**2)*(acos(c1))**2+0.25*c1*(1-c1**2)**0.5*
     1
              acos(c1)*(7+2*c1**2)-(1/8)*(1-c1**2)*(5*c1**2+4)
              t4=-acos(c1)+c1*(1-c1**2)**0.5
              t_{5=-(1-c_{1}*2)-(acos(c_{1}))*2+2*c_{1}*(1-c_{1}*2)*0.5*acos(c_{1})}
              t6=t2
              t7=-((1/8)+c1**2)*acos(c1)-(1/8)*c1*(1-c1**2)**0.5*
     1
              (7+2*c1**2)
              t8 = -(1/3) * (1 - c1 * * 2) * * 0.5 * (2 * c1 * * 2 + 1) + c1 * acos(c1)
              t9=0.5*((1/3)*(1-c1**2)**1.5+ah*t4)
              t10=(1-c1**2)**0.5+acos(c1)
              t11=acos(c1)*(1-2*c1)+(1-c1**2)**0.5*(2-c1)
              t12=(1-c1**2)**0.5*(2+c1)-acos(c1)*(2*c1+1)
              t_{13=0.5*(-t_7-(c_1-a_h)*t_1)}
              t14=(1/16)+0.5*ah*c1
C required constants
С
      OPEN (UNIT=10, FILE='LAPLACE.DAT', STATUS='OLD')
           do 25 i=1,2
      READ (10,*) n1(i,1),n1(i,2),n1(i,3),n1(i,4),n1(i,5),
     1
                   n1(i,6), n1(i,7)
 25
           continue
      READ (10,*) b1,b2,b3,b4,err
      CLOSE (UNIT=10)
С
```

```
C FOUR LAG TERMS
С
C PADE APPROXIMANT
                      THEODORSEN FUNCTION
                                                C(k)
                                                       AND
                                                            C(k) * ik
С
C term Q(1,1): C(k) *ik
C
              aO(1,1) = n1(2,1)
              al(1,1) = nl(2,2)
              a2(1,1) = n1(2,3)
              a3(1,1) = n1(2,4)
              a4(1,1) = n1(2,5)
              a5(1,1) = n1(2,6)
              a6(1,1) = n1(2,7)
С
C term Q(1,2): C(k)*ik*(0.5-ah)+C(k)
              a0(1,2) = n1(2,1) * (0.5-ah) + n1(1,1)
              a1(1,2) = n1(2,2) * (0.5-ah) + n1(1,2)
              a2(1,2) = n1(2,3) * (0.5-ah) + n1(1,3)
               a3(1,2) = n1(2,4) * (0.5-ah) + n1(1,4)
              a4(1,2) = n1(2,5) * (0.5-ah) + n1(1,5)
               a5(1,2) = n1(2,6) * (0.5-ah) + n1(1,6)
              a6(1,2) = n1(2,7) * (0.5-ah) + n1(1,7)
С
C term Q(2,1): C(k)*ik*(-(ah+0.5))
С
              a0(2,1) = n1(2,1) * (-0.5-ah)
              al(2,1) = nl(2,2) * (-0.5-ah)
              a2(2,1) = n1(2,3) * (-0.5-ah)
               a3(2,1) = n1(2,4) * (-0.5-ah)
               a4(2,1) = n1(2,5) * (-0.5-ah)
              a5(2,1) = n1(2,6) * (-0.5-ah)
               a6(2,1) = n1(2,7) * (-0.5-ah)
С
C term Q(2,2): C(k)*ik*(-(ah+0.5)*(0.5-ah))-(ah+C.5)*C(k)
C
              a0(2,2) = n1(2,1) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,1)
              a1(2,2) = n1(2,2) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,2)
              a2(2,2) = n1(2,3) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,3)
              a3(2,2) = n1(2,4) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,4)
              a4(2,2) = n1(2,5)*(-0.5-ah)*(0.5-ah)-(ah+0.5)*n1(1,5)
              a5(2,2) = n1(2,6) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,6)
              a6(2,2) = n1(2,7) * (-0.5-ah) * (0.5-ah) - (ah+0.5) * n1(1,7)
С
С
            flap control
С
C term Qc(1,1): ik*C(k)*t11/(2*pi)+t10*C(k)/pi
С
              a0c(1,1) = n1(2,1)*t11/(2*pi)+n1(1,1)*t10/pi
              alc(1,1) = n1(2,2)*t11/(2*pi)+n1(1,2)*t10/pi
              a2c(1,1) = n1(2,3)*t11/(2*pi)+n1(1,3)*t10/pi
              a3c(1,1)= n1(2,4)*t11/(2*pi)+n1(1,4)*t10/pi
               a4c(1,1)= n1(2,5)*t11/(2*pi)+n1(1,5)*t10/pi
              a5c(1,1) = n1(2,6)*t11/(2*pi)+n1(1,6)*t10/pi
```

•..

