

INFRARED INTENSITIES OF THE CO STRETCHING MODES OF SOME TRANSITION METAL CARBONYL COMPLEXES

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By

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ABSTRACT

The work in this thesis constitutes an additional study of the infrared intensities of the carbonyl stretching absorptions of transition metal carbonyl complexes. Q-intensity data are reported for the following complexes: $Cr(\Omega)_6$, $(Cycloheptatriene)H(\Omega)_3$ and (Bicyclohepta $diene)M(\Omega)_4$ (H = Cr, Mo, H), <u>fac-Mo(Ω)_3[P(C_6H_5)_{3-x}Cl_x]_3 (x = 1-3), cis-Mn(Ω)_4 LBr [L = P(C_6H_5)_3, As(C_6H_5)_3, Sb(C_6H_5)_3], Π -RC_5H_4Mn(Ω)_3 (R = H, CH₃), and Π -C₅H₅Mn(Ω)₂CS.</u>

The integrated intensity results are used as a basis for a further investigation of the o- and π -bonding abilities of the various ligands in these complexes. The first measurement of the infrared intensity of a thiocarbonyl stretching mode in a metal thiocarbonyl complex, π -c₅H₅Mn(∞)₂CS, is also reported.

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RESUME

Le contenu de cette thèse est une étude addition:melle des intensités d'absorption par élongation, en spectroscopie infrarougé, des groupes carbonyles des complexes carbonylés des mótaux de transition. Les valeurs des intensités (Ω) ont été déterminées pour les complexes suivants: $Cr(\Omega)_6$, $(Cycloheptatriene)M(\Omega)_3$ et (Bicycloheptadiene)- $M(\Omega)_4$ (M = Cr, ho, W), <u>fac-Mo(\Omega)_3[P(C_6H_5)_{3-X}Cl_X]_3</u> (x = 1-3), <u>cis-Mn(\Omega)_4LBr [L = P(C_6H_5)_3</u>, As(C_6H_5)_3, Sb(C_6H_5)_3], π -RC_5H_4Mn(\Omega)_3 (R = H, CH_3), et π -C_5H_5Mn(Ω)₂CS.

Les résultats obtenus par intégration des intensités servent de base pour l'étude des possibilités qu'ont les liaisons 0 et π de former différents ligands chez ces complexes. La première mesure de l'intensité, en spectroscopie infrarouge, d'un type d'élongation d'un groupement thiocarbonyle du complexe métallique thiocarbonylé, π -C₅H₅Mn(∞)²CS, est également relatée.

Master of Science Degree

Department of Chemistry HoGill University To my parents, Mr. & Mrs. H. T. Johansson, for their unfailing support

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Chapter 1

INTENSITY OF AN INFRARED ABSORPTION BAND

General

Every distinct infrared absorption band is characterized by the frequency at which maximum absorption occurs and by the integrated absorption intensity. Both band characteristics are related in a fundamental way to the electronic properties of the molecule and consequently provide very useful information concerning molecular properties. However, although absorption frequencies of molecules have been studied extensively, absorption intensities have received comparatively little attention. This apparent lack of interest in absorption intensity studies is due largely to the difficulties and tediousness of the experimental procedures and to the theoretical problems encountered in the interpretation of results. However, now that high resolution infrared spectrophotometers with linear absorbance facilities have become generally available, the practical difficulties, at least, have been minimized. Consequently, as predicted by Brown¹ in 1958 in his review of infrared intensities, much more work has been carried out in recent years.

According to the classical model, a molecule consists of a set of point masses connected to one another by a set of elastic forces. The system is considered to undergo vibrational motion in accordance with the laws of mechanics;^{2,3} the 3N-6 vibrational degrees of freedom (3N-5 for a linear molecule) of an N-atom molecule are termed the normal modes.

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If the vibrational motion associated with a normal vibration produces a periodic shift in charge distribution and this shift results in a change in molecular dipole moment, the vibration is said to be infraredactive. Since any vibrational motion of the molecule can be resolved as a superposition of the normal vibrations, the frequency at which the absorption occurs is the normal vibration frequency, and the intensity of the absorption is the intensity corresponding to the normal vibration.¹

Expressions for Band Intensity

The true integrated absorption intensity of an infrared absorption band, A, corresponding to the ith fundamental vibrational transition, is defined as

$$A_{i} = \int_{\text{band}} \alpha_{v}(i) dv = \frac{1}{cl} \int_{\text{band}} \ln(\frac{I_{o}}{I}) v dv \qquad (1)$$

where $\alpha_v(1)$ is the adsorption coefficient of the band corresponding to the ith normal vibration at frequency v, I_{0v} and I_v are the incident and transmitted intensities of monochromatic radiation of frequency v, c is the concentration of solute in moles liter⁻¹, and l is the cell path length in centimeters. Such a definition follows readily from the well-known exponential law of absorption,

$$I_{v} = I_{ov} e^{-\alpha} v^{cl}$$
 (2)

The absorption intensity is related to the polar properties of a molecule through its direct proportionality to the square of the molecular dipole moment derivative²

$$\int_{\text{band}} \alpha_{\nu}(\mathbf{i}) d\nu = \frac{N}{3c^2} \left[\left(\frac{\partial \mu}{\partial \mathbf{Q}_{\mathbf{i}}} \right)_0 \right]^2$$

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where N is the number of molecules per unit volume of sample, c is the vwinsity of light and \mathbf{Q}_{i} is the normal co-ordinate for the ith vibration. The vector $\boldsymbol{\mu}$ is the molecular dipole moment and $(\partial \boldsymbol{\mu}/\partial \mathbf{Q}_{i})_{0}$ is its rate of change with respect to \mathbf{Q}_{i} . The quantity $(\partial \boldsymbol{\mu}/\partial \mathbf{Q}_{i})_{0}$, the molecular dipole moment derivative, is evaluated at the molecular equilibrium configuration. This equation expresses the measured intensities in terms of the normal co-ordinates of a molecule. In order to interpret the intensities in terms of bond properties, one must reduce the dipole moment derivatives, $(\partial \boldsymbol{\mu}/\partial \mathbf{Q}_{i})_{0}$, to quantities which are characteristic of individual bonds. The difficulties encountered in such an operation are best appreciated by considering the assumptions which are involved.

Equation (3) is derived on the assumption that the molecular dipole moment can be expanded as a Taylor series in terms of nuclear displacements from the equilibrium positions and that all but the first derivatives can be neglected; that is, *

 $\mu = \mu_0 + \frac{\xi}{i} (\partial \mu / \partial Q_i)_0 Q_i + \text{higher terms (negligible)}$ (4)

This is the assumption of electrical harmonicity. It is true only to a first approximation, since intensities of infrared combination and overtone bands ought to be zero in this approximation. Although this is certainly not the case, it is a fairly reasonable assumption since the intensities of combinations and overtones are usually far less than those of the fundamentals (unless a combination or overtone band gains intensity by resonance from a fundamental of the same symmetry). At any rate, the error involved in the assumption of harmonicity is probably much less than the present experimental errors and those arising from the uncertainty in the normal co-ordinate Q_1 .^{4,5}

A simpler way of representing Q_1 is to express it as a linear ', combination of molecular symmetry co-ordinates, S_1 ,²

 $Q_1 = \xi L_{11}^{-1} S_1$ (5) where the coefficients L_{ij} are obtained from the molecular geometry, the masses of the atoms and the potential energy function of the molecule. It is here that the real difficulty arises. Although atomic masses are accurately known and molecular geometries can often be independently determined or estimated fairly accurately, the force field (i.e., the force constants connecting all atoms in the molecule), in terms of which the potential energy function is expressed, constitutes the unknown in a vibrational problem. Unfortunately, force constants are known with precision for very few molecules. Most molecules have more force constants than frequencies, and so their complete potential energy function cannot be determined. The usual procedure is to either assume that some of the force constants can be neglected while some can be carried over from other molecules, or to adopt a model force field which decreases the number of distinct force constants by postu-

lating that some symmetry unrelated force constants have related values (generally a simple proportion). Thus, the major source of uncertainty in the interpretation of infrared absorption bands in terms of bond properties is the form of the force field.⁴

Once a reasonable force field has been deduced, and hence the form of the normal co-ordinates approximated, the next step is to visualize what the resulting $(\partial \mu / \partial Q_i)_0$ mean. Because Q_i is a normal co-ordinate of the molecule, it is conceptually difficult to appreciate the significance of the derivatives $(\partial \mu / \partial Q_i)_0$ and to utilize them. If one could translate the molecular properties into the sum of a set of bond properties, one could reduce the derivatives $(\partial \mu / \partial Q_i)_0$ into those characteristic of individual bonds. In order to accomplish this, two assumptions are necessary:

- (1) The stretching of a bond by dr produces a change of dipole moment along the bond of $(\partial \mu / \partial r) dr$;
- (2) Changes in one bond do not result in changes in another bond.

This means essentially that the total molecular dipole moment change resulting from the simultaneous displacement of several bonds is assumed to be the vector sum of the moments produced by each individual bond. This assumption is quite reasonable for relatively simple molecules and for larger molecules possessing sufficient symmetry.

Now, since the molecular symmetry co-ordinates, S_j , may be expressed as linear combinations of internal bond co-ordinates, R_L ,²

$$s_{j} = \bigotimes_{k} U_{jk} R_{k}$$
 (6)

one can obtain [by utilizing equations (1), (4), (5), and (6)] the desired result: an expression relating the intensity of a normal mode to bond moment derivatives.⁴

$$A_{i} \propto \left(\frac{\partial \mu}{\partial Q_{i}}\right)^{2} = \left[\underset{j}{\leq} \underset{k}{\leq} L_{ji} U_{jk} \left(\frac{\partial \mu}{\partial R_{k}}\right) \right]^{2}$$
(7)

It follows from equation (7) that one should be able to gain useful information concerning the nature of the bond dipole moment changes upon alteration of bond lengths during a normal stretching vibration by studying the integrated absorption intensity of the normal mode.

Determination of Band Intensity

The expression for the true integrated absorption intensity, A, as cited earlier, is

$$A = \int_{\text{band}} \alpha_v dv = \frac{1}{cl} \int_{\text{band}} \ln(\frac{I_0}{I})_v dv$$
 (8)

where I_{0_V} and I_v are the incident and transmitted intensities of monochromatic radiation of frequency v, α_v is the adsorption coefficient for frequency v, c is the concentration of solute in moles liter⁻¹, and l is the cell path length in centimeters.

Unfortunately, owing to the presence of finite slit widths, the radiation is not monochromatic and the quantity that is actually measured is the apparent integrated absorption intensity, B, given by

$$B = \frac{1}{cl} \int_{band} \ln(\frac{T_o}{T}) \sqrt{d\nu}$$
 (9)

where T_{O_V} and T_V are the apparent intensities of the incident and transmitted radiation when the spectrophotometer is set at frequency V.

The true integrated absorption intensity, A, can be determined essentially in two ways:

- (1) by measuring the area under the experimental absorption
 band and extrapolating the B values so determined as
 a function of various parameters;⁶⁻⁹
- (2) by direct calculation from principal band characteristics, such as the peak height and the half-intensity band width, assuming certain functions for the slit and for the shape of the absorption band.^{8,10-17}

Direct calculation methods have the advantage that they are generally extremely rapid compared to band-area extrapolations. However, a serious limitation of direct calculation methods lies in the fact that they usually depend on measurements at only a few points on the experimental band while band-area methods involve measurements over the whole of the experimental curve. Assumptions regarding the slit function and the shape of the absorption band are also fundamental to direct calculation methods. Therefore, if a quick crude estimate of intensity is all that is desired, direct calculation methods are sufficient. However, for accurate work, band-area extrapolation procedures are usually the methods of choice. The extent of this discussion, therefore, will be limited to a consideration of this technique.

The first two extrapolation procedures for the determination of true integrated absorption intensities from experimental absorption curves were described by Bourgin⁶ and by Wilson and Wells⁷ and were applied to gases. The methods involve extrapolations of measurements of band areas made at a series of different concentrations or path lengths to zero concentration or path length and are equally applicable in the case of liquids and solutions.

1. The Fourgin Extrapolation

Bourgin⁶ plotted the absorption curve as fractional absorption, $1-(T/T_0)_{v}$, against frequency, v, in cm⁻¹; the area, A', under the curve is then given by

$$A' = \int_{\text{band}} \left[1 - \left(\frac{T}{T_0}\right)_{\nu}\right] d\nu \qquad (10)$$

Since it has been shown that this area is independent of the resolving power, we may write

$$A^{\prime} = \int_{\text{band}} \left[1 - \left(\frac{I}{I_0}\right)_{V_0}\right] dv \qquad (11)$$

Substituting in the exponential law of absorption and expanding we get

$$\frac{A^{\prime}}{cl} = \int_{\text{band}} \left[\alpha_{\nu} - \frac{\alpha_{\nu}^{2}}{2!} cl + \frac{\alpha_{\nu}^{3}}{3!} (cl)^{2} - \dots \right] d\nu \qquad (12)$$

or

$$\lim_{\substack{n \to 0}} \frac{A^*}{cl} = \int_{\text{band}} \alpha_v dv = A$$
(13)

Hence, the true integrated absorption intensity may be obtained by plotting A'/cl against cl and extrapolating to cl = 0. Since the area, A', is independent of the resolving power used, the nature of the extrapolation is also independent of the resolving power and is only dependent on the band shape.

2. <u>The Wilson-Wells Extrapolation</u> Wilson and Wells⁷ demonstrated that

$$\lim_{cl\to 0} (A-B) = 0$$

or in other words

$$\lim_{c \to 0} B = A$$
(15)

(14)

Therefore, by extrapolating the apparent integrated absorption intensities taken at a number of concentrations or path lengths to zero concentration or path length, the true integrated absorption intensity is obtained. The method involves the assumption that (i) the incident

A

intensity is constant over the slit width and (ii) the resolving power is constant over the width of the band. The nature of the extrapolation is dependent on the true band shape and on the resolving power used.

3. Ramsay's Methods

Ramsay⁸ modified both the Bourgin and the Wilson-Wells procedures by introducing an assumption concerning the form of the absorption³ curves. He assumed that the true shape of an absorption band could be expressed as a Lorentz function

 $\ln(\frac{I_0}{I})_{v} = \frac{a}{(v - v_0)^2 + b^2}$ (16)

where v_0 is the frequency of the band centre, a and b are constants, and $v-v_0$ is the interval on either side of the band centre over which the integration is performed. His assumption regarding the slit function was the usual one <u>viz</u>., a triangular slit function.

Nameay also employed correction factors for the "area under the wings." Although, in principle, the integration of equation (8) should be carried out to an infinite distance on either side of the band maximum, in practice, the area under the absorption curve is measured only to finite limits on either side of the band centre. In the absence of overlapping bands, the limits chosen are those at which the experimental error becomes of the same order of magnitude as the apparent intensity. Ramsay felt, by virtue of his calculations, that the residual unmeasured area could be appreciable, since although the absorption is small, the remaining frequency range is large. For this reason he applied "wing corrections" to all his areas, both calculated and measured. The wing are available for a series of values of the characteristic band quantity . $(v-v_{0})/b$.

a. Extension of the Bourgin Extrapolation

If the shape of the true absorption curve is represented by equation (16) then Bourgin's expression, equation (12), becomes

$$\frac{A^{\prime}}{cl} = A \left\{ 1 - \frac{1}{2!} \times \frac{1}{2} \ln(\frac{I_0}{I})_{\nu_{\text{max}}} + \frac{3!}{3!} \times \frac{3!}{4!2} \left[\ln(\frac{I_0}{I})_{\nu_{\text{max}}} \right]^2 - \frac{1}{4!} \times \frac{5!}{6!4!2} \left[\ln(\frac{I_0}{I})_{\nu_{\text{max}}} \right]^3 + \dots \right\}$$
(17)

Therefore, the true integrated intensity may be obtained by plotting A'/cl against $\ln(I_0/I)_{v_{max}}$ and extrapolating to $\ln(I_0/I)_{v_{max}} = 0$.

Alternatively, by defining a quantity, \emptyset , such that

$$\emptyset = 1 / \left\{ 1 - \frac{1}{2!} \times \frac{1}{2} \ln(\frac{I_0}{I})_{\nu_{\text{max}}} + \frac{1}{3!} \times \frac{3 \cdot 1}{4 \cdot 2} \left[\ln(\frac{I_0}{I})_{\nu_{\text{max}}} \right]^2 \dots \right\}$$
(18)

equation (18) can be more conveniently expressed as

$$A = \emptyset, \frac{A'}{cl}$$
(19)

By tabulating values of \emptyset for various values of $\ln(I_0/I)_{V_{\text{max}}}$, it is then possible to determine the true integrated absorption intensity from a single measurement of the apparent integrated absorption intensity by means of equation (19).