28

 \mathbf{Z}

```
a6c(1,1) = n1(2,7) *t11/(2*pi)+n1(1,7) *t10/pi
С
C term Qc(2,1) : -(0.5+a)*(ik*C(k)*t11/(2*pi)+t10*C(k)/pi)
             a0c(2,1)= (n1(2,1)*t11/(2*pi)+n1(1,1)*t10/pi)*(-0.5-ah)
             a1c(2,1)= (n1(2,2)*t11/(2*pi)+n1(1,2)*t10/pi)*(-0.5-ah)
             a2c(2,1)= (n1(2,3)*t11/(2*pi)+n1(1,3)*t10/pi)*(-0.5-ah)
             a3c(2,1)= (n1(2,4)*t11/(2*pi)+n1(1,4)*t10/pi)*(-0.5-ah)
             a4c(2,1) = (n1(2,5)*t11/(2*pi)+n1(1,5)*t10/pi)*(-0.5-ah)
             a5c(2,1)= (n1(2,6)*t11/(2*pi)+n1(1,6)*t10/pi)*(-0.5-ah)
             a6c(2,1) = (n1(2,7)*t11/(2*pi)+n1(1,7)*t10/pi)*(-0.5-ah)
С
С
С
             b1234=b1+b2+b3+b4
             b123==b1+b2+b3
             b124=b1+b2+b4
             b134=b1+b3+b4
             b234=b2+b3+b4
             bt123=b1*b2*b3
             bt124=b1*b2*b4
             bt134=b1*b3*b4
             bt234=b2*b3*b4
             ba1234=b1*b2*b3*b4
             bt1234=b1*b2*b3+b1*b2*b4+b1*b3*b4+b2*b3*b4
             bd1234=b1*b2+b1*b3+b2*b3+b1*b4+b2*b4+b3*b4
             bd123=b1*b2+b1*b3+b2*b3
             bd124=b1*b2+b1*b4+b2*b4
             bd134=b1*b3+b1*b4+b3*b4
             bd234=b2*b3+b3*b4+b2*b4
С
С
C titles and initial parameter listing to file
 66
             write(4,34) t(1,1),t(1,2)
             write(4,34) tst(1,1),tst(1,2)
             write(4, 32)
             write(4,121)
             write(4,121)
             write(4,33)
             write(4,121)
             write(4,34) u,r,ah,wrat,velndnc
             format(6x, ' LAPLACE METHOD
 32
                                           INCOMPRESSIBLE (THEODORSEN) ')
             format(6x,'
 33
                            u
                                        r(alp)
                                                   ah
                                                          wben/wtor
     1
                     'no control velocity')
 34
             format(6x, f7.1, 3f10.1, f8.2)
                        ')
121
             format('
С
C iterate velocity
             vel=velndnc+0.1
             velck=velndnc*1.3
 60
             vel=vel+step
             if(vel/(semic*wtor).gt.velck) goto 800
             wf=vel/semic
             xvar=2*wf**2/u
С
```

** *

```
C formation of matrices :
                            [M], [C], [K], [D] and [E]
             m(1,1) = 1 + 1/u
             m(1,2) = xtor - ah/u
             m(2,1) = xtor-ah/u
             m(2,2) = r * *2 + (0.125 + ah * *2)/u
              C(1,1)=0.0
              c(1,2) = wf/u
              c(2,1)=0.0
              c(2,2) = (0.5-ah) * wf/u
              k(1,1)=wrat**2*wtor**2
             k(1,2)=0.0
              k(2,1)=0.0
              k(2,2)=r**2*wtor**2
              mc(1,1) = xbeta - t1/(pi*u)
              mc(2,1)=rbeta**2+(c1-ah)*xbeta+(-t7-(c1+0.5)*t1+
     1
              (0.5-ah) *t1)/(pi*u)
              cc(1,1) = -wf + t4/(u + pi)
              cc(2,1) = (wf/u*pi)*(-(t4-(2/3)*(1-c1**2)**1.5)+(0.5+ah)*t4)
              kc(1,1)=0.0
              kc(2,1) = (wf**2/u)*(t4+t10)/pi
С
C formation of polynomial coefficients :
                                             [F8], [F7], [F6], [F5], [F4]
С
                                             [F3], [F2], [F1]
С
С
    [F8]s7+[F7]s6+[F6]s5+ \dots + [F2]s+[F1] = 0
          do 200 i=1,2
          do 100 j=1,2
С
              calculation of [F8]
                      trm1=m(i,j)+xvar*(1/wf)**2*a2(i,j)
                      trm2=mc(i,1)+xvar*(1/wf)**2*a2c(i,1)
              f(8,i,j) = trm1 + trm2 + t(1,j) + ds + trm2 + tst(1,j)
С
              calculation of [F7]
                      trm3=c(i,j)+wf*b1234*m(i,j)+xvar*(1/wf)*
     1
                       (a1(i,j)+b1234*a2(i,j))
                      trm4=cc(i,1)+wf*b1234*mc(i,1)+xvar*(1/wf)*
     1
                       (alc(i,1)+b1234*a2c(i,1))
              f(7,i,j)=trm3+trm4*t(1,j)+ds*trm4*tst(1,j)+cst*trm1+
              cst*trm2*t(1,j)
     1
С
              calculation of [F6]
                      trm5=k(i,j)+wf*b1234*c(i,j)+wf**2*bd1234*
     1
                      m(i,j)+xvar*(a0(i,j)+b1234*a1(i,j)+bd1234*
                      a2(i,j)+a3(i,j)+a4(i,j)+a5(i,j)+a6(i,j))
     1
                      trm6=kc(i,1)+wf*b1234*cc(i,1)+wf**2*bd1234*
     1
                      mc(i,1)+xvar*(a0c(i,1)+b1234*a1c(i,1)+bd1234*
     1
                      a2c(i,1)+a3c(i,1)+a4c(i,1)+a5c(i,1)+a6c(i,1))
              f(6,i,j)=trm5+trm6*t(1,j)+ds*trm6*tst(1,j)+cst*
     1
              trm3+cst*trm4*t(1,j)
С
              calculation of [F5]
                      trm7=wf*b1234*k(i,j)+wf**2*bd1234*c(i,j)+wf**3*
     1
                      bt1234*m(i,j)+xvar*wf*(b1234*a0(i,j)+bd1234*
     1
                      al(i,j)+bt1234*a2(i,j)+b234*a3(i,j)+b134*a4(i,j,+
     1
                      b124*a5(i,j)+b123*a6(i,j))
                      trm8=wf*b1234*kc(i,1)+wf**2*bd1234*cc(i,1)+
     1
                      wf**3*bt1234*mc(i,1)+xvar*wf*(b1234*a0c(i,1)+
```

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A = 119
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	1	bd1234*a1c(i,1)+bt1234*a2c(i,1)+b234*a3c(i,1)+b134*a4c(i,1)+b124*a5c(i,1)+b123*a6c(i,1))
	Ŧ	f(5,i,j) = trm7 + trm8 * t(1,j) + ds * trm8 * tst(1,j) + cst*
_	1	trm5+cst*trm6*t(1,j)
С		calculation of [F4]
	-	trmy=WI**2*DQ1234*K(1, J)+WI**3*Dt1234*C(1, J)+
	1	$WI \sim 4 \sim DQI234 \sim M(1, J) + XVal \sim WI \sim 2 \sim (DQI234 \sim dQ(1, J) + b+ 1024 + a)(i - i) + b < a < a < a < a < a < a < a < a < a <$
	1	bd134 * a4(i, j) + bd124 * a5(i, j) + bd123 * a6(i, j)
	1	trm10=wf**2*bd1234*kc(i,1)+wf**3*bt1234*cc(i,1)+
	1	wf**4*bg1234*mc(i,1)+xvar*wf**2*(bd1234*a0c(i,1)+
	ī	bt1234*a1c(i,1)+bg1234*a2c(i,1)+bd234*a3c(i,1)+
	ī	bd134*a4c(i,1)+bd124*a5c(i,1)+bd123*a6c(i,1))
		f(4,i,j)=trm9+trm10*t(1,j)+ds*trm10*tst(1,j)+cst*
	1	trm7+cst*trm8*t(1,j)
С		calculation of [F3]
		trm11=wf**3*bt1234*k(i,j)+wf**4*bq1234*c(i,j)+
	1	xvar*wf**3*(bt1234*a0(i,j)+bq1234*a1(i,j)+
	1	bt234*a3(i,j)+bt134*a4(i,j)+bt124*a5(i,j)+
	1	bt123*a6(i,j))
		trm12=wf**3*bt1234*kc(i,1)+wf**4*bq1234*cc(i,1)+
	1	xvar*wf**3*(bt1234*a0c(i,1)+bq1234*a1c(i,1)+
	1	bt234*a3c(1,1)+bt134*a4c(1,1)+bt124*a5c(1,1)+
	1	bt123*a6c(1,1))
	-	I(3,1,]) = trm11 + trm12 * t(1,]) + ds * trm12 * tst(1,]) + cst * trm0 + cst + tst(1, -1) + cst * tst(1, -1) + cst(1, -1