The extrapolations and \emptyset values depend entirely on the assumption of a Lorentzian function for the true band shape, but they are independent of the slit function. Likewise, the wing corrections provided by Ramsay for this method depend mainly on a Lorentzian band; however, they also depend to a very slight extent on the slit function as well.

b. Extension of the Wilson-Wells Extrapolation

Assuming a Lorentzian form for the true band shape, and a triangular slit function, Ramsay calculated, for a fixed value of $\frac{1}{2}\sqrt{\frac{1}{2}}$ (the ratio of the spectral slit width to the true half-intensity band width), a series of true band shapes for several values of the true peak optical density. He then measured the areas under the corresponding apparent absorption curves and plotted the ratios

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 $\left[\int_{band}\ln(T_0/T)_v dv\right] / \left[\int_{band}\ln(I_0/I)_v dv\right]$

against the apparent peak optical densities, $\ln(T_0/T)_{v_{max}}$. The plots he obtained were very nearly linear and had a slight negative slope, 9.

It follows that, by tabulating values of Θ for a range of $s/\Delta v_{\pm}^{T}$ between 0.1 and 1.0, the true integrated intensity of an absorption band can be obtained by plotting the apparent integrated absorption intensity measured at several different concentrations or path lengths against the apparent peak optical density and determining the best straight line with intercept A and slope AO. Ramsay felt this procedure of determining the intercept, A, was more accurate than the Wilson-Wells extrapolation based on the method of least-squares, since the latter method is dependent on measurements of weak bands where experimental errors are apt to be large.

One advantage of the 0 tabulations is that it is also possible to determine the true integrated absorption intensity from a single measurement of the apparent integrated absorption intensity by means of the relation

$$B = A + A \Theta \ln(T_0/T)$$
(20)

The Θ values, and hence the slopes of the extrapolations, depend on the assumption of a Lorentzian function for the true band shape and on a triangular slit function. The wing corrections provided by Ramsay for this method also depend on these two assumptions.

c. Method for Partially Overlapping Band Systems

Although Ramsay's methods are still valid for overlapping band systems, the slopes of the extrapolations can no longer be calculated. Moreover, Ramsay felt that attempts to fit extrapolation lines by the method of least-squares could result in considerable error unless a large number of observations was carried out.

However, in his studies of isolated absorption bands, Ramsay found that the true integrated absorption intensity of a band differed very little from the apparent quantity, even though the spectral slit width of the spectrophotometer was of the same order as the widths of the absorption bands. He showed that, for a slit width equal to onehalf the true width of the absorption band, the true intensity exceeded the apparent intensity by only $\sim 2-3\%$. Apparently the decrease in peak height produced by the finite slit width is roughly compensated by an increase in band width. As a general rule, Ramsay and colleagues found that the true integrated intensity of an isolated absorption band usually exceeded the mean of the apparent intensity determinations (i.e., B values were usually determined at a number of concentrations) by 2%. Since they felt that the uncertainty in the evaluation of the extrapolation slope for overlapping band systems was likely to exceed this considerably, they simply estimated the true intensity by adding 2% to the mean of the apparent intensity values.

4. The Russell-Thompson Extrapolation

Russell and Thompson² extrapolated the apparent integrated absorption intensities determined at a number of concentrations or path lengths to zero apparent peak optical density,

$$\lim_{n \to 0} B = A$$
(21)

as outlined by Ramsay in his extension of the Wilson-Wells technique; however, they evaluated the slope of the extrapolation by the method of least-squares instead of fitting it to Ramsay's predetermined • values (dependent on the assumption of a Lorentzian function for the true band shape and on a triangular slit function), and they omitted Ramsay's wing corrections (also dependent on a Lorentzian function). So, as was the case for the Wilson-Wells extrapolation, the Russell-Thompson technique is independent of any assumption regarding the true band shape, but it is quite sensitive to measurements at dilute concentrations. However, if a fair number of concentrations are measured, the resulting error should not be too serious.

Several studies of integrated intensities by the Russell-Thompson technique^{9,20,21} have shown that the extrapolation is essentially linear and that, if performed for several series of measurements made at different slit widths, it converges to a reasonably constant value for the true integrated absorption intensity. An interesting observation is that the slope of the Russell-Thompson extrapolation is larger than that predicted by Ramsay.^{9,20,21}

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5. Comparison of Methods

The Wilson-Wells extrapolation of B values presents certain advantages over the Bourgin extrapolation of A' values:

- (i) the extrapolation of B as a function of cl is approximately linear, especially if the band is wide enough, while that of A' is a shallow curve;
- (i1) since A' is independent of the resolving power and B is not, the use of a spectrophotometer with greater resolving power will assist the extrapolation of B but not of A'.

For these two reasons, the Wilson-Wells extrapolation of B values is more precise than the Bourgin extrapolation of A' values. However, an important weakness common to both extrapolation procedures is that they are quite sensitive to measurements at small concentrations or path lengths where experimental errors in determining band areas are apt to be large. In order to obtain satisfactory extrapolations, measurements at a considerable number of concentrations are necessary.

The most serious limitation of the Ramsay procedures is their dependence on a Lorentzian function for the true band shape. Although absolute integrated intensities obtained by Ramsay's methods are dependent on a Lorentzian function only to the extent of the wing correction and the slope of the extrapolation, both of which are fairly small quantities, the Lorentzian dependence may result in Significant error, since it has been shown that the assumption of a Lorentzian form for the true band shape of an infrared absorption band is uncertain and often produces unsatisfactory results.^{10,19} Since the Wilson-Wells and Russell-Thompson procedures are independent of any assumption regarding the true band

shape (i.e., they do not utilize wing corrections in evaluating the area under an absorption band, and they analyse the slopes of the extrapolations by the method of least-squares instead of fitting the slopes to predetermined values dependent on some assumption for the true band shape), it is felt that the Wilson-Wells and Russell-Thompson extrapolations are somewhat more reliable for evaluating the true integrated intensity of an infrared absorption band than the Ramsay procedures.

Also, even though the Wilson-Wells and Russell-Thompson techniques are sensitive) to measurements at dilute concentrations, it is expected that any error resulting from using the least-squares technique in evaluating the extrapolation slope (especially if measurements at a considerable number of concentrations are used in the intensity plot) will be less than the error incurred in utilizing Ramsay's e and \emptyset values which were determined assuming a Lorentzian function for the true band shape. Likewise, since the procedures for determining Ramsay's wing corrections are quite imprecise, ¹⁹ particularly in the case of overlapping bands, any error made in preserving the apparent intensity values in the Wilson-Wells and Russell-Thompson procedures is probably less than that incurred in the Ramsay methods by applying wing corrections.

However, an important advantage of Ramsay's procedures over the Wilson-Wells and Russell-Thompson extrapolations is that it is possible by means of Ramsay's calculated Ø and e values to estimate the true integrated, intensity from a single, measurement of the apparent integrated absorption intensity. Therefore, when a quick estimate of the true integrated intensity is required, Ramsay's methods are convenient techniques.

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As far as a comparison of the Wilson-Wells and Russell-Thompson procedures is concerned, no advantage of one technique over the other may be demonstrated because of their similarities: (1) the nature of the extrapolations [i.e., apparent integrated intensity as a function of concentration (Wilson-Wells) or apparent peak optical density (Russell-Thompson)] as long as the intensity measurements are taken over a concentration range in which the Beer's-Law plots are linear; (2) freedom from any assumptions regarding the true band shape; (3) analysis of extrapolation slopes by the method of least-squares; and (4) omission of wing corrections. So the two procedures were considered to be equivalent.

Chapter 2

BONDING IN METAL CARBONYLS

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Molecular Orbital Description of the Bonding

Before considering the infrared intensities of the carbonyl stretching vibrations in metal carbonyl complexes, it is necessary to briefly describe the bonding between a terminal carbonyl group and a transition metal atom and to consider the sensitivity of the M- β -O bond towards replacement of other carbonyl groups in the molecule by different ligands. A molecular orbital description of the bonding between a terminal CO group and a transition-metal atom is presented in several texts; $^{22-24}$ a simplified version of the bonding is sufficient for present purposes.

In the carbon monoxide molecule, the bonding is represented by a 0 system, consisting of a filled 0-bonding orbital and two 0-nonbonding lone pairs; and a π system, comprised of two mutually perpendicular filled π -bonding orbitals and the corresponding unfilled π -antibonding orbitals (Figure 2-1). When co-ordinated to a metal atom, the less electronegative atom of the CO group, carbon, donates lone-pair electron density to a 0-oriented metal orbital. However, since it is apparent from the general chemistry of carbon monoxide that the donor properties. of carbon monoxide are very slight (i.e., CO forms a weak complex with BH₃ and no compounds with metals in which there are no d electrons for back bonding), it is believed that additional π bonding makes a significant contribution to the molecular stability. Various





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Figure 2-1. Molecular Orbital Description of the Bonding in Carbon Monoxide.

^{*}F. A. Cotton and G. Wilkinson, "Advanced Inorganic Chemistry," 3rd ed, Interscience Publishers, Inc., New York, N. Y., 1972, p 685. 18

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o II C empty metal or sp hybrid filled orbital orbital

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Figure 2-2. Molecular Orbital Description of the Bonding in Transition-Metal Carbonyls."

"F. A. Gotton and G. Wilkinson, "Advanced Inorganic Chemistry," 3rd ed, Interscience Publishers, Inc., New York, N. Y., 1972, p 685. experimental techniques (x-ray crystallography, infrared spectroscopy, etc.,) provide physical evidence for the double bond character of the M-C bond.

To summarize then, the essentials of the bonding between a terminal CO group and a transition-metal atom are:

- (1) o donation from the unshared pair of electrons on carbon to an unfilled o orbital on the metal;
- (2) IT back-donation of electrons from the dT orbitals on the metal to the empty pT+-antibonding orbitals on the ∞ group.

Thus, the two components of the metal-carbon bond complement one another in the sense that charge is shifted in opposite directions in the two bond systems. π bonding has the offect of inducing still stronger 0-bond interaction between C and M and vice versa.

A pictoral description of the bonding is presented in Figure 2-2. For present purposes, the M-C-O skeleton will be considered linear because, although small deviations from linearity (rarely more than 5°) often wccur, they are not sufficiently large to effect present bonding arguments.^{4,25}

Vibrational Stretching Distortion of the CO Bond

Stretching the bond of a co-ordinated CO group will lead to a change in electron density within the M-C-O bonding system. Calculations using Slator atomic orbitals indicate that relatively little change in Oelectron density occurs during a small vibration of carbon monoxide, but that there are quite large changes in the M-electron density.^{4,26} It would seem, therefore, that M-bonding effects within a bound vibrating CO group probably predominate. The relative charge densities on the atoms in a co-ordinated CO group appear to be ${}^{b-}$ C \rightarrow O ${}^{b+}$. Thus, a vibrational atretching distortion most probably leads to a transfer of electron density C \rightarrow O, 25 If this C \rightarrow O electron transfer occurs and is located in the T system, then, as the CO bond is stretched, the empty pT orbitals will be increasingly lowered in energy and will localize more and more on the carbon atom. The expected result would be enhanced M \rightarrow C dT-electron donation.⁴ The converse of the above argument is also useful, i.e., that M \rightarrow C dTelectron donation should enhance C \rightarrow O pT-electron transfer.

It would appear then that the flow of electrons C->O in the π system of the CO group during a CO vibrational stretching distortion should serve as an indication of the availability of the metal d π -electron density. The dipole moment change created by this charge transfer will be reflected in the frequencies and in the intensities of the CO stretching absorptions of metal carbonyl complexes. Hence, the foregoing arguments may be usefully applied to a discussion of the effects of substitution of CO groups in metal carbonyl complexes by other ligands.

If a CO group in a metal carbonyl complex is replaced by a ligand of greater 0-donor but poorer π -acceptor ability, the effective nuclear charge of the metal atom will decrease, providing a greater amount of π -electron density for donation to the π^* orbitals of the remaining CO ligands. The resultant increase in the C-+O π -electron transfer during a ∞ bond stretching vibration and consequent weakening of the ∞ bond should result in decreased ∞ -frequencies and enhanced ∞ -intensities relative to those of the parent complex.

Conversely, replacement of a CO group in a metal carbonyl complex with a highly electronegative group, such as a halogen atom, will lead

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to an increased effective nuclear charge on the metal atom. This in turn will lead to a decrease in T-electron transfer to the remaining (O groups during a CO bond stretching vibration and hence to a strengthening of the CO bond. The expected result for such a substitution would therefore be increased CO frequencies and decreased CO intensities

The foregoing discussion will be used as a basis for a further investigation of the σ - and π -bonding abilities of various ligands in substituted metal carbonyls, since replacement of a CO group by a ligand of differing σ -donor and π -acceptor ability will be reflected in a change in CO-intensity relative to that of the parent complex.
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PART II

EXPERIMENTAL

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Chapter 3

6

1

PREPARATION METHODS

Physical Measurements

Carbon and hydrogen microanalyses were performed by the Schwarskopf Microanalytical Laboratory, Woodside, New York, except for the analysis of <u>fad</u>-tricarbonyltris(chlorodiphenylphosphine)molybdenum(O) which was done by Mr. N. J. Coville in the Department of Chemistry, McGill University, on a model 185 Hawlett Packard carbon-hydrogen-nitrogen analyzer equipped with a Cahn ratio electrobalance, model G, and a Honeywell recorder. Phosphorus and halogen microanalyses were carried out by Schwarzkopf.

Melting points were taken in unsealed capillaries on a Gallenkamp melting point apparatus design number 889339 and are uncorrected. A Zeiss Opton Nri21133 Abbe refractometer was used for refractive index mensurements. Values are reported for the D line of the spectrum of sodium; the temperature of determination was 20° . Infrared spectra were recorded in the metal carbonyl stretching region (2200-1800 cm⁻¹) on a model 337 Perkin Elmer double-beam grating spectrophotometer.

Materials

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Alfa Inorganics, Peverly, Massachusetts, was the source of most starting materials: hexacarbonylchromium(0), hexacarbonylmolybdenum(0), hexEcarbonyltungsten(0), decacarbonyldimanganese(0), m-cyclopentadienyltricarbonylmanganese(I), m-methylcyclopentadienyltricarbonylmanganese(I), trichlorophosphine, dichlorophenylphosphine, and chlorodiphenylphosphine.

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Triphenylphosphine, triphenylarsine, triphenylstibine, cycloheptatriene and bicyclo[2.2.1]heptadiene were obtained from the Aldrich Chemical Company, Milwaukee, Wisconsin. π -Cyclopentadienyldicarbonylthiocarbonylmanganese(I) was generously provided by Mr. A. E. Fenster²⁷ of the Department of Chemistry, McGill University.

A & C American Chemicals, Ville St-Laurent, Quebec, supplied the common solvents; <u>n</u>-hexane, benzene, methylenechloride, chloroform and carbon tetrachloride. The remaining reagents were purchased from various companies: acetonitrile from Brickman and Company, Montreal, Quebec; mesitylene from the J. T. Baker Chemical Company, Phillipsburg, New Jersey; methylcyclohexane from Eastman Organic Chemicals, Rochester, New York; and bromine and petroleum ether (bp 30-60°) from Anachemia Chemicals, Montreal, Quebec.

$lexacarbonylchromium(0)[cr(c0)_6]$

The white, crystalline solid was purified by sublimation (50°/0.1 mm Hg), mp 150-153°; lit.^{28a,29} 152-153°, 154-155°, 152-155°, 151-152°.

<u>n-Cyclopentadienylmanganese(I)</u> Complexes

1. π -Cyclopentadienyltricarbonylmanganese(I) [π -C₅H₅Mn(∞)₃]

Purification of the yellow crystals was accomplished by sublimation (50°/0.1 mm Hg), mp 73-77°; lit.^{28b,30a} 76.8-77.1°, 77°.

π-Methylcyclopentadienyltricarbonylmanganese(I) [(π-CH₃C₅H₄)Mn(CO)₃]
The yellow liquid was used as received, n²⁰D 1.5862; lit.^{28c,30b}
1.5868. Further purification was deemed unnecessary.

3. π -Cyclopentadienyldicarbonylthiocarbonylmanganese(I) $[\pi$ -C₅H₅Mn(∞)₂CS]

The yellow, crystalline solid was sublimed $(45^{\circ}/0.1 \text{ mm Hg})$ before use, mp 52-54°; lit.²⁷ 52-53°.

Group VI Metal-Olefin Carbonyls

1. Tricarbonylcycloheptatrienechromium(0) $\left[C_{7}H_{8}Cr(\Omega)_{3}\right]$

 $C_{7}H_{8}Cr(CO)_{3}$ was prepared according to King's³¹ method by refluxing hexacarbonylchromium(0) (4.4 g, 20 mmol) and excess cycloheptatriene (5 ml, 4.0 g, 44 mmol) in acetonitrile (50 ml). The resulting red crystals were purified by sublimation (100°/0.1 mm Hg), mp 126-128° dec; lit.^{28d},31-33 128-130° dec, 129-130°, after having first sublimed out (70°/0.1 mm Hg) any unreacted hexacarbonylchromium(0).

2. <u>Tetracarbonylbicyclo[2.2.1]hepta-2.5-dienechromium(0)</u> $[\underline{nor}-C_{7}H_{8}Cr(\infty)_{4}]$

<u>nor</u>- $C_{7}H_{8}Cr(\infty)_{4}$ was prepared by refluxing hexacarbonylchromjum(0) (5.0 g, 23 mmol) and excess bicycloheptadiene (8 ml, 7.3 g, 79 mmol) in methylcyclohexane (40 ml), as described previously in the literature.^{28e,34} Purification of the yellow, crystalline product was accomplished by sublimation (80°/0.1 mm Hg), mp 91-93°; lit.^{28e,34} 92-93°, after having first sublimed out (60°/0.1 mm Hg) any unreacted hexacarbonylchromium(0).

3. Tricarbonylcycloheptatrienemolybdenum(0) $[C_{\gamma}H_{8}Mo(\infty)_{3}]$

The reported procedure $^{28f, 32}$ for preparing $C_{7}H_8MO(\infty)_3$ by refluxing hexacarbonylmolybdenum(0) (26.4 g, 0.1 mol) and excess cycloheptatriene (50 ml, 40.3 g, 0.44 mol) in methylcyclohexane (200 ml) was employed. The resulting red crystals were dried (50°/0.1 mm Hg) for 12 hr to

remove any unreacted hexacarbonylmolybdenum(0) or other volatile impurities, and were sublimed $(100^{\circ}/0.1 \text{ mm Hg})$ before use, mp 98-102°; lit. $^{281}, 32, 35, 36$ 100.5-101.5°, 100°. Anal. Calcd for $C_{10}H_80_4Mo_3$ C, 44.1; H, 3.0; Mo, 35.3. Found: C, 44.4; H, 3.1; Mo, 35.2.

4. <u>Tetracarbonylbicyclo[2.2.1]hepta-2,5-dienemolybdenum(0)</u> $[\underline{nor}-C_{7}H_{8}Mo(\infty)_{4}]$

<u>nor-C₇H₈Mo(∞)₄ was prepared according to the published methods^{28f, 34}</u> by refluxing hexacarbonylmolybdenum(0) (26.4 g, 0.2 mol) and excess bicycloheptadiene (40 ml, 36.0 g, 0.4 mol) in methylcyclohexane (200 ml). The resulting yellow crystals were purified by sublimation (100°/0.1 mm Hg), mp 76-78°; lit.^{28f,33,34} 77-78°, 76-77°, after having first sublimed out ($60^{\circ}/0.1$ mm Hg) any unreacted hexacarbonylmolybdenum(0).

5. <u>fac</u>-Tricarbonyltris(acetonitrile)tungsten(0) [(CH_3CN)₃ $W(\infty)_3$]

The published method^{37,38} for preparing $(CH_3CN)_3W(CO)_3$ by refluxing hexacarbonyltungsten(0) (7.0 g, 20 mmol) in actonitrile (50 ml) was used. Purification of the yellow, crystalline product was not attempted; the crude form was adequate for subsequent reaction with the elefins.

6. Tricarbonylcycloheptatrienetungsten(0) $[C_7H_8W(\infty)_3]$

 $C_7H_8N(\infty)_3$ was prepared by refluxing <u>fac</u>-tricarbonyltris(acetonitrile)tungsten(0) (1.2 g, 3 mmol) and excess cycloheptatriene (2 ml, 1.6 g, 18 mmol) in <u>n</u>-hexane (50 ml) as described previously in the literature.³⁸ The resulting red crystals were purified by sublimation (100°/0.1 mm Hg), mp 114-118°; lit.³⁸⁻⁴¹ 115-121° dec, 117-118°, 110°, after having first sublimed out (70°/0.1 mm Hg) any hexacarbonyltungsten(0) present.

7. <u>Tetracarbonylbicyclo[2.2.1]hepta-2,5-dienetungsten(0)</u> $[\underline{nor}-C_{\gamma}H_{B}W(\infty)_{4}]$

The procedure³⁸ for preparing <u>nor-C₇H₈W(∞)₄ was the same as the</u> one used for the preparation of C₇H₈W(∞)₃ except that bicycloheptadiene was substituted for the cycloheptatriene. Purification of the yellow, crystalline product was accomplished as outlined above for the cycloheptatriene analogue, mp 88-91°; lit.^{38,39} 90-92°

Kolybdenum-Phosphine Carbonyls

1. Tricarbonylmesitylenemolybdenum(0) $[[1,3,5-(CH_3)_3C_6H_3]MO(CO)_3]$ The published method ^{42,43} for preparing $[1,3,5-(CH_3)_3C_6H_3]MO(CO)_3$ by refluxing hexacarbonylmolybdenum(0) (4.0 g, 0.015 mol) in excess mesitylene (20 ml, 17 g, 0.14 mol) was used. The infrared spectrum of the yellow, crystalline product was identical to that reported previously⁴³ for $[1,1,5-(CH_3)_3C_6H_3]MO(CO)_3$. No attempt was made to purify the product; the crude form was adequate for subsequent reaction with the phosphines.

2. <u>fac-Tricarbonyltris(trichlorophosphine)molybdenum(0)</u> [<u>fac-Mo(∞)₃(PCl₃)₃</u>]

 $fac-Mo(\infty)_3(FCl_3)_3$ was prepared by heating tricarbonylcycloheptatrienemolybdenum(0) (1.2 g, 0.002 mol) and excess trichlorophosphine (20 ml, 0.23 mol), as described previously in the literature.⁴⁴ The resulting yellow crystals were purified by recrystallization from petroleum ether (bp 30-60°), mp 82-85° dec; lit.⁴⁴ 85° dec. <u>Anal</u>. Calcd for $C_3P_3O_3Cl_9Moi$ C, 6.1; Cl, 53.9; P, 15.7. Found: C, 6.1; Cl, 53.3; P, 15.6.

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3. $\frac{\text{fac-Tricarbonyltris(dichlorophenylphosphine)molybdenum(0)}}{[\text{fac-Mo(CO)}_3[P(C_6H_5)Cl_2]_3]}$

The procedure for preparing $\underline{fac}-Mo(\mathfrak{O})_3[P(C_6H_5)Cl_2]_3$ was analogous to that of Poilblanc and Bigorgne⁴⁵ for preparing $\underline{fac}-Mo(\mathfrak{O})_3P(C_6H_5)_3$. Tricarbonylmesitylenemolybdenum(0) (0.3 g, 1 mmol) and excess dichlorophenylphosphine (1.0 ml, 1.3 g, 7 mmol) were dissolved in benzene and allowed to stand for 3 hr under nitrogen. The solution was then evaporated to dryness and the residue extracted with petroleum ether (bp 30- 60°) (4 x 50 ml), as outlined by Abel, Bennett and Wilkinson.⁴⁴ Filtering, concentrating (50 ml) and cooling (-78°) the extract deposited yellow crystals. After filtration, washing with petroleum ether (40 x 25 ml) and drying (25°/0.1 mm Hg), the crystals were identified as $\underline{fac}-Mo(\mathfrak{O})_3[P(C_6H_5)Cl_2]_3$ by a comparative infrared spectrum,⁴⁴ mp 115-117° dec; lit.⁴⁴ 106-110° dec.

4. fac-Tricarbonyltris(chlorodiphenylphosphine)molybdenum(0)

#54 2

$[\underline{fac}-Mo(\infty)_{3}[P(C_{6}H_{5})_{2}Cl]_{3}]$

The preparation of $\underline{fac}-Mo(\infty)_3[P(C_6H_5)_2Cl]_3$ was analogous to Poilblanc and Bigorgne's⁴⁵ preparation of $\underline{fac}-Mo(\infty)_3P(C_6H_5)_3$. A solution of tricarbonylmestitylenemolybdenum(0) (0.3 g, 1 mmol) and excess chlorodiphenylphosphine (1.0 ml, 1.2 g, 5 mmol) in benzene was allowed to stand for 3 hr under nitrogen, during which time yellow crystals deposited. Filtration, washing with benzene (4 x 20 ml) and drying (25°/0.1 mm Hg)⁴⁴ yielded the desired product, mp 155-157° dec; lit.⁴⁴ 148° dec. <u>Anal</u>. Calcd for $C_{39}H_{30}P_{3}O_3Cl_3MO$: C, 55.7; H, 3.6. Found: C, 55.6; H, 3.4.

Manganese-Mixed-Ligand Carbonyls,

1. Bromopentacarbonylmanganese(I) $[Mn(0)_5Br]$

 $Mn(CO)_5$ Br was prepared by reacting bromine (1.3 ml, 3.7 g, 23 mol) with decacarbonyldimanganese(0) (6.9 g, 18 mol) in carbon tetrachloride (100 ml), as described previously in the literature.^{28g,46} The resulting yellow-orange crystals were identified by a comparative infrared spectrum ^{28g,46} and were purified by sublimation ($60^{\circ}/0.1$ mm Hg).

2. <u>cis-Bromotetracarbonyl(triphenylphosphine)manganese(I)</u>

 $[\underline{cis}-Mn(\infty)_{4}[P(C_{6}H_{5})_{3}]Br]$

<u>cis-Mn(∞)</u> $[P(C_6H_5)_3]$ Br was prepared according to Angelici and Basolo's method⁴⁷ by reacting bromopentacarbonylmanganese(I) (0.54 g, 2 mmol) and triphenylphosphine (0.53 g, 2 mmoI) in chloroform (25 ml) at 40°. The orange-yellow product was recrystallized from <u>n</u>-hexane by dissolution at room temperature and recrystallization at dry-ice temperature, mp 130°. The infrared spectrum is identical to that reported previously⁴⁷ for <u>cis-Mn(∞)</u> $[P(C_6H_5)_3]$ Br. <u>Anal</u>. Calcd for C₂₂H₁₅PO₄BrMn: C, 51.9; H. 3.0. Found: C, 52.1; H, 3.0.

3. $\frac{\text{cis-Bronotetracarbonyl(triphenylarsine)manganese(I)}}{[\text{cis-Mn}(\Omega)_{4}[\text{As}(C_{6}!!_{5})_{3}]\text{Br}]}$

The procedure for preparing and purifying <u>cis-Mn(CO)</u> $As(C_{6}H_{5})_{3}Br$ was similar to that used for the preparation of <u>cis-Mn(CO)</u> $P(C_{6}H_{5})_{3}Br$. Identification of the triphenylarsine product (mp 130-134°) was confirmed by a comparative infrared spectrum. ⁴⁷ <u>Anal</u>. Calcd for C₂₂H₁₅AsO₄BrHn: C, 47.8; H, 2.7. Foundi C, 47.8; H, 3.0. 4.

<u>cis</u>-Bromotetracarbonyl(triphenylstibine)manganese(I)

 $\frac{[ois-Mn(\infty)_4[Sb(C_6H_5)_3]Br]}{[Sb(C_6H_5)_3]Br]}$

 $\underline{\operatorname{cis}}-\operatorname{Mn}(\operatorname{CO})_{4}[\operatorname{Sb}(\operatorname{C_{6}H_{5}})_{3}]$ Br was also prepared and purified according to the method described above for the triphenylphosphine analogue, mp 131-134°. A comparative infrared spectrum⁴⁷ was used for identification purposes. <u>Anal</u>. Calcd for $\operatorname{C_{22}H_{15}SbO_{4}BrHn_{1}}$ C, 44.0; H, 2.5; Br, 13.3. Found: C, 44.3; H, 2.6; Br, 14.2. Chapter 4

INTENSITY MEASUREMENTS

Instrumentation

Infrared spectra were recorded on a Perkin-Elmer model 337 doublebeam grating spectrophotometer equipped with a 70-min linear frequency scan (4000-1200 cm⁻¹, i.e., a recording rate of 40 cm⁻¹ min⁻¹), and a variable slit program (10 settings ranging from narrower than normal for high resolution scans to wide slit operation where high energy is required). The spectral slit width⁴⁸ in the metal carbonyl CO stretching region (2200-1800 cm⁻¹) ranged from 3.21 cm⁻¹ at 2200 cm⁻¹ to 2.55 cm⁻¹ at 1800 cm⁻¹; for the manganese thiocarbonyl CS stretching absorption, 1264 cm^{-1} , it was equal to 2.14 cm⁻¹ (Appendix B).

A Texas Instruments Servo/Riter II recorder, and later a Sargent-Welch model SRIG recorder, was installed to expand the frequency scale $(286 \text{ cm}^{-1} \text{ inch}^{-1})$ of the standard chart paper. As much frequency scale expansion as possible was employed because of the narrowness of the metal carbonyl peaks. On the Texas Instruments recorder, maximum expansion corresponded to a chart paper frequency scale of 5 cm⁻¹ inch⁻¹; on the Sargent-Welch instrument, it was equivalent to a chart paper frequency scale of 8 cm⁻¹ inch⁻¹. The carbon monoxide absorption at 2147.08 cm⁻¹ and the polystyrene peak at 1601.4 cm⁻¹ were used for calibration purposes.⁴⁹

In order to ensure accurate measurements, proper attention to double-beam balancing, zero transmission setting, etc., is essential.

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The employment of narrow slit widths is also important in the measurement of the sharp metal carbonyl absorptions.

Optics

Matched, sealed, standard-type liquid cells (1 mm and 0.1 mm path length) with NaCl windows were used for all measurements. International Crystal Laboratories (Irvington, New Jersey) calibrated the 1 mm cells by comparison of the benzene absorption doublet peak at 3.89 miorons with that of a known 1 mm cell (previously calibrated against a 0.2 mm cell, the path length of which had been obtained by means of interference fringes between 8 and 13 microns). The 0.1 mm cells were calibrated prior to each set of intensity measurements by means of interference fringes⁵⁰ between 4.8 and 5.2 microns; the path lengths were found to be the nominal values (determined by ICL by the interference fringe pattern method from 4.5 to 8 microns) + 2% (Appendix C).

Sample Preparation

In general, solutions were prepared from individually weighed samples; however, for the three manganese-mixed-ligand carbonyls, volumetric dilution techniques were employed giving results which fitted Beer's-Law plots for the individually prepared samples fairly well.

The solid material was weighed on a Mettler balance accurate to \pm 0.1 mg, transferred carefully to a 50-ml volumetric flash and then the solution was diluted accurately to the mark with solvent. Whenever possible, <u>n</u>-hexane was used as solvent because of the optimum resolution of the metal carbonyl absorption spectra in this solvent.⁵¹ However, where low solubility of the complexes in <u>n</u>-hexane precluded its use, methylene chloride (molybdenum-phosphine carbonyls and metal-cycloheptatriene

carbonyls) or carbon disulphide [π -cyclopentadienyldicarbonylthiocarbonylmanganese(I)] was substituted. Spectrograde n-hexane, methylene chloride and carbon disulphide were obtained from Fisher Scientific Company, Fair Lawn, New Jersey and were used without further purification.

Recording Procedure

Spectra were recorded as soon as possible after complete dissolution of the sample, usually within 30 min. Decomposition of the samples was minimized by using fresh solutions protected from light by aluminum foil. Generally, decomposition was not extensive enough (usually less than 1%) to affect the intensity results.

Spectra were recorded for at least four, usually five, different concentrations of each sample, providing a smooth distribution of bands within the 15 to 65% transmittance range (i.e., $1.89 \ge \ln(T_0/T)_{v_{max}} \ge$ 0.427). This required sample concentrations ranging from 0.1 to 15.5 mM. The relative error in the integrated intensities is expected to be less than 10% under these conditions.⁵² Four individual recordings were taken at each concentration and the average values of $\ln(T_0/T)_{v_{max}}$ and of the area for each band at each concentration were used in the intensity calculations.

Area Determination

The area under the absorbance curve

$$\int_{\text{band}} \log_{10}(T_0/T)_v dv$$
 (21)

where T $_{V}$ and T $_{V}$ are the apparent intensities of the incident and transmitted radiation when the spectrophotometer is set at frequency V, was determined by two techniques, depending on which recorder was in use (Appendix A). The Sargent-Welch instrument presented the spectrum in linear absorbance; this permitted direct evaluation of the integral in (21) by cutting out the band profile and weighing it. Recordings were made on Sargent recorder chart paper catalog No. S-72167 and a Sartorius-Werke model AG balance accurate to ± 0.01 mg was used for the weighings. The paper was calibrated for each spectrum and was found to be quite uniform.

On the Texas Instruments recorder, a linear transmittance spectrum was obtained; here conversion of the experimental absorption band from a linear transmittance to a linear absorbance curve necessarily preceded evaluation of the integral in (21). Readings were taken usually at intrements of about 1/7 the true half-intensity band width (Appendix B), providing an equidistantly tabulated function of $\log_{10}(T_0/T)_v$ versus v. Evaluation of the aren under the absorbance curve was then accomplished by the application of Simpson's rule. This procedure is quite reliable for reasonably broad bands. However, for very narrow bands, where the increments of the argument values were sometimes greater than $\frac{1}{4}$ the halfintensity band width, the errors encountered were somewhat larger, making the procedure less useful in these cases. A comparison of intensity results for Cr(∞)₆ determined by the weighing method and by Simpson's rule (Table 5-1) affords good evidence for the equivalence of the two techniques.

The integration was generally performed over an interval of approximately five times the true half-intensity band width (Appendix B) on either side of the band centre, $v_{\max} \pm 5\Delta v_{\frac{1}{2}}^{t}$ (i.e., the region from the band centre to the point at which the absorption appeared complete). The area under the solvent background curve was subtracted, but no

correction for the area under the wings was applied since the wing correct, signa provided by Namsay⁸ seemed to be everly large for the sharp CO stretching bands of most substituted metal carbonyls. In order to convert the areas to the maperian base (i.e., to obtain the area under the optical density curve, $\int_{\text{band}} \ln(T_0/T)_v dv)$, the results were multiplied by 2.303. Overlapping bands were resolved by extrapolation procedures assuming symmetric band shapes. Such an extrapolation for $C_7H_8Mo(CO)_3$ is illustrated in Figure 4-1.

Treatment of Data

Three methods were employed for the determination of the true integrated absorption intensity

$$A = \frac{1}{cl} \int_{band} \ln(\frac{I_0}{I}) v^{dv}$$
 (22)

where c is the concentration of solute in moles per liter, l is the cell path length in cm, and I_{0y} and I_{y} are the incident and transmitted intensities of monochromatic radiation of frequency v. In the first method, the Wilson-Wells extrapolation,⁷ A is obtained by extrapolating measurements of apparent integrated absorption intensity

$$B = \frac{1}{cl} \int_{band} \ln(\frac{T_o}{T})_v dv$$
 (23)

made at a number of concentrations to sero concentration. The plots for $\pi-C_{5}H_{5}$ $Fn(CO)_{2}$ CS are shown in Figure 4-2.

The second method used to determine A is the Russell-Thompson extrapolation.⁹ In this procedure, A is obtained by plotting B as a function of apparent peak optical density, $\ln(T_0/T)_{v_{max}}$ and extrapolating to $\ln(T_0/T)_{v_{max}} = 0$. The extrapolations for $\pi - C_5 H_5 Nn(CO)_2 CS$ are illustrated in Figure 4-3. 「「「「「「「」」、「「」」、「」、「」、「」、



 $C_{\gamma}H_{8}Ho(\infty)_{3}$ by Extrapolation Procedures.

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Figure 4-2. Apparent Intensity versus Concentration Plots for the CO and CS Vibrational Modes in the $T-C_5H_5Mn(CO)_2CS$ Complex. CO Modes: e, A' Vibration; s, A" Vibration. CS Mode: A, A' Vibration.



Figure 4-3. Apparent Intensity versus Apparent Peak Optical Denaity Plots for the CO and CS Vibrational Modes in the n-C_H_Mn(CO)2CS Complex. CO Modes: •, A' Vibration; , A" Vibration. CS Mode: A, A' Vibration.

In both extrapolation procedures, the data were analyzed by the method of least-squares using a program written for the IBM 360/65 computer (Appendix A). Since the number of observations was small (usually five), a statistical treatment of data was considered meaningless, ⁵³ and so all points were retained in the analyzes and were assigned equal weight. For the most part, negative slopes which generally increased in magnitude with increasing half-intensity-width were obtained by both extrapolation techniques (Appendix R).

The third method used to determine A is the method described by Ramsay and colleagues^{8,18} for partially overlapping band systems. In this method, A is determined simply by increasing the mean of the B values by 2% (Appendix A). B values for at least four (usually five) different concentrations were used in the calculation of the mean.

A comparison of intensity results determined by the three methods outlined above (Appendix C) shows that the values obtained by the two extrapolation procedures agree very closely (usually to within 1%); this affords good evidence for the equivalence of the two techniques. However, the extrapolation results tend to be somewhat higher than the results obtained by Ramsay's method (generally about 7% higher). A similar increase in A value (over that predicted by Ramsay on the basis of the assumption of a Lorentz band contour and a triangular slit function) was noted by Russell and Thompson⁹ with their extrapolation method.

In conclusion, therefore, it is felt the two extrapolation procedures are the most reliable methods for determining the true integrated intensity of an infrared absorption band. However, an important advantage of Ramsay's method is that it is possible by means of his relation (i.e., A = 1.2 B)

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to estimate the true integrated intensity from a single measurement of the apparent integrated absorption intensity. Therefore, when a quick estimate of the true integrated intensity is required, Ramsay's method for overlapping band systems is a convenient technique.

Results

Complete details of the integrated intensity calculations are presented in Appendix C. The tabulated data used in the discussion were extracted from this Appendix.



Chapter 5

INFRARED INTENSITIES IN NETAL CARBONYL COMPLEXES

Interpretational Approach

Integrated infrared intensities of the carbonyl stretching absorptions of metal carbonyl compounds and their derivatives were investigated by Noack⁵⁴ as far back as 1962; however, the results of relatively few studies have appeared subsequently.^{25,55-72} This lack of activity is unfortunate, since the intensities of the CO stretching vibrations of metal carbonyl complexes can provide very useful information concerning the electronic properties of the molecules.

Metal carbonyl ∞ stretching vibrations are many times more intense than those of free ∞ (i.e., the intensity of the allowed ∞ stretching mode in $\operatorname{Cr}(\infty)_6$, $\sim 65 \times 10^4 \text{ M}^{-1} \text{ cm}^{-2}$, is approximately 110 times the value of the fundamental for ∞ in the gas state, $0.58 \times 10^4 \text{ M}^{-1} \text{ cm}^{-2}$).²⁵ Explanations for such intensity enhancement have been given in terms of π electron transfer from the metal to the ∞ groups.^{25, 62} Stretching of a ∞ bond lowers the energy of the $p\pi^*$ orbitals which act as acceptors for metal $d\pi$ -electron density (<u>vide supra</u>, Chapter 2). The dipole moment change created by this charge transfer accounts for the large increase in intensity. Since the effects on the ∞ intensity due to variations in o-bonding during a ∞ stretching vibration are very small,²⁵ it is assumed that they can be ignored and that the intensity of metal carbonyl ∞ stretching vibrations is determined solely by π -bonding effects. This assumption appears to be in accord with many of the results which have been published.

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Thus, the CO intensity of a metal carbonyl complex can be taken as a measure of the extent of π -bonding present in the molecule. Replacement of one or more of the CO groups in a metal carbonyl complex by other ligands of varying G-donor and π -acceptor ability should result in a change of π bonding to the remaining CO groups (vide supra, Chapter 2). This should in turn lead to different values for the CO intensities. So, it should be possible to gain an appreciation of the sensitivity of metal carbonyl complexes to substitution by studying the integrated intensities of the CO stretching absorptions of the various molecules.

Unfortunately, it is not clear how a change in π -electron density in a metal carbonyl complex will be distributed among the various CO stretching modes, since both electronic coupling among several CO modes or vibronic contributions to individual CO modes often occur. Thus, a method of analyzing the measured intensity data wherein coupling could be disregarded was described.^{4,25} The total integrated CO intensity of any particular metal carbonyl complex is independent both of the angle between any of the symmetry related dipole moment derivatives and the symmetry axis and of any vibrational coupling, provided that the bond moment derivatives of all carbonyl groups in a molecule are identical and that there is only electronic interaction between the carbonyl groups. These requirements are fairly reasonable^{4,25} and appear to be quite realistic for most of the metal carbonyls and their derivatives which have been investigated.

Dividing the total absolute integrated CO intensity of a metal carbonyl complex by the number of carbonyl groups present in the molecule affords the integrated intensity per carbonyl group, or specific intensity. ⁵⁴ Using specific CO-intensity data, it is possible to compare

CO-intensities of metal carbonyl complexes of varying molecular symmetries and with differing numbers of carbonyl groups.⁴ In their recent review. article on the CO-intensities of metal carbonyl complexes, Kettle and Paul⁴ utilized specific CO-intensities in order to reduce the published data (reported using several systems of presentation) to a common basis. In this manner, a ready comparison of existing results could be made.

The work in this thesis constitutes an additional study of the integrated intensities of the carbonyl stretching absorptions of metal carbonyl complexes. The ∞ -intensities of the individual stretching vibrations as well as the specific ∞ intensity results for the following series of compounds will be considered: $Cr(\infty)_6$, $(Cycloheptatriene)M(\infty)_3$ and $(Bicycloheptadiene)M(\infty)_4$ (M = Cr, Mo, W). <u>fac</u>-Mo(∞)₃[P(C₆H₅)_{3-x}Cl_x]₃ (x = 1-3), <u>cis</u>-Mn(∞)₄LBr [L = P(C₆H₅)₃, As(C₆H₅)₃, Sb(C₆H₅)₃], and Π -RC₅H₄Mn(∞)₃ (R = H, CH₃) and Π -C₅H₅Mn(∞)₂CS. The first measurement of the infrared intensity of a thiocarbonyl stretching mode in the metal thiocarbonyl complex, Π -C₅H₅Mn(∞)₂CS, will also be discussed. The features of Kettle and Paul's report which are pertinent to the research carried out in this thesis will be reviewed at the same time; where applicable, the results of more recent publications will be summarized as well.

Hexacarbonylchromium(0)

The $Cr(\infty)_6$ molecule is octahedral (Figure 5-1) and possesses O_h symmetry. Group theory²³ predicts one infrared-active CO stretching mode (T_{1u}) for compounds of this type (Figure 5-2). As expected, the infrared spectrum of $Cr(\infty)_6$ displays a single isolated peak in the CO stretching region.

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Figure 5-1. Structure of the $Cr(00)_6$ Molecule.



Figure 5-2. Infrared-Active CO Stretching Vibration of the $Cr(CO)_6$ Molecule.

Since single peaks which are essentially free of neighbouring absorptions make convenient standards for comparison of intensity measurement techniques, we measured the integrated intensity of the CO stretching mode of $Cr(\infty)_6$ to check the reliability of our measuring techniques (Tables 5-1 and 5-2). The first point that emerges from an examination of the results is that the two integration methods used to measure band area produce almost identical results. Thus, integration of band area by the weighing method and by Simpson's rule are considered to be essentially equivalent procedures.

Second, if our specific CO intensity results are compared with the published values for $Cr(\infty)_6$ (Table 5-5), it is evident that our results obtained by Ramsay's method, 4.71 and 4.69 x 10^4 M⁻¹ cm⁻² (integration performed by Simpson's rule and the weighing method, respectively), agree very closely with the reported values measured in the same solvent (nhexane), 4.65^{54} and $4.61^{60} \times 10^4 \text{ M}^{-1} \text{ cm}^{-2}$. Since Abel and Butler's result was also determined by using one of Ramsay's methods (Noack's 54 value is not corrected for the finite slit width of the spectrophotometer and, as such, is only an apparent intensity and so should tend to be somewhat lower than the corrected results), it is apparent that our technique for determining integrated intensities by Ramsay's method compares well with the Ramsay techniques of other researchers, and so our Ramsay results may be reliably compared with equivalent values in the literature. Also, the specific intensity results obtained by the Wilson-Wells and Russell-Thompson extrapolations (Table 5-2) tend to be somewhat higher than both our Ramsay values and the published results. This is the expected trend because Russell and Thompson? noted a similar increase

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	Infr	ared Frequenc	y, Half-Inteni the CO Strete	sity Band Wi chirg Yibrat	dth, and Absolute In ion of $Cr(CO)_6$ in <u>n</u> -	Regrated Absorption Regrate Solution	Intensity for	6
			V BAX	Cu ^{t•}		E,** 10 ⁴	H ⁻¹ ca ⁻²	
Complex	Integration method	Normal mode	ca ⁻¹	cm ⁻¹	Ransay	Vilson- Vells	Russell- Thompson	Neas
در ه)و	Sizpson's	Tlu	1986.9	1.92	28.3 (* 0.6)	29.8 (* 1.9)	29.8 (* 2.0)	29.3 (+ 1.5)
ဇ . (ထ) ₆	Weighing -	Ťlu	1997.2	2.00	28.2 (* 0.6)	29.9 (* 1.3)	30.1 (* 1.4)	29.4 (* 1.1)

"These are the true half-intensity band widths computed from the observed values by applying Ramsay's corrections for the finite slit width of the instrument [D. A. Ramsay, <u>J. Am. Chem. Soc.</u>, <u>74</u>, 72 (1952)].

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"Defined as $\frac{1}{cl} \int_{band} \log_{10} (\frac{\Gamma_0}{l})_v dv$, where c equals the concentration of solute in H^{-1} , 1 is the cell path length in cm, and Γ_{0y} and Γ_{y} are the incident and transmitted intensities of monochromatic radiation of frequency V. Bassay: E is computed by taking the mean of the enserved intensities (defined as $\frac{1}{cl} \int_{band} \log_{10} (\frac{\Gamma_0}{T})_v dv$, where T_{0y} and T_y are the observed intensities of the incident and transmitted radiation when the spectrophotometer is set at frequency V) obtained at a number of concentrations and increasing it by 24. The numbers in parentheses are the standard deviations from least-squares analyses of the intensity versus concentration plots. Reseall-Thompson: E is determined by extrapolating the observed intensities taken at a number of concentrations from least-squares analyses of the intensity versus peak optical density plots. Reseall-Thompson: E is determined by extrapolating the observed intensities taken at a number of concentrations to sero apparent seek optical density. The numbers in parentheses are the standard deviations from least-squares analyses of the intensity versus peak optical density plots. Reseal: the E value is the mean for the three methods of determination. The numbers in parentheses are the standard deviations from least-squares analyses are the average standard deviations for the three methods.

TAPLE 5-1

Specific Intensity for the ∞ Stretching Vibration of $\operatorname{Cr}(\infty)_6$

Tom 1			Specific Intensity Wilson-	y, 10 [°] K ⁻¹ cm ⁻² Massil-	
8	vent	Ranary	¥ells	Thompson	
]		4.71 (± 0.10)	4.97 (± 0.31)	4.97 (± 0.34)	4 , 8
희	enane	4.6 9 (± 0.10)	4.93 (± 0.21)	5.01 (± 0.24)	+) 8,4

burdes the integrated interstry per CD group, $\{\mathbf{z}_{ij}^{\prime}|_{n}$, where n is the number of CD groups present in the molecule and \mathbf{z}_{ij}^{\prime} seats the istegrated intensity of the 1th normal mode. A definition of 2 and descriptions of the methods for obtaining it are provided in Table 5-1.

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in integrated intensity with their extrapolation procedure over the value predicted by Ramsay on the basis of the assumption of a Lorents band contour and a triangular slit function. Therefore, it is reasonable to assume that our extrapolation techniques are comparable to other extrapolation techniques.

Group VI Metal-Olefin Carbonyls

Early in the study of the ∞ intensities of metal carbonyl complexes, Beck and Nitzmann⁵⁵ observed that the specific ∞ intensities of isoelectronic and isostructural series of metal-carbonyls increase steadily from cationic through neutral to anionic species (Table 5-3). The most striking example of this trend is for the series $Nn(\infty)_6^+$, $Cr(\infty)_6$ and $V(\infty)_6^-$ (Table 5-3). Such trends are thought to reflect the greater availability of metal dm-electron density with increasing negative charge. A considerable body of data supporting Beck and Nitzmann's original observations is now available (Tables 5-3 and 5-4).

Another early observation was that for related species in a given group, an increase in atomic number of the transition metal atom results in an increase in the CO intensities. This observation has also been consistently reported (Tables 5-5 to 5-7) and is in good agreement with T-bonding arguments since one would expect the availability of metal dT-electron density to increase with the size of the metal atom electron cloud. In the case of $Mn(CO)_6^+$ and $Re(CO)_6^+$ (Table 5-6) and some complexes containing iodo ligands (Table 5-6), however, some anomalous behaviour has been reported.

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A comparison of CO-frequency and band-width data for the two series of group VI metal-olefin carbonyls studied in this work shows that there





Figure 5-4. Infrared-Active CO Stretching Vibrations of the C₇H₈H(CO)₃ Molecule (H = Cr. Ho, W).

·	and Isostructural Se	ries of Metal-Carbonyle"		. A
Complex	Solvent	Specific Intensity ^b 10 ⁴ H ⁻¹ cm ⁻²		Beference
N1(00)4	Tetrahydrofuran	4.95		• '
ω(∞) ,	Diethyleneglycol diethylether	7.50		¢
Hn(00)6+	Tetrahydrofuran	2.52		•
cr(∞) ₆	Tetrahydrofuran	6.10		•
۷(∞) ₆	Tetrahydrofuran	9.47	Ø	e
$Mn_2(\infty)_{10}$	Tetrahydrofuran	3.85		,
$cr_2(\infty)_{10}^{2-}$	Tetrahydrofuran	4.95		٩.
Re(ŵ) ₆ +	Tetrahydrofuran	· 2.11		•
w(∞)6	Tetrahydrofuran	7.75	3	•,

Specific Intensities for the CO Stretching Vibrations of Some Isoelectronic

^AS. F. A. Kettle and I. Paul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. West, Ed., Academic Press, Inc., New York, N. Y., 1972, pp 224, 230.

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^bAs defined in Table 5-2.

V. Beck and R. E. Mitsmann, Z. Maturforach., 178, 577 (1962).

^dR. M. Wing and D. C. Crocker, <u>Inorg. Chem.</u>, <u>6</u>, 289 (1967).

Specific Intensities of the ² CO Stretching Vibrations of Some Isosloctzonia and Isostructural Series of Notal Carbonyl Complexes ²					
Complex	Solvent	Specific Intensity ^b 10 ⁴ N ⁻¹ cm ⁻²	Neference		
Mn(00)5C1	Chloroform	3.06	•		
ന്തുംപ_	Chloroform	4.47	•		
Mn(00) ₅ Br	Chloroform	3.76	ت ا		
Cr(00)5Br	Chloroform	4,52	ł 9		
ي Kn(00)51	Chloroform	4.23	•		
α _r (α) ₅ 1 ⁻	Chloroform	6.28	• U		
Ne(O) ₅ Br	Chloroform	4.25	•		
W(@)5 ^{Br"}	Chloroform	6.08	•		
Re(CO)gI	Chloroform	4.19	•		
¥(∞),51-	Chlorofora	6 .20	٠		
dig-Fe(CO)aBra	D-Hexane	1.84	, 4		
019-Mn(00)4Br2	Chloroform	3.47	•		
<u>eis</u> -Fe(00) ₄ I ₂	Chloroform ,	2.51	٠		
019-Hn(00)412	Chloroform	3.99	.		
Co(CO)3110	Tetrahydrofwran	4.15	•		
r.(00),10°	Methyleneglycol diethylether	5.16	•		

⁸5. F. A. Kettle and I. Paul in "Advances in Organometallie Chemistry," Vel. X, F. G. A. Stone and R. West, ML. Academic Press, Inc., New York, N. Y., 1972, pp 220-226. ^bAs defined in Table 5-2.

^CE. W. Abel and I. S. Butler, <u>Trans, Faraday Soc</u>., <u>63</u>, 45 (1967).

4. Roack, <u>Helv. Chin. Acta</u>, 45, 1847 (1962)

"V. Bock and R. E. Mitzmann, E. "aturforsch., 177, 577 (1962).

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· ·	Specific Intensities for the CD Stretching Vibrations of a Series of Group VI Metal-Carbonyls in Various Solvents ^a .						
			Specific Intensit	y, ^b 10 ⁴ H ⁻¹ cm ⁻²	ζ		
Complex	n-fexane ^C	<u>p-::exame</u> d	Tetrachloro- ethylone ^d	Carbon tetrachloride [®]	fetrahydro furan [®]		
œ(œ) ₆	4.61	4.65	5.09	5.14	6.10		
‰(∞) ₆	5-39	5.01	² 5.19				
¥(∞)	5.34	5.53	5-37	6.62	7.75 2		

⁸S. F. A. Kettle and I. Paul in "Advances in Organometallic Chemistry," Vol. I, F. G. A. Stone and B. West, Ed., Academic Press, Inc., New York, N. Y., 1972, p 224.

bAs defined in Table 5-2.

& C. V. Abel and I. S. Butler, Trans. Fareday Soc., 63, 45 (1967).

d. Noack, <u>Helv. Chis. Acts</u>, 45, 1847 (1962).

"W. Bock and R. E. Mitsmann, Z. Waturforsch., 173, 577 (1962).

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Complex	Solvent	Specific Intensity ^b 10 ^b H ⁻¹ cm ⁻²	Baferene
Kn(@)6	Totrahydrofuraa	2.52	•
Re(∞) ₆ +	Tetrahydrofuran	2.11	۰ •
œ(∞) ₅ ¤(c ₆ × ₅) ₃	<u>n-Hexane</u>	5.25	4
10(00)5r(C6115)3	D-Hexane	6,14	4
F=CsHsKn(∞)3	Carbon tetrachloride	4.65	٠
n-05H5He(00)3 -	Carbon tetrachloride	5.36 /	٠
(n-c6K500C5H4)Hn(00)3	Carbon tetrachloride	4.59	٠
(n-C6H500C5H4)No(00)3	Carbon tetrachloride	° 5.14	٠
h(@)s ^{hr} ·	Chloroform	3.31 (3.76)	4 (f)
le(∞) ₅ Br	Chloroform	4.15 (4.25)	4 (r)
ín(∞) _s I	Chloroforn	4.23	f
No(00)51	Chloroform	4.19	, t
∝(∞) ₅ c1 [−]	Chloroform	4.47	f
((00) ₅ c1 -	Chloroform ,	6,24	r
≈(∞) _s Br	Chleroform	4. 52	e f
v(@) ₅ ar"	Chloroform	6.08	f
αr(ω) _s 1	Chloroform	6.20	1. 1
w(∞),1	Chloroforn	6.20	

"8. F. A. Kettle and I. Paul in "Advances in Organometallic Chemistry," Vol. X, P. G. A. Stone and R. Nest, Ed., Academic Press, Inc., New York, W. Y., 1972, pp 224, 225, 227.

As defined in Table 5-2

⁰W. Bock and R. E. Mitsmann, <u>Z. Maturforach., 178</u>, 577 (1962).

^dT. L. Brown and D. J. Darenabourg, <u>Inorg. Chem.</u>, <u>6</u>, 971 (1967).

A. N. Nesneyanov, G. G. Dvoryantseva, Kh. P. Pustyl'nik, Yu. H. Sheinker, H. E. Kelebeva, and K. M. Anleinov, Dokl. Akad. Hauk 5253, 174, 368, (1967); Acad. Sci. USSR (hem, 174, 557 (196

¹E. V. Abel and I. S. Butler, <u>Trans. Taraday Soc.</u>, <u>63</u>, 45 (1967).

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LABLE	_ <u>}*7</u>

Specific Intepsities for the CO Strutching Vibrations of Some Section of Felynuclear Metal-Garbonyle ⁴						
Conplex	Solvent	Bpecific Intensity ^b 10 ⁴ M ⁻¹ cm ⁻²	Reference			
Mag(00)10	Cy clohézane	2.87	•			
102(00)10	Cyclohoxane	3.57	6 b			
m ₂ (00) ₁₀	Carbon totrachiorido	3.36	4			
RoKn(03)10	Carbon tetrachloride	. 3. 44	4			
Ne ₂ (00) ₁₀	Carbon tetrachloride	J.6J	4			
n=C_N_NoKn(00)g	Carbon tetrachloride	3.60				
n-0,H,NoRe(00)8	Carbon tetrachloride	4.08	4			
n-C ₅ H ₅ VHn(00) ₈	Carbon Setrachloride	.). ?4	¢			
n-c_H_vne(00)8	Carbon tetrachloride	4,15	4 `			
n-C5H5Nonn(00)8	Carbon totrachloride	3.68	4			
n-05H5Wn(00)8	Carbon tetrachloride	3.74	4			
n-C ₅ N ₅ KoRe(00) ₈	Carbon tetrachloride	4.08	1			
n-C5H5VRe(00)8	Onrben tetrachloride	4,15	đ.			
NCr2(00)10-	r Tetrahydrofuran	4.95	· C			
HW2(00)10-	Totrahydrofuran	6.33	٠			

⁶E. T. A. Kettle and I. Paul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. Vest, Ed., Academic Press, Inc., New York, W. Y., 1972, p 230.

"As dofined in Table 5-2.

^CR. H. Wing and D. C. Crocker, <u>Inorg. Chem.</u>, <u>6</u>, 289 (1967).

⁶A. N. Yesseyanev, G. G. Dvoryaniseva, Yu. Y. Sheinker, N. M. Kolobeva, and K. N. Amielnev, <u>Dokl. Akad. Mauk SSIN, 169</u>, 843 (1966): <u>Acad. Sci. USSN Chem.</u>, 169, 751 (1966). is little change in either variable in going from chromium to tungsten (Tables 5-8 and 5-9). Because of this, a study of the OD-intensities of thy two series of compounds was undertaken to see whether the expected increase in CO-intensity with increasing size of the metal atom would occur or whether the two series of metal-olefin carbonyls would be further exceptions to the general trend.

Owing to the incomplete delocalisation of the H-electron system in cycloheptatriene. Free rotation about the metal-ring axis does not occur³² (i.e., cycloheptatriene-metal complexes have the ring fixed with respect to the carbonyl groups). The structure of the $O_2H_BM(CD)_3$ (M = Or, Mo, W) molecules is a plano-stool type arrangement (Figure 5-3), with the exact disposition of the carbonyl groups relative to the ring as yet undecided. The symmetry is O_1 for which group theory predicts three infrared-active (D stretching modes (3A). Each carbonyl group should give rise to a symmetric A stretching mode of reasonable intensity (Figure 5-4). As expected, infrared spectra exhibiting three well-resolved bands of comparable intensity were obtained.

Noiseulas of the type <u>nor-C₇H_HM(CO)₄ (M = Cr. Mo. W) have an</u> octahedral structure (Figure 5-5) and possess formal G_{2v} symmetry. Four infrared-active CO stretching modes (2A₁ + B₁ + B₂) would be expected for such an arrangement. The <u>trans</u> pair of carbonyls should produce a weak $A_1^{(2)}$ mode (which may be more intense than anticipated because of coupling with the other A_1 mode) and a strong antisymmetric B_1 stretching mode, while the other two carbonyls should give rise to a symmetric $A_1^{(1)}$ mode and to an antisymmetric B_2 stretching mode of comparable intensity (Figure 5-6).





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Figure 5-6. Infrared-Active CO Stretching Vibrations of the <u>nor-C₇H₈M(CO)₄ Molecule [M = Cr. Mo, W].</u>

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The spectrum of $\underline{nor}-C_{\gamma}H_{\beta}Cr(\infty)_{\mu}$ does display four well-resolved bands in the CO stretching region as expected; however, the spectra of the molybdenum and tungsten analogues exhibit only three distinct peaks. It is likely that the strong middle band in the spectra of these two compounds actually consists of two coincident absorptions. The CO-intensity data for the series (Table 5-9) help to corroborate this view; it is clear that the intensity of the strong middle band of both the molybdenum and tungsten complexes is comparable to the sum of the intensities of the two middle bands of the chromium complex.

Probable assignments of the four peaks in the <u>nor-C₂H₈Cr(CO)₄ spectrum in order of decreasing frequency are $A_1^{(2)}$, $A_1^{(1)}$, B_1 , and B_2 , respectively (Table 5-9). The highest frequency band is readily assigned to the $A_1^{(2)}$ mode since the local dipoles of the $A_1^{(2)}$ mode have their similar poles facing one another and hence there should be some cancellation giving rise to a weak absorption. The assignment of the B_1 mode to the absorption at 1944 cm⁻¹ is supported by the CD-intensity data; the B_1 mode should give rise to the most intense band and the 1944 cm⁻¹ absorption is in fact the most intense. The two remaining unassigned peaks at 1959 cm⁻¹ and 1915 cm⁻¹ were assigned to the $A_1^{(1)}$ mode and to the B_2 mode, respectively, since symmetrical stretching modes of vibration are expected to be at higher frequencies than the corresponding asymmetrical modes of vibration. These assignments are also supported by previous force constant calculations on isostructural molecules, 73, 74 e.g., <u>cis</u>-Fe(CO)₄X₂.</u>

The CO-intensity data for the group VI metal-olefin carbonyls are presented in Tables 5-8 to 5-10. From a comparison of the specific CO-intensities of the metal-olefin carbonyls with the values for the

TAPLE	5-8
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		VBAX	∆v ^t •	. /	E,** 10 ⁴	R ⁻¹ cm ⁻²	
Cosplex	Normal	ca ⁻¹	ca ⁻¹	Bancay	Vilson- Vells	Russell- Thompson	Zeen
с _{7^H3} ст(<i>Ф</i>) ₃	A ⁽¹⁾	1982.1	13.1	6.12 (* 0.22)	6.93 (+ 0.53)	7.14 (+ 0.75)	6.74 (+ 0.50
с _{уна} ко(а),	A ⁽¹⁾	1983. ₆	11.9	7.23 (* 0.22)	7.62 (+ 0.60)	7.63 (* 0.71)	7.49 (* 0.51
с ₇ н ₃ ч(∞) ₃	A ⁽¹⁾	1991.4	11.0	6.72 (* 0.22)	7.90 (* 0.45)	8.09 (- 0.56)	7.57 (* 0.41)
<i>د</i> ې ^{يړ} وک ر (۵۵)	A ⁽²⁾	1916.0	19.1	6.65 (* 0.17)	7.50 (* 0.45)	7.64 (* 0.56)	7.27 (* 0.39
C,HgHo(OO)	A ⁽²⁾	1912.,	22.0	7.93 (* 0.35)	9.17 (* 1.07)	9.24 (* 1.33)	8.75 (* 0.92
د _ع برهه(۳)	A ⁽²⁾	1908. ₃	20.6	7.50 (* 0.13)	8.06 (* 0.53)	. 8.13 (* 0.60)	7.90 (* 0.42
с _{уна} ст(∞),	▲ (3)	1886.a	26.0	6.51 (* 0.35)	8.17 (* 0.61)	8.31 (* 0.73)	~ 7.66 (+ 0.56
c,Haro(00)3	▲ (3)	1880.2	27.0	7.81 (* 0.13)	8.37 [/] (+ 0.14)	. 8,48 (- 0,13)	8.22 (* 0.13
C,,,,,,,,,,(00),	A(2)	1873.	25.8	7.37 (* 0.04)	7.33 (* 0.27)	7.32 (* 0.29)	7.7 (* 0.20

","" As defined in Table 5-1.

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TABLE 5	-9
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Infrared Frequen	cies, Half-Intensity	Band Widths, (Bicyclohe	and Absolute ptadiene)N(C	Integrated Absorptio	n Intensities for the xane Solution	0 Stretching Vibr	ations of the
		V max	 ∆v‡*		E,** 10 ⁴	H ⁻¹ ca ⁻²	<u> </u>
Complex	Normal	-1 	-1 ca ⁻¹	Ransay	Wilson- Wells	Russell- Thompson	Kan
<u>nor</u> -C,H _B Cr(O),	A, ⁽²⁾	2033.6	1.62	3.14 (+ 0.11)	3.54 (* 0.42)	3.58 (= 0.54)	3.42 (+ 0.36)
nor-CyHgMo(CO)	A1 ⁽²⁾	2044.4	1.50	2.62 (* 0.06)	2.57 (* 0.25)	2.55 (* 0.25)	2.58 (* 0.19)
nor-C7H8V(00)4	A ₁ ⁽²⁾	2043.9	1.95	2.81 (* 0.06)	2.79 (- 0.08)	2.79 (* 0.09)	2.80 (* 0.03)
<u>م.</u> -۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲.,۲	A1 ⁽¹⁾	1958.,	2.86	1.92 (* 0.11)	2.26 (+ 0.52)	2.26 (* 0.58)	2.15 (* 0.40)
<i>,</i> • • •	В ₁	1944.4	3.38	9.29 (+ 0.26)	9.57 (* 1.24)	9.49 (* 1.42)	9.45 (* 0.97
nor-Cyteno(00)	$[A_1^{(1)}, B_1]$	1959.2	3.74	12.2 (+ 0.4)	13.8 (+ 0.4)	14.0 (+ 0.4)	13.3 (* 0.4)
nor-CyHgV(00)	[A ₁ ⁽¹⁾ , B ₁]	1957.3	3.50	12.2 (+ 0.4)	14.4 (+ 1.0)	14.8 (* 1.1)	13.8 (* 0.8
<u>807</u> -C,H8Cr(00)4	B ₂	1914. ₈	3.74	7.76 (* 0.13)	7.99 (+ 0.28)	8.02 (+ 0.34)	7.92 (* 0.25
nor-Cyligho(00)	^B 2	1913.6	3.86	9.19 (* 0.35)	9.90 (* 0.52)	9.97 (* 0.64)	9.69 (* 0.50)
nor-C,HgV(O)	B ₂	1910.	3.40	8.13 (+ 0.17)	8,37 (* 0.43)	8.41 (* 0.47)	8.30 (* 0.36

**** As defined in Table 5-1.

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		~		Specific Intensity	r, 10 ⁴ H ⁻¹ cm ⁻²	ST.
Complex		Solvent	Rapsay	Vilson- Vells	Ressell- Thompson	7.00.3
دمه ⁸ وتد(۳)	2	Sethylene chloride	6.43 (* 0.25)	7.55 (+ 0.53)	7.69 (* 0.68)	" 7.22 (⁺ 0.49)
د م الع نده(۵۵)ع	-	Hethylene chloride	7.66 (+ 0.23)	8.39 (-0.60)	8.45 (- 0.72)	8.17 (+ 0.52)
دؠؾ 8ؠ(∞) ³	U	Kethylene chloride	7.20 (* 0.13)	7.76 (* 0.42)	7.85 (* 0.43)	7.60 (* 0.34)
<u>مەت</u> تى _{اۋ} دە(تى)		,-liexane	5.53 (+ 0.15)	5.84 (* 0.62)	5.54 (- 0.72)	5.74 (* 0.50)
nor-C,HgKo(CO)		n-Hexane	6.01 (* 0.20)	6.56 (+ 0.31)	6.62 (⁺ 8.72)	6.40 (= 0.23)
hor-C,Hav(00)		n-Hexane	5.78 (* 0.16)	6.40 (* 0.35)	6.49 (* 0.41)	6.22 (⁺ 0.31)

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As defined in Table 5-2.

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TABLE 5-10

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parent hexacarbonyl, $M(\Omega)_{6}$ (Tables 5-24 and 5-5), it is evident that a marked enhancement of CO-intensity occurs upon replacement of some of the carbonyl groups with either cycloheptatriene or bicyclo[2.2.1]-hepta-2,5-diene. However, the intensity measurements of the cycloheptatriene complexes were taken in a rather polar solvent, methylene chloride, while those for the hexacarbonyl compounds were determined in non-polar <u>n</u>-hexane. Strictly, therefore, one should not compare the two results, as part of the increase in the CO-intensity value for the cycloheptatriene species could be explained on the basis of solvent effects along [i.e., the marked band broadening in methylene chloride (Tables 5-8 and 5-9) indicates that considerable solvent effects were present]. In the case of the bicyclo[2.2.1]hepta-2,5-diene derivatives, a comparison with the parent hexacarbonyls is valid because both sets of intensity measurements were taken in <u>n</u>-hexane.

The increase in specific CO-intensity for the bicyclo[2.2.1]hepta-2.5-diene derivatives may be interpreted as evidence that there is an increase in M-CO T-electron transfer from metal to CO in the <u>nor-</u> $C_7H_8M(CO)_4$ species over that in the $M(CO)_6$ molecule. The enhanced M-CO The bonding upon substitution may arise from the decrease in M-L T bonding or the increase in M-L σ bonding that occurs when a ligand with lesser T-acceptor or greater T-donor ability replaces CO as donor group. M_{12} Undoubtedly both these factors are important. A similar enhancement of intensity has been reported for substituted octahedral carbonyls. 60,25,65

A comparison of CO-frequency, band-width, and intensity data for all of the group VI metal-olefin carbonyls studied shows that there is little change in the corresponding values as the metal is varied both

when cycloheptatriene and when bicyclo[2.2.1]hepta-2,5-diene act as ligand. No correlation between (0-frequency or half-intensity band width and size of the metal atom is apparent; however, there seems to be a slight tendence for the CO-intensity to increase in the order Cr < W < Mo. Although the variation is small and considerable overlapping of values is present, the trend was considered to be significant because it is observed in both studies. Thus, it would appear that the extent of metal-CO m bonding is greater in molybdenum than in tungsten in the olefin complexes studied here, and hence the results are not in accord with the general trend (i.e., that for any particular type of complex, as one increases the atomic number of the transition metal atom the CO-intensities increase). So, it appears that these two series of group VF metal-olefin carbonyls are also exceptional in their behaviour in this regard.

Nolybdenum- and Manganese-Fhosphine Carbonyls

Another observation which has been consistently reported in the metal carbonyl ∞ -intensity literature is that if there is a halogen atom bonded directly to a transition metal atom, a relatively low specific ∞ -intensity compared to the parent carbonyl will result (Tables 5-11 to 5-13). The trend which has been generally observed is that the specific ∞ -intensity of a metal-fieldgen carbonyl complex decreases in the order I > Br > Cl; although some conflicting results have been reported, they are the exception rather than the rule. The established trend is in good agreement with π -bonding arguments since one would expect that the highly electronegative halogen atoms would withdraw charge from the central metal atom leaving less metal d π -electron density available for backbonding with the ∞ groups.

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	Rectfle Intensities for the CD i of Group VI Yetal-Ha	Birolahing Vibrations of Sono Soria lages Carbony:1 Complemes [®]	10
Complem	Bolvent	Specific Intensity ^b 10 ⁴ H ⁻¹ cm ⁻²	Beferene
œ(œ) _é	Totrahydrofugan	6. io	- • *
0r(00),01 [*]	Chlerefern	4.47	4
Qr(Q),Br	Chlerofern	4.52	` 4
œ(œ),ï¯	- Chloroforn	6.20	▲ <i>1</i>
v(œ)6	Tetrohydrofuren	7.75	٠
₩(00) ₅ 01	Chloroform	° . 6.24	4
¥(@) ₅ 82	- Chleroforn	6.08	4
¥(@),1"	Chlorofern	6.20	د ا
•			

TABLE 5-11

^aS. F. A. Kettle and I. Paul in "Advances in Organometallic Chemistry," Vol. X. F. G. A. Stone and R. Vest, Ed., Academic Press, Inc., New York, N. Y., 1972, p 225. ^bAs defined in Table 5-2.

W. Book and R. E. Mitsmann, <u>L. Maturformoh</u>., <u>178</u>, 577 (1962).

⁴E. V. Abol and I. S. Butler, <u>Trans. Farmday Sec. 53</u>, 45 (1967).

		Specific Intensity ^b	
Complex	Bolvent	10" X" cm"	Meterence
¥n(∞)6 ⁺	Tetrahydrofuran	2.52	٠
Hn(00)501	Chloroform	3.06	۰ د
Mn(00)_Br	Chloroform	3.76	. 🛋
Hn(00)51	Chloroform	4.23	4 *
<u>cis</u> -Mn(00) ₄ ClBr ⁼	Chloroform	2.77	
oie-Hn(@)4Br2	Chloroform	3.47	, (
<u>oir</u> -Mn(CO) ₄ BrI	Chloroform	3,48	4
<u>cla</u> -Kn(@) ₄ I2	Chloroform	3.99	4
$\operatorname{Hn}_2(\infty)_{10}$	Chloroform	3.53	•
Mn ₂ (00)8012	Chloroform	2.51	, 4
Mn ₂ (@) ₈ 8r ₂	Chloroform	2.70	4
Mn ₂ (00)812	Chloroform	3.54	4,
Re(00)6 ⁺	Chloroform	2,11	•
Re(@) ₅ Br	Chloroform	4.25	đ
Re(00) 5	Chloroform	4.19	4

TABLE 5-12

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[•]S. F. A. Kettle and I. Faul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. West, Ed., Academic Press, Inc., New York, N. Y., 1972, pp 224, 225, 226, 230.

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^bAs defined in Table 5-2.

/ ⁰W. Beck and R. E. Mitsmann, <u>B. Naturforach., 178</u>, 577 (1962).

dr. W. Abel and I. S. Butler, <u>Trans. Farmday Soc</u>., <u>63</u>, 45 (1967).

⁹R. H. Wing and D. C. Grocker, <u>Inorg. Chem., 6</u>, 289 (1967).

Complex	Solvent	Specific Intensity ^b 10 ⁴ H ⁻¹ cm ⁻²	Reference
Fe(ထ) ₅	<u>n</u> -Hexane	3.46	¢
<u>019</u> -Fo(00)301.4H2C=CH-CH2	<u>n-</u> Hexane	2.83	0
<u>cie</u> -Fe(O) ₃ Br·H ₂ C=CH-CH ₂	n-Hexane	2.83	٩
<u>cie-</u> Fo(@)301.H2C-CH-CHCH3	<u>n</u> -Xexane	2.90	e
ole-Fe(CO) Brok 2C-CH-CHCH	<u>n-Hexane</u>	2,80	
<u>ets-Fe(</u> Ø) ₄ Cl ₂	Tetrachloroethylene	1,66	، ۲
<u>cis</u> -Fe(@) _{\$} Br ₂	<u>n-Hexane</u>	1.84	,0
<u>cis</u> -Fe(C) ₄ I2	<u>n-Hexane</u>	2.30	٠
π-C ₅ H ₅ Fo(∞) ₂ Cl	n-Hexane	2.7	, d
n-C ₅ H ₅ Fe(CO) ₂ Br	<u>n-Hexane</u>	2.9	, d
π-C ₅ H ₅ F•(∞) ₂ I	n-Hexane .	3.1	đ
π-c ₅ H ₅ F●(∞) ₂ α	Chloroform	3.69	٠
T-C,H, Ye(CO),I	Chloroform	3.64	•

TABLE 5-13

. "S. F. A. Kettle and I. Paul in "Advances in Organometallic" Chemistry," Vol. X.

F. G. A. Stone and R. West, Ed., Academic Press, Inc., New York, N. Y., 1972, pp 226-228. *
^bAs defined in Table 5-2.

^oK. Noack, <u>Helv, Chim. Acta</u>, <u>45</u>, 1847 (1962)

^dJ. Dalton, I. Paul and F. G. A. Stone, <u>J. Chem. Soc</u>., <u>A</u>, 2744 (1969).

[•]D. J. Darenabourg, <u>Inorg. Chen.</u>, <u>11</u>, 1606 (1972).



Specific Int	tensities for (the CO Stretching	Vibrations of	Some Berles of
Metal-Carl	bonyl Complexes	with Phosphorus	Silicon, and	Tin Liganda

Complex	A Bolvent	Specific Intensity ^a 10 ⁴ H ⁻¹ cm ⁻²	Reference
mo(co) ₆	-Hexane	5.33	· b
Ko(00)5PC13	n-l'exane	5.12	•
Mo(00)5P(C6H5)3	<u>n-Hexan</u> e	6:14	•
. Ko(có)6	Cyclohexane	4.2 👳	C
fac-ho(00);(PCl.;);	Cyclohexane	4.2	,
co(α) ₄ sici ₃	<u>n</u> -Hexane	3.15	, 4
100 (00)481 (C6Hg)3	<u>n</u> -liezane	3.80	. 🔺
/π-C5H5 to (@)28nCl 3 ;	Chloroform .	3.1	• •
n-c5H5Fe(0)2Sn(c6H5)3	Chloroform	3.1	, •

As defined in Table 5-2.

^bT. L. Frown and D. J. Darensbourg. <u>Inorg. Chem.</u>, <u>6</u>, 971 (1967); D. J. Darensbourg and T. L. Brown, <u>Inorg. Chem.</u>, <u>7</u>, 959 (1968).

⁰D. Benlian and H. Bigorgne, <u>Bull. Foc. Chim. France</u>, 4106 (1967).

^dD. J. Darensbourg, <u>Inorg. Chim. Acta</u>, <u>4</u>, 597 (1970).

⁹J. Dalton, I. Faul, and F. G. A. Stone, J. Chen. Sog., A. 2744 (1969).

Noack has shown that the halogen effect was additive for mome of the iron halogen derivatives (Table 5-13) that he studied. It would be interesting to see if his observation could be extended to other halogencontaining metal carbonyl complexes where the halogen is not bonded directly to the central metal atom. In metal-halophosphine carbonyls, for instance, one would expect that the halogen effect, although somewhat diminished in magnitude, would be transmitted through the phosphorus atom to the central metal atom and so on to the CO groups.

With this idea in mind, a preliminary investigation of a series of $\underline{fac}-Mo(co)_3[P(C_6H_5)_{3-x}Cl_x]_3$ (x = 1-3) was undertaken. The confrequencies and half-intensity band widths were examined to see if any definitive trends could be observed (Table 5-15). Both band characteristics seemed to be affected in the expected manner with successive substitution of chlorine atoms onto the phosphine ligand. Moreover, a comparison of the Lv_1^t values for the $\underline{fac}-Mo(co)_3(PCl_3)_3$ complex measured in a rather polar solvent, methylene chloride, and in non-polar <u>n</u>-hexane indicated a rather strong solvent offect on band characteristics as well. The results for the solvent and substitution effects act in the same way.

It is difficult to define the properties which mainly affect the $\mathcal{L}v_{\frac{1}{2}}^{t}$ values. It may be supposed that the extent of polarization of the carbonyl groups and their ability to interact with their environment depend upon a combined action of the O-donor and m-acceptor properties of the ligands.⁶⁷ However, steric effects should not be overlooked; the largest $\mathcal{L}v_{\frac{1}{2}}^{t}$ values were obtained in the more strongly reacting solvent as well as with the bulkier substituent ligands. Likewise, the CO

frequencies reflect the combined action of 0-donor and m-acceptor properties of the substituents; that is, the steady decrease in ∞ frequencies along the series <u>fac-Mo(∞)₃(PCl₃)₃; <u>fac-Mo(∞)₃[P(C₆H₅)Cl₂]₃</u>, and <u>fac-Mo(∞)₃[P(C₆H₅)₂Cl]₃ implies that there is either a net decrease in M- ∞ 0-bond strength, or an increase in M- ∞ m bonding, or both.</u></u>

In order to gain more insight into the o and π contributions to the bonding in the various phosphine ligands, we decided to measure the ∞ intensities of the compounds. Since the preliminary investigation of the <u>fac-No(∞)₃[P(C_{6}h_5)_{3-x}Cl_x]₃ (x = 1-3) complexes revealed such definite trends in ∞ frequencies and half-intensity band widths with successive substitution of chlorine onto the phosphine, and since specific ∞ intensities for a similar series of complexes demonstrated the expected dependence on chlorine (Table 5-14), a study of the CO-intensities for this cories of compounds was undertaken to see if a similar trend would occur.</u>

The fac-ho(∞)₃[P(C₆H₅)_{3-x}Cl_x]₃ (x = 1-3) molecules are octahedral with the three carbonyl groups arranged <u>cis</u> to one another (Figure 5-7); the formal symmetry is C_{3v}. The three carbonyls are expected to give rise to two infrared-active ∞ stretching modes: a symmetric A₁ stretching mode and a degenerate E mode (Figure 5-8). All three compounds display two strong bands in the ∞ stretching region as predicted.

The CD-inter**pre** data for the molybdenum-phosphine carbonyls are presented in Tables 5-15 to 5-17. The results are somewhat surprising. Firstly, although there is a marked solvent effect upon the $\Delta v_{\frac{1}{2}}^{t}$ values, no solvent effect upon the integrated intensities is apparent in this study. Intensity results for <u>fac-Mo(CO)₃(PCl₃)₃</u> determined in polar methylene chloride and in non-polar <u>n</u>-hexane show no significant change;



		5	V	Lv1 .		e, ** 10 ⁴	"." c=-2	
Cosplex	Solvest	Mormal mode	a-1	డ ్న	Razsay	Vilson- Vells	Eussell- Thompson	Kean
<u>۲۹۰-</u> ۳۰(۵۵)ر ۲۵۱-۲۵	n-Hexane	Å _{1.}	2040.2	1.43	6.64 (* 0.22)	7.17 (- 0.34)	7.29 (* 0.35)	7.03 (+ 0.30)
f_{BC} -Ho($(\infty)_{3}(P_{2})_{3}$	Nethylene chloride	٨,	2039.0	10.5	6.10 (- 0.13)	6.07 (* 0.40)	6.64-(- 0.43)	6.07 (+ 0.32)
<u>fac-%(0)</u> 3[P(c6H5)C12]3	Methylene _chloride	Α,	2006.	12.9	5.77 (* 0.09)	5.90 (- 0.32)	5.90 (- 0.36)	5.96 (- 0,26)
<u>fec</u> -m(0) ₃ [P(c ₆ 1 ₅) ₂ 0] ₃	Methylens chloride	, ^A 1	1970.5	13.6	6.20 (* 0.17)	6.17 (⁺ à.00)	5.92 (* 0.79)	6.10 (* 0.65)
<u>fac-Ho(</u> (0) ₃ (PCI ₃) ₃	n-Hexane	- E	1991.1	4,46	10.1 (* 0.3)	10.7 (* 0.5)	10.7 (* 0.6)	10.5 (* 0.5)
<u>fac-Ko(00)3(PC13)3</u>	 Methylens chloride 	E	1996.	26.8	11.3 (+ 0.3)	11.2 (+ 1.1)	11.1 (* 1.1)	11.2 (* 0.8)
<u>fac-xo(0)3[b(c^{H2})01]</u>	Kethylehe chloride	I	1940	33.2	11.0 (= 0.1)	10.8 (+ 0.5)	10.8 (+ 0.6)	10.9 (* 0.4)
<u>fac-</u> но(Ф) ₃ [Р(С6 ^Н 5)2 ^{СІ}]	Kethylene chloride	E	1889. ₄	37.0	12.1 (- 0.5)	11.4 (* 3.0)	10.8 (+ 2.7)	11.4 (- 2.1)

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**** As defined in Table 5-1.

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TABLE 5-15

	LITCOULTY MELLO OF LOS SYMPLETC	$(\alpha)_{3}(rosphise)_{3}$ Ouplems		•
	•	E. 10 ^b f ⁻¹	ہے۔ 1 مع - 2	n)
Camplex	Solvert	4	- 100	·Ing
ر (م)ر (م)ر (م)ر (م)ر (م)ر (م)ر (م)ر (م)	- interest	7.03 (2 0.30)	20.5 (2 0.5)	0.670
	Methylene chloride	6.07 (± 0.32)	11.2 (2 0.5)	6. 9
، کُمْ ت (دی ۲۹ د) ۲۰ (۵۰ میلون) میکور (۵۰ میلون) در مارد میلوند.	Kethylene chloride	5.86 (2 0.26)	· 10.9 (2 0.4)	0°230
ر تم ₂ (ی ⁴ ی)۴)د(۵)م اد م ف	Methylene chloride	6.10 (2 0.65)	11.4 (2 2.1)	6.5 35

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TIELZ 5-16

• • •		-	Specific Intensity	, 10 ⁴ z ⁻¹ cz ⁻²	_
Complex	Solvent	Bassay	Vilson- Vells	Russell- Thompson	Nees
<u>fac-tio(</u> () ₃ (PCl ₃) ₃ -	<u>p-Sezane</u>	5.57 (= 0.17)	5.96 (- 0.29)	6.61 (* 0.32),,	5.85 (* 0.26)
(ac-Ho(CC)) (PC))	Kethylete chloride	5.81 (= 0.14)	5.74 (= 0.49)	5.71 (* 0.51)	5.75 (* 0.39)
<u>fec-Bo(CO) (C6H5)C12</u>]	Hethylete chloride	5.53 (+ 0.06)	5.56 (+ 0.28)	5.55 (= 0.31)	5.57 (= 0.22)
<u>tac-m(00) (c6=5)2</u> ,	Nethylene chloride	6.11 (* 0.22)	5.85 (* 1.34)	5.55 (= 2.24)	5.8 (* 0.91)
<u>ets-Ra(00) [P(C_K,)] br ~</u>	1-Tezane	3,50 (* 0,12)	. 3.70 (* 0.25)	3-73 (* 0-26)). # (- e. 22)
eta-Ra(0),[As(CgHg)])T	t-lemme	3.62 (* 0.11)	3.60 (- 0.30)	3.59 (* 0.32)	3.60 (* 0.24)
ete-Ra(00) [50(C,K,),) 37	<u>n-Sexane</u>	3.70 (+ 0.12)	3.93 (* 0.18)	3.96 (+ 0.20)	3.66 (* 0.27)

TABLE 5-17

"As defined in Table 5-2.

apparently, the increase in band width with the polar solvent is roughly compensated by the decrease in peak height. It should be noted, however, that Penlian and Bigorgne's⁶³ value (determined by the Wilson-Wells method) for fac-Mo(CO)₃(POl₃)₃ in cyclohexane solution, 4.2×10^4 M⁻¹ cm⁻², is considerably lower than our results (also obtained by the Wilson-Wells extrapolation method) for both <u>n</u>-hexane and methylene chloride solution, 5.96 and 5.74 $\times 10^4$ M⁻¹ cm⁻², respectively. This may reflect either a difference in measuring technique or a solvent effect or both.

Second, if our result (obtained by Ramsay's method) for fac-Mo(∞)₃-(PCl₃)₃ in <u>n</u>-hexame solution is compared with the published value (also determined by Ramsay's method) for Mo(∞)₆ (Table 5-5) measured in the same solvent, it is evident that there is no significant change in the specific CO-intensity an one goes from the parent hexacarbonyl, 5.39 x 10⁴ M⁻¹ cm⁻², ⁶⁰ to the phosphine complex, 5.57 x 10⁴ M⁻¹ cm⁻² [a point observed by Penlian and Bigorgne⁶³ as well in their measurements on Mo(∞)₆ and Mo(\cos)₃(PCl₃)₃ (Table 5-14)]. This may be interpreted as evidence that there is little change in the M- ∞ N-electron transfer from metal to ∞ in the phosphine derivative compared to the Mo(∞)₆ molecule. Evidently any increase in the M-L o bonding due to withdrawal of charge by the highly electronegative chlorine atoms is roughly compensated by a decrease in N-L N bonding so that the expected decrease in the M- ∞ N bonding upon substitution does not occur.

Moreover, unlike the $\Delta v_{\frac{1}{2}}^{t}$ values and the CO frequencies, no effect of ligand upon the integrated intensities is evident, indicating that there is essentially no change in the extent of M-CO T bonding in the three substituents, PCl₃, P(C₆H₅)Cl₂ and P(C₆H₅)₂Cl. This observation strongly suggests that the decrease in CO frequency as one goes from PCl₃ to

- 73

 $P(C_6H_5)_2$ Cl stems mainly from a weakening of the CO 0 bond, and not from an increase in π bending. The fact that there is no appreciable change in the intensity ratios either with the three ligands (Table 5-16) further supports this conclusion. Although this is not the expected result, similar conclusions concerning M-P bonding characteristics have been reported.^{65,75-78} It should be emphasized, however, that when bulky ligand groups are present, intensity results should not be interpreted simply on the basis of the π -acceptor and o-donor ligand abilities alone. Geometrical factors, particularly as far as steric distortion by the substituent ligand is converned, should also be considered.

Another interesting aspert of CO-intensity studies on metal-phosphine carbonyls is with phosphine-like ligands such as arsines and stibines. Several studies have been carried out on series of metal carbonyl complexes with phosphorus, arsenic, and antimony (or with silicon, germanium and tin) donor ligands (Tables 5-19 and 5-20); however, no general trend is apparent.

From a comparison of the (∞) frequencies and half-intensity band widths for the series of <u>cis-Mn((∞))</u> [LBr complexes containing triphenylphosphine, -arsine, and -stitute ligands (Table 5-18), it is clear that there is no obvious trend. It appears, therefore, that there is either no difference in the o- and m-bonding abilities in the three ligands or that any differences in o and m bonding complement one another (i.e., a decrease in 0 bonding in any one ligand may be compensated by an increase in m bonding and <u>vice versa</u>). In the hope of gaining some more information concerning the bonding characteristics of the three ligands, we decided to measure the (∞ intensities of the tetracarbonyl complexes.

•

	-	Y HAX			E, ** 10 ^k	E ⁻¹ ca ⁻²	
Complex	[¶] ormal Röde	_ -1 [°]	ca-I	Bansay /	Wilson- Wells	Russell- Thompson'	Meen.
<u>ets-Ma(@)</u> [P(Cgfis)]Br	A.(2)	2089.2	2.66	1.61 (* 0.03)	1.62 (- 0.11)	1.63 (= 0.12)	1.62 (- 0.09)
$\underline{c1s} - R_{\pi}(\omega)_{a} [As(C_{6}H_{5})_{3}]Br$	A.(2)	2091.8	2.60	1.58 (-0.02)	1.51 (- 0.06)	1.51 (- 0.06)	1.53 (= 0.05)
cis-Ma(O) [Sb(C635)]br	A.(2)	2087.5	3.06`	1.87 (* 0.06)	2.00 (- 0.11)	2.02 (- 0.12)	1.96 (* 0.10)
ete-m(@),[?(CgEg),]er	A.(1a)	2024.6	5.83	3.25 (* 0.11)	3.42 (* 0.26)	3.45 (⁺ 0.29)	3.35 ([±] 0.22)
ta-Ma(00) [As(C6H5)] Pr	A*(1=)	2023.7	6.00	3.91 (+ 0.17)	4.20 (- 0.42)	4.19 (- 0.48)	A.10 (- a.36)
<u>ers-m(@) [sb(c675)]Br</u>	At(In)	2021.6	5.48	3.34 (* 0.11)	3.62 (* 0.09)	3.65 (* 0.1 1)	3.54 (* 0.11)
) <u>cia-En(CC), [P(CcH,),]</u> Br	٨-	2006.	6.53	5.33 (+ 0.26)	5-75 (-* 0-52)	5.82 (- 0.59)	5.63 (- 0.46
<u>ets-Rm(O)</u> [As(C ₂ E ₂)] Br	A	2011.	6.48	5-21 (-0-13)	5.15 (- 0.35)	5.13 (= 0.37)	5.16 (= 0.28)
<u>ers-200</u> (0) (20(062) 3)	A -	2009. ₂	6.91	5.57 (0.22)	6.01 (- 0.27)	6.09 (- 0.30)	5.89 (- 6.26)
cia-Ma(O) [F(C, R,), br	<u>"</u> .(ъ)	1960.8	11.0	3_50 (+ 0_09)	4.01 (⁺ 0.11)	4.04 (⁺ 0.12)	3.95 (⁺ 6.11)
cta-Ha(00) [40(C6H4)] br	▲-(р)	1962.7	10.5 ,	3.77 (- 0.12)	3-54 (* 0.36)	3.52 (* 0.36)).£ (- 6.25
<u>da-m(a)</u> [So(cgs,),]br	*.(p)	1960.	10.5	4.01 (- 0.09)	4.07 (- 9.26)	4.07 (⁺ 0.28)	A.65 (* 0.21

TABLE 5-18

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"" As defined in Table 5-1.

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TAI	12	5-19	

Complex	Selvent	Specific Intensity ^b 10 ⁴ H ⁻¹ cm ⁻²	Reference
νω(ω) ₆	D-Hezene	5.33	•
No(00)5P(C6H5)3	A-Hezane	6, 14	•
Ho(00)5AB(C6H5)3	D-Hemane	5.67	· •
мо(co)526(c6H2)3	<u>n</u> -Hexane	5.76	, • •
Mo(@)6	Cyclohezane	~ 4.20	¢.
<u>918</u> -Ko(00)4[F(C2H5)3]2	Cyclohezane	4.50	4
<u>ein-Ko(</u> (),[Sb(C2H5))]	Cyclohezane	4.80	•
œ(œ) ₃ %	Cyclohezane	3.2	•
∞(∞) ₂ .∞⊾(c ⁶ n ²) ³	Cyclohezane	4.0 ·	• v ●
00(00)2NOA=(C6H5)3	Cyclohezane	3.9	•
ω(ω) ₂ 1056(c6H3)3	Cyclohezans	· 3·3	` ₽
n-c5H5Hn(@)	Carbon disulphide	4.2	، ۲
π-C5^H5^{Hn}(∞)2^{P(C6H}5)	Carbon disulphide	a . 4.7	• • f
n-c5H5Hn(@)2As(C6H5)3 ,	Carbon disulphide	4.8	ſ
π-C5H5Kn(^(C) 2 ^{8b(C6H5)})	- Carbon disulphide	4.7	f

^aS. F. A. Kettle and I. Faul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. West, Ed., Academic Press, Inc., New York, N. Y., 1972, pp 225, 229, ^bAs defined in Table 5-2.

⁴T. L. Brown and D. J. Darenabourg, <u>Inorg. Chem.</u>, <u>6</u>, 971 (1967); D. J. Darenabourg and T. L. Brown, <u>Inorg. Chem.</u>, <u>7</u>, 959 (1968).

^dD. Benlian and H. Bigorgne, <u>Bull, Soc. Chim. France</u>, 4100 (1967).

A. Poletti, A. Foffani, R. Cataliotti, Spectrochim. Acta, 264, 1063 (1970).

^TW. P. Anderson, T. B. Brill, A. R. Schoenberg, and C. V. Stranger, Jr., <u>J. Organometal</u>. <u>Chem., AA, 161 (1972).</u>

Complex	Bolvent	Specific Intensity ⁴ 10 ⁴ H ⁻¹ cm ⁻²	Reference
ν	D-Hezane	3.46	6
ro(00)4r(C6H5)3	n-Hexane	4,00	•
n-c5H5 *** (00)251C13	Chloroform	· 3.2	• 🖕 •
m-C5H5Fe(00)2CeCl3	Chlorofora	3.0	۰ 🗳
n-C5H5Fe(00)23nC1)	Chloroform	3.1	4

^aAs defined in Table 5-2.

bx. Noach, Helv. Chin. Acta. 45, 1847 (1962).

⁰D. J. Darensbourg. Inorg. Chin. Acta. 4, 592 (1970). #

dJ. Dalton, I. Paul, and F. G. A. Stone, <u>J. Chem. Bog</u>., <u>A</u>, 2744, (1969).

TABLE 5-20

The <u>cia-Mn((∞))</u> LBr $[L = P(C_6H_5)_3, As(C_6H_5)_3, Sb(C_6H_5)_3]$ molecules have an octahedral structure with a <u>cis</u> arrangement of the L and Br groups (Figure 5-9). This configuration results in formal C_5 symmetry for which group theory predicts four infrared-active CO stretching frequencies $(3A^{+} + A^{+})$. The four modes are very similar to those in the <u>nor-C_7H_8M((∞))</u> species. The <u>trans</u> pair of carbonyls should give rise to a weak symmetric $A^{+(2)}$ mode (analagous to the $A_1^{-(2)}$ mode of the <u>nor-</u> $C_7H_8M(CO)_4$ molecule) and to a strong antisymmetric A^{+} stretching mode (analagous to the <u>nor-C_7H_8M((∞))</u> B_1 mode). The carbonyl <u>trans</u> to the L group should give a symmetric $A^{+(1a)}$ mode of medium intensity and the carbonyl <u>trans</u> to the Br group should produce another symmetric A^{+} (1b) mode of comparable intensity (Figure 5-10). The infrared spectra show four well-ronolved bands as expected.

Pousible assignments for the infrared spectra of the <u>cis-Mn(∞)</u>₄LBr species were made on the basis of the similarity of the vibrational modes to those of <u>nor-C₂H₈Cr(∞)₄. Thus, the highest frequency absorption was assigned to the A¹⁽²⁾ mode and the most intense peak was assigned to the A" mode. The greater inductive character and lower n-bonding ability of the bromide compared with the phosphine, arsine and stible led to the assignment of the lowest frequency peak to the A^{1(1b)} stretching mode of the carbonyl <u>trans</u> to the bromide, and leaving the A^{1(1a)} mode to the romaining unassigned peak. These assignments have now been verified by an infrared spectrascopic study-of the exchange reactions of <u>cis-Mn(∞)</u>₄LBr with ¹³CO and subsequent approximate force constant calculations.⁷⁹</u>

The CO intensities for the <u>cis-Mn(CO)</u>₄LBr complexes are presented in Tables 5-17 and 5-18. A comparison of the results shows that, as for the



Figure 5-10 Infrared Active CO Stretching Vibrations of the $\underline{\text{cis}}-\text{Mn}(\text{CO})_{\mu}\text{LBr}$ Molecule $[L = P(C_6H_5)_3, As(C_6H_5)_3, Sb(C_6H_5)_3].$

.

(0) frequencies and half-band widths, the intensity values with $L = P(C_6H_5)_3$. As $(C_6H_5)_3$ and Sb $(C_6H_5)_3$ are very close. So, it may be taken that the extent of metal-ligand π bonding in phosphorus, arsenic and antimony is essentially the same. In conclusion, therefore, it appears that both the o-donor and the π -acceptor properties of phosphorus, arsenic and antimony towards manganese are equivalent.

Other workers have reached similar conclusions concerning the M-P, M-As and M-Sb bond strengths in molybdenum and nickel complexes. Specifically, (1) 60-frequency studies⁴⁴ on tricarbonylmolybdenum compounds containing triphanylphosphine, -arsine and -stibine ligands and (2) dipole-moment measurements⁸⁰ on complexes containing Ni-P and Ni-As bonds have indicated that the differences in multiple bonding capacities of phosphorus, arsenic and antimony towards a transition metal atom are virtually negligible. (However, as mentioned earlier in the discussion of the molybdenum-phosphine carbonyls (vide supra), steric effects may also be important in the interpretation of infrared results when bulky ligand groups are present. Thus, geometrical factors as well as N-acceptor and o-donor ligand abilities should be considered when interpreting the results.

T-Cyclopentadienylmanganese(I) Complexes

Clarification of the ability of metal-arene carbonyls to transfer an electronic effect through the transition metal atom is a subject of considerable interest. With this aim in mind, several workers have attempted to correlate the CO intensities of metal-arene carbonyl complexes with the electron acceptor ability of the ring substituents. Nesmeyahov et al. 5^8 have demonstrated that for a series of

Complex	Solvent	Specific Intensity ^b 10 ⁴ M ⁻¹ cm ⁻²	Reference
CpMn(CO)	Carbon tetrachloride	1	; G
Cp = n-C ₅ H ₅	, a _	4.85	
$Cp = \pi - C_3 H_{\gamma} C_5 H_{\phi}$	Ū	4,86	•
Cp = T-C6H5CH2C5H4		4.83	
Cp = π-COCH2C5H4	•	4.50	
Cp = n-C6H500C5H4	-	4.59	
$Cp = \pi - CF_3 COC_5 H_4$	a * *	4.14	• • /
$C_{p} = \pi - (C_{2}H_{5})_{2}C_{5}H_{3}$	<u> </u>	5.05	•
Cp = n-C(CH3)3(C6H50)C5H3		4.75	, s ²
$Cp = \pi - (C_2 H_5)_2 (CH_3 CO) C_5 H_2$		4.82	· • •
$C_p = \pi_{-}(C_2H_5)_3(CH_3OC_5H)$	•	4.94	. r
Cp = π-(C2H5)4(CH300)C5	· · ·	4.97	•
$C_{p} = \pi - (C_{2}H_{5})_{4}C_{5}H$	1	5.30	
cpπ•(∞) ₃	Carbon tetrachloride		c,
Cp = π-C ₅ H ₅	-	5.36	· · · ·
$C_p = \pi - CH_3 OC_5 H_4$	•	4.92 k	ب حد ر
Cp = T-C6H500C5H4	•.	, 5.14	•

Specific Intensities for the CO Stretching Vibrations of Semi Series, of Group VII Ketal-Cyclopentadienyl Cambonyl Complexes

TABLE 5

^aS. F. A. Kyttle and I. Paul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. West, Ed., Academic Press, Inc., New York, M. Y., 1972, p.227.

^bAs defined in Table 5-2.

^CA. N. Nesseyanov, G. G. Dvoryantseva, Xh. P. Fustyl'nik, Yu. N. Sheinker, N. E. Kolobova, and K. N. Anlainov, <u>Dokl. Akad. Nauk SSSR</u>, <u>174</u>, **368** (1967); <u>Acad. Sci. USSR</u> <u>Chem.</u>, <u>176</u>, 457 (1967).

Spec	ific.Iptensities for the one Series of Arene-Chro	0 OStretching Vibrations mium Carbonyl Complexes	
Complex	Solwnt	Specific Intensity ^b 10 ⁴ N ⁻¹ cm ⁻²	Reference
RC6H4000CH3Cr(00)3	Carbon tetrachloride		· · · · · · · · · · · · · · · · · · ·
R = =-000CH3	}	6.07	2
R = 2-000CH3	./	4.73 · · · :	1
R = 0-Cl	· · · ·	5.13	
R = <u>n</u> -Cl ,	• .	4.53	
R = p-CL		4.23	
R = H ⁻¹		··· 5.03	•,
R = o-CH3		· 5.07	•
R = a-CH3	` - ,	\$.3 0	.*
R = P-CH3	,	5.73	•
R = 0-OCH3	<u>.</u>	5.60	ŋ
, R = 2-0CH3	•	5.20	
e * 5 -	•	•	\$?
céhêcr(∞)3	Cyclohexane	4.50	· d
$(C_{H_{\varepsilon}})_{\zeta}C_{\zeta}Cr(\infty)$	Cyclohexate	, 5.31	, · · · i

CABLE 5

^BS. F. A. Kettle' and I. Paul in "Advances in Organometallic Chemistry," Vol. X, F. G. A. Stone and R. West, Ed., Academic Press, Inc., Wew York, N. Y., 1972, p 227. ^bAs defined in Table 5-2.

G. Klopman and K. Noack, Inorg. Ches., 7, 579 (1968).

^dR. D. Flacher, <u>Spectrochim. Acta</u>, <u>19</u>, 842 (1963).

The cyclopentadienyltricarbonylmanganese(I) and -rhenium(I) derivatives the specific ∞ intensities decrease with the electron acceptor ability of the ring substituents (Table 5-21). That is, the electronic effects on the cyclopentadienyl ring are transmitted to the central metal atom and then to the ∞ groups. This is in contrast to the observations of Klopman and Moack⁶⁴ on a series of substituted benzoic-methyl-ester tricarbonyl chromium derivatives (Table 5-22). These workers observed only small variations in specific intensity which could not be correlated with the properties of the substituents.

In this work, the CO intensities of two ring-substituted derivatives of π -cyclopentadienyltricarbonylmanganese(I) were measured. At the same time, the CO and CS intensities of the thiocarbonyl complex, π -C₅H₅Mn(CO)₂CS, were studied to see if we could say something about the relative bonding properties of \Im and CS. Thiocarbonyls have been known since 1966, when Baird and Wilkinson⁸¹ synthesized trans-Rh(CS)[P(C₆H₅)₃]₂X(X = Cl, Br). However, as yet, there have been no intensity studies to determine how CS compares to CO as a ligand. Physical and chemical evidence indicate that CS is both a better σ - and π -bonding ligand than CO.^{82,83}

The Π -RC₅H₄Mn(∞)₃ (R = H; CH₃) molecules have a piano-stool-type structure ³² (Figure 5-11). Moreover, since there is free rotation of the cyclopentadienyl ring about the metal-ring axis (irrespective of the presence of substituents on the ring³²), the molecules possess C₃v symmetry. Two infrared-active ∞ stretching modes (A₁ + E) are expected for structures of this type. The <u>cis</u> carbonyls should give rise to a symmetric A₁ stretching mode and a degenerate E mode (Figure 5-12). The spectra show two well-resolved bands in accordance with the group theoretical predictions.





The thiocarbonyl complex, $\Pi-C_5H_5Mn(\Omega)_2CS$, also has a piano-stooltype structure⁸⁴ (Figure 5-13) and the local symmetry of the $Mn(\Omega)_2(CS)$ molety is C_s . Two infrared-active ∞ stretching frequencies (A' + A") and one infrared-active CS stretching frequency (A') should be observed for such an arrangement. The <u>cis</u> carbonyls should give rise to a symmetric A' mode and an antisymmetric A" mode. The thiocarbonyl must of necessity produce a symmetric A' stretching mode (Figure 5-14). The infrared spectrum exhibits the anticipated results: two well-resolved bands in the ∞ stretching region and one isolated peak in the CS stretching region (1350 - 1200 cm⁻¹).

The ∞ -intensity data for π -RC₅H₄Mn(∞)₃ (R = H, CH₃) and π -C₅H₅Mn(∞)₂CS are given in Tables 5-23 and 5-24. The specific intensity of the tricarbonyl molety in π -C₅H₅Mn(∞)₃ obtained by Ramsay's method in <u>n</u>-hexane (4.92 x 10⁴ M⁻¹ cm⁻²) agrees slosely with Nesmeyanov et al.'s⁵⁸ value (4.85 x 10⁴ M⁻¹ cm⁻²) obtained by Ramsay's method for a more polar solvent, carbon tetrachloride; however, it is somewhat higher than the value reported by Anderson et al.⁷² in carbon disulphide solution (4.19 x 10⁴ M⁻¹ cm⁻²). The low value in the Anderson study is not surprising since they evaluated their absolute integrated intensities by taking the mean of several apparent intensity measurements, while the Ramsay results in this study were obtained by taking the mean of several' apparent intensity determinations and then adding 2% (vide supra, page 39). In order to make a meaningful comparison between the two studies, Anderson's value was increased by 2%.

Anderson's adjusted value for $\pi - C_5 H_5 Mn(CO)_3$ (5.01 x $10^4 M^{-1} cm^{-2}$) compares very well with the corresponding result in this study, even though



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<u> </u>	•				
Infrared Frequencies,	Half-Intensity Band	Widths, and Absolute	Integrated Absorption	Intensities for the 00	Stretching Vibrations of
the (T-Cyclopentadien	yl)Xn(CO), Complemes	and for the CO and C	5 Stretching Vibrations	of the (n-Cyclopentad)	enyl)Kn(CO)_CS Complex

	Carbon		<u> </u>	<u></u>	Ransay ,	Yells	Thospson	Nenz
(2003).co		Normal	1		•	Wilson-	Russell-	
			V . BRX	∆v ^{t•}	\sim	e , ** 10 ⁴	H ⁻¹ cm ⁻²	
X ¹	، ۳			cs ra	DE			Υ.
n-C5H5Kn(@)2CS	Carbon disulphide	٨.	1956.9	8.00	6.14 (* 0.13)	6.56 (* 0.23)	6.68 (- 0.29)	6.46 (* 0.22)
(n-CH3C5H4)Ka(CO)3	n-Hexane	E	¹⁹⁴⁴ .5	6.60	11.2 (+0.4)	12.6 (+1.3)	12.8 (+1.5)	12.2 (+1.1)
T-C5H5Kn(Q)3	<u>n-Hexane</u>	2	¹⁹⁴⁸ .6	4.11	11.1 (* 0.2)	12.0 (+ 0.3)	12.1 (* 0.4)	11.6 (* 0.3)
π-C ₅ 3 ₅ 8n(α) ₂ α	Carbon disulphide	۸۰ .	2009.7	5.65	4.62 (* 0.09)	4.92 (- 0.17)	5.01 (+ 0.20)	4.85 (* 0.15)
(T-CH3C5H4)XB(@)3	<u>n-Hezane</u>	*1	2026.7	1.91	3.77 (= 0.05)	3.55 (+ 0.36)	·3.55 (+ 0.37)	3.62 (* 0.26)
n-05H5Xn(@)3	<u>n-Hexane</u>	A ₁ .	2030. 8	1.86	3.65 (* 0.05)	3.77 (+ 0.28)	3.80 (* 0.32)	3.74 (* 0.22)
Complex	, Solvent	Mormal mode	_a ⁻¹	<u>ca</u> -1	Ransay	Vilson- Vells	Bussell- Thompsont	Kenn
			V.BAX	" Cv+	·	E , [•] 10 ⁴	R ⁻¹ cm ⁻²	
0				<i>w rv</i> .		•		1

TA3L3 5-23

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••••	т	A NOTES			
	• • •		Specific Istemuity,	• 10 ⁴ K ⁻¹ cm ⁻²	
Complex	Solvent (, Kenned	Vilaco- Yella	Prasell- Thompsen	Ţ
*-c ₅ it ₅ *=(∞) ₃		*. 92 (± 0.03)	5.25 (* 0.21)	5/71 (± 0.23)	5.16 (2 0.17)
(m-cl ⁵ c ₅ H ₄)m ₄ (m) ₃	i B-Hotane	5.00 (± 0.15)	5.38 (* 0.55)	5.46 (± 0.63)	5.28 (* 0.44)
=c2#2m(@)~G2	Carten disulphide	5.26 (ڭ ٥.07)	5.76 (2 0.20)	5.25 (2 0.25)	5.66 (2 0.17)
		2002			•
Ť.			Specific Intendity,	• 10° K ⁻¹ m ⁻²	-
Complex 7	Solvent		Tilsen Fells	finesell- Thorpson	Land
-c ² t ² z=(ω) ⁵ αχ	Carton di mijbhide	5.42 (± 0.13)	5.88 (± 0.12)	5.99 (± 0.22)	5.76 (± 0.18)
in defined at	be 52. Do	۰. -	4.1	,	** *

FAILS 5-24

the measurement was taken in a more polar solvent (carbon disulphide). The close agreement in intensity values in all three solvents (i.e., <u>n</u>hexane, carbon tetrachloride and carbon disulphide) suggests that there is little if any difference in solvent effects upon the integrated COintensities of Π -C₅H₅Mn(CO)₃ in these solvents. Apparently, any increase in band width with the more polar solvents is roughly compensated by a decrease in peak height. This appears to be the case for other systems studied here as well (<u>vide supra</u>, page 68). In view of these results, it is evident that our CO-intensity values for Π -C₅H₅Mn(CO)₃ in <u>n</u>-hexane solution may be reliably compared with those for Π -C₅H₅Mn(CO)₂CS in carbon disulphide solution.

However, a word of caution regarding comparison of ∞ -intensity measurements in different solvents is in order. A notable exception to the foregoing conclusions regarding solvent effects upon ∞ -intensities is the behaviour of cyclohexane. Fischer's⁵⁶ value for π -C₅H₅Mn(∞)₃ (4.01 x 10⁴ M⁻¹ cm⁻²) obtained by Ramsay's method in cyclohexane solution is considerably lower than any of the results cited above for more polar solvents. Since similar behaviour in cyclohexane solution was observed in other systems as well (<u>vide supra</u>, page 73), the decrease is considered to be significant and indicates that solvent effects upon ∞ -intensities should be regarded when comparing measurements in different solvents, especially if one of the solvents is cyclohexane.

The specific (0)-intensity seems to increase in going from π -C₅H₅-Mn((0))₃ (5.16 x 10⁴ M⁻¹ cm⁻²) to π -CH₃C₅H₄Mn((0))₃ (5.28 x 10⁴ N⁻¹ cm⁻²) in agreement with the conclusions of Nesmeyanov et al.⁵⁸ that electronic effects on the cyclopentadienyl ring are transmitted to the central metal atom and so to the (0) groups (i.e., if transmission of electronic effects

on the ring to the CO groups occurs, then one would expect from T-bonding arguments that the availability of metal dT-electron density would increase with the donor ability of the ring subatituents). Since a methyl group is a better electron donor than a hydrogen atom, one would anticipate a small increase in CO-intensity with the methyl-substituted derivative. However, it should be kept in mind that the errors in the intensity measurements are such that the two values for $TT-C_5H_5Mn(-CO)_3$ and $TT-CH_3C_5H_4-Mn(-CO)_3$ overlap and so no definitive conclusions may be drawn from these results. Unfortunately, Nesmeyanov et al. did not give, any estimate of the errors present in their measurements.

The specific ∞ -intensity of π - $c_5H_5Mn(\infty)_3$ (5.16 x 10⁴ M⁻¹ cm⁻²) is significantly lower than that of π - $c_5H_5Mn(\infty)_2CS$ (5.66 x 10⁴ M⁻¹ cm⁻²). This may reflect a greater o-donor ability of CS over ∞ or a lesser π acceptor ability, or both. However, the specific CS-intensity (5.76 x 10⁴- M^{-1} cm⁻²) of π - $c_5H_5Mn(\infty)_2CS$ is comparable to the specific ∞ -intensity (5.66 x 10⁴ M⁻¹ cm⁻²). This may be taken as direct evidence that the π -acceptor ability of CS is at least as good as that of ∞ . It necessarily follows that the increase in specific ∞ -intensity in going from π - c_5H_5 - $Mn(\infty)_3$ to π - $c_5H_5Mn(\infty)_2CS$ is due solely to the greater σ -donor ability of CS over CO and not to a lesser π -acceptor ability.

Dipole moment derivatives were calculated as previously described by Darensbourg²⁵ using the equation

 $I_{T} = n G_{tt} \mu_{ML}^{2}$ (24)

where I_T is the total measured intensity for the CO or CS stretching vibrations (obtained by the Wilson-Wells extrapolation) and is defined as

Alt of Y .

2.303 E x 10⁻⁴, and n equals the number of CO or CS groups in the molecules. C_{tt} is the term representing the inverse mass of the CO or CS grouping (based on atomic mass units) and is equal to 0.14585 and 0.11444 for the CO and CS groups, respectively. μ'_{ML} is the respective change in group dipole moment with symmetry coordinates (defined as the effective group dipole moment derivative involving both MC and CO or CS stretching motions) and is expressed in arbitrary units.²⁵.

The μ_{MCO} 's, 9.10 and 9.51 for the Mn(∞)₃ (where $I_T = 2.303 [E_{A1} + E_E]$ x 10⁻⁴ = 2.303 [3.77 + 12.0] = 36.28) and the Mn(∞)₂CS (where $I_T = 2.303$ $[E_{A^+} + E_{A^+}] \times 10^{-4} = 2.303 [4.92 + 6.56] = 26.4$) moleties, respectively, agree closely with Darensbourg's values obtained in molybdenum (9.00), cobalt (7.56), iron (7.74) and iridium (7.78) substituted carbonyls. $\frac{65}{5}$.68.69 However, the μ_{MCS} for the Mn(∞)₂CS (where $I_T = 2.303 E_A$, $\times 10^{-4} = 2.303$ [5.88] = 13.5) species, 10.9, is substantially larger than these. This indicates that the U-acceptor ability of the CS group is considerably greater than that of ∞ in agreement with molecular orbital calculations on the CS and ∞ molecules.⁸² and also with many other experimental data.⁸³

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PART IV

SUMMARY AND PROSPECTUS
SUMMARY AND PROSPECTUS

Integrated infrared intensity measurements of the OO stretching modes in transition metal-carbonyl complexes have provided useful information concerning the electronic character of the coordinated OO group.^{25,54-73} The intensities have been shown to be highly dependent on the T-electron density in the M-OO bond.^{25,65} The work in this thesis constitutes an additional study of the infrared intensities of the OO stretching absorptions of some group VI and VII metal-carbonyl complexes. The integrated intensity results were used as a basis for a further investigation of the O- and T-bonding abilities of the various transition metals and ligands in these complexes.

The results of this study also demonstrate the utility of integrated intensity measurements in assessing T-electronic delocalization in transition metal-thiocarbonyl compounds. The first measurement of the infrared intensity of a CS stretching mode in a metal-thiocarbonyl complex was reported here. The intensity measurements and the calculated dipole moment derivatives reflect the demands made for T-electron density as well as the availability of T-electron density during a ligand stretching vibration. From a comparison of the CO and CS results in the T-C₅H₅-. $Mn(CO)_2$ CS molecule, it is clear that CS has a greater T-acceptor ability than CO, in accordance with molecular orbital calculations on the CO and CS molecules⁸² and many other experimental findings.⁸³

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An intriguing aspect of transition metal-thiocarbonyl chemistry, in view of the great similarity in bonding properties of CO and CS, is the surprising difference in infrared frequency shifts (bonded ligand versus free ligand) of M-CO and M-CS bonds." Whereas the CO stretching frequency " of free 00 decreases from 2143 cm⁻¹ to approximately 2000 cm⁻¹ (sometimes to as low as 1850 cm⁻¹ in the case of terminal CO bonding) upon coordination to a transition metal atom, it appears that the CS stretching frequency can both increase and decrease (i.e., the range of v(CS) stretching frequencies in metal-thiocarbonyls is 1381-1193 cm⁻¹) with respect to the frequency of "free" CS (1274 cm^{7} when trapped in a CS, matrix⁸⁵).⁸³ Since n-backbonding is almost certainly present in metal-CS bonds, it is apparent that frequency values alone cannot lead to meaningful conclusions as to the relative m-acceptor ability of the CS and CO groupings. ۸n interesting project for future consideration would be to attempt to discuss the relative extent of T-electron delocalization in transition metalthiocarbonyl complexes on the basis of CS intensity measurements. CS-intensity studies would also permit comparisons between analogous series of metal-carbonyl and metal-dimitrogen complexes.

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ABPENDICES

PART V

Appendix A

PROGRAMS FOR INTEGRATED INTENSITIES

11

Program I PURPOSE (1) TO COMPUTE THE AREA UNDER THE APPARENT OPTICAL DEWEITY CURVE (FROM BAYD WEIGHT DATA) C C (2) TO DETERMINE THE APPARENT INTEGRATED ABSORPTION INTENSITY C g C DESCRIPTION OF PARAMETERS C C - # BANDS 1 - + CONCAS C Ľ. CALIBRATION OF RECORDSR (CN-1/INCH) C . R SAMPLE PATPLENGTH (NI) C PL -- MOLECULAR VEIGHT SOLUTE (CM) C M - CONC (CH SOLUTE/LITER SOLUTION) - CONC (HILLINOLES SOLUTE/LITER SOLUTION) C CC C C C CP - CALIBRATION OF CHART PAPER (GN/INCH) C ¥ YEIGHT OF SAND (GN) . - AREA UNDER APPARENT OPTICAL DENSITY CURVE (CH-1) C AREA C - ADSORBATICE AT BATTO MAXIMUM (TO MASE 1) AT - LN