<u> </u>	T	trm9+CSt*trm10*t(1,])
C.		Calculation of [r2]
	1	CIMIS-WIAAADQI234AK(I,J)+XVarAWIAA4ADQI234A
	T	trm14=wf**4*bal234*kc(i 1)+xvar*wf**4*bal234*
	1	a0c(i,1)
	-	f(2,i,i) = trm13 + trm14 + t(1,i) + ds + trm14 + tst(1,i) + cst + ds + trm14 + tst(1,i) + cst + ds +
	1	trm11+cst*trm12*t(1,j)
С		calculation of [F1]
		f(1,i,j)=cst*(wf**4*bq1234*k(i,j)+xvar*wf**4*bq1234*
	1	a0(i,j))+cst*(wf**4*bq1234*kc(i,1)+xvar*wf**4*bq1234*
	1	a0c(i,1))*t(1,j)
100		continue
200		continue
С		
С		calculation of inverse (inv[F9])
		f8det=f(8,1,1)*f(8,2,2)-f(8,1,2)*f(8,2,1)
		if(f8det.eq.0.0) goto 210
210		goto 220
210		Write(4,215)
212	1	IORMAT(' DETERMINANT EQUALS ZERO - INVESTIGATE INPUT ',
	Ŧ	rete 800
220		finv(9, 1, 1) - f(9, 2, 2) / f9dot
~~0		f(9, 1, 1) = f(0, 2, 2) / 100 = 0
		finv(8,2,2) = f(8,1,1) / f8det
		f(9,2,2) = finv(8,2,2)
		finv(8, 1, 2) = -f(8, 1, 2) / f8det
		f(9,1,2)=finv(8,1,2)

Ľ
```
finv(8,2,1) = -f(8,2,1)/f8det
             f(9,2,1) = finv(8,2,1)
С
C formation of matrices [A] and [B] to
C solve eigenvalue problem given by [A] \{x\} = lambda [B] \{x\}
          do 400 i=1,n
          do 300 j=1,n
             a(i,j) = dcmplx(0.0,0.0)
             if(i.eq.j+2) a(i,j)=dcmplx(1.0,0.0)
             b(i,j) = dcmplx(0.0,0.0)
              if(i.eq.j) b(i,j)=dcmplx(1.0,0.0)
 300
          continue
 400
          continue
          do 500 i=1, n/2
             y=n/2+1-i
             atempl=-(finv(8,1,1)*f(y,1,1)+finv(8,1,2)*f(y,2,1))
              a(1,2*i-1)=dcmplx(atemp1,0.0)
              atemp2=-(finv(8,1,1)*f(y,1,2)+finv(8,1,2)*f(y,2,2))
              a(1,2*i) = dcmplx(atemp2,0.0)
              atemp3=-(finv(8,2,1)*f(y,1,1)+finv(8,2,2)*f(y,2,1))
              a(2,2*i-1)=dcmplx(atemp3,0.0)
              atemp4=-(finv(8,2,1)*f(y,1,2)+finv(8,2,2)*f(y,2,2))
              a(2,2*i) = dcmplx(atemp4,0.0)
 500
          continue
С
С
C subrountine for eigenvalue calculation
C abar : real part of eval()
C bbar : imaginary part of eval()
      CALL EIGZC(A, IA, B, IB, N, IJOB, EIGA, EIGB, Z, IZ, WK, INFER, IER)
          do 5 I=1,n
 5
              EVAL(I) = EIGA(I) / EIGB(I)
С
C printout of velocity and associated eigenvalues
C flutter velocity check
              vndcrit=0.0
          do 6 i=1,n
С
               write(4,22) eval(i)
              if(dreal(eval(i)).gt.0.0001) vndcrit=vel/(semic*wtor)
С
      1
               dreal(eval(i)),dimag(eval(i))/wtor
 6
           continue
               FN=velndbig-vndcrit
               perc=((vndcrit-velndnc)/velndnc)*100
               if(vndcrit.gt.0.0) goto 779
 21
              format(6x,' performance index is', 2f15.5)
              format (6x, ' ND VELOCITY ', f6.3, ' Ta ', f6.4)
 23
 22
              format(6x,8f17.5)
              format(6x,f17.5,',',f17.5,',',f17.5)
 28
              format(6x,'NO CONTROL VELOCITY ',f5.2,'
 41
                                                         CRITICAL ',
              'VELOCITY ', f5.2,' % INCREASE ', F5.2,' %')
     1
 42
              format(6x, ' Ta ', f7.4, '
                                         Tb ',f7.4)
              format(6x,' [ t ]
 43
                                   ',2f14.4)
              format(6x,' [ tst ] ',2f14.4)
 44
С
```

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A = 121
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Acres 1

700		goto 60
779		write(5,41) velndnc,vndcrit,perc
		write(5,42) ta,tb
		write(5,43) t(1,1),t(1,2)
		write(5,44) tst(1,1),tst(1,2)
		write(5,121)
789	RETURN	

800 END

1

INPUT FILE : LAPLACE.DAT

 .99828226
 -.00003562
 .00000474
 -.04120375

 -.16289333
 -.22908745
 -.06459371
 -.00003629
 .00061076

 -.00001355
 .50037475
 -.00003629
 .00061076

 .01380489
 .05848218
 .04963449
 .000011782

INPUT FILE : CONST6.DAT

14 14 14 14 2

OUTPUT FILE : OPTIMAL.OUT

NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.78	<pre>% INCREASE</pre>	5.52 %
		3000			
	[tst] 4.0000	3.6000			
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.68	<pre>% INCREASE</pre>	3.31 %
		- 3000			
	$[t_{st}] = 4.0000$	3,6000			
		510000			
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.68	% INCREASE	3.31 %
	Ta .0012 Tb .5484				
	[t] .0000	4000			
	[tst] 4.0000	3.6000			
NO		VELOCITY	1 79	9 INCOFASE	5 52 8
NO	Ta .0012 Tb .5484	VEDOCITI	4.70	J INCKEASE	J.J2 7
	[t] .0000	3000			
	[tst] 4.1000	3.6000			
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.78	% INCREASE	5.52 %
	Ta .0012 Tb .5484				
	[t].0000	3000			
	[tst] 4.0000	3.7000			
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.73	% INCREASE	4.42 %
	Ta .0012 Tb .5484				
	[t] .0500	3000			
	[tst] 4.0500	3.6000			
					•
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.88	% INCREASE	7.73 8
		- 3500			
	[tst] 4.0500	3.6000			
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.78	% INCREASE	5.52 %
	Ta .0012 Tb .5484				
	[t].0000	3000			
	[tst] 4.1000	3.6000			
NO	CONTROL VELOCITY 4 53 CRITICAL	VELOCITY	A 79	S INCREASE	5 52 ¥
110	Ta .0012 Tb .5484	VELOCI11	4./0	A THCKENDE	J.J2 8
	[t] .0000	3000			
	[tst] 4.0500	3.6500			
				_	
NO	CONTROL VELOCITY 4.53 CRITICAL	VELOCITY	4.78	<pre>% INCREASE</pre>	5.52 %
	Ta .0012 Tb .5484				
		3000			
	[TST] 4.0500	2.0000			

1%. •

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NO	CONTROL VELOCITY 4.53 CRITIC	CAL VELOCITY	4.73	INCREASE	4.42 %
		- 3000			
		3 6000			
	[[]] 4.0500	3.0000			
NO	CONTROL VELOCITY 4.53 CRITIC	CAL VELOCITY	4.88	% INCREASE	7.73 %
	Ta .0012 Tb .5484				
	[t] .0000	3500			
	[tst] 4.0500	3.6000			
NO	CONTROL VELOCITY 4.53 CRITI	CAL VELOCITY	4.78	* INCREASE	5.52 %
	$T_{a} = 0.012$ The 5484		4170		5.52 0
		- 3000			
		5000			
		3.0000			
NO	CONTROL VELOCITY 4.53 CRITI	CAL VELOCITY	4.78	% INCREASE	5.52 %
	Ta .0012 Tb .5484				
	[t] .0000	3000			
	[tst] 4.0500	3.6500			
NO	CONTROL VELOCITY 4.53 CRITI	CAL VELOCITY	4.78	* INCREASE	5.52 %
	Ta .0012 Tb .5484				
		- 3000			
		3 6000			
	[[[]]] 4.0500	J.0000			

OUTPUT FILE : DATA.OUT

•#**\$** ** SEQUENTIAL SIMPLEX : PROBLEM MINIMIZATION

12 TRIALS USED

ESTIMATES

PARAMETER	NEXT-TO-BEST	BEST	
1	.000000	.0000000	
2	3000000	3500000	
3	4.1000000	4.0500000	
4	3.5999999	3.5999999	

FUNCTION VALUES

NEXT-TO-BEST	BEST
--------------	------

4.280000206 4.180000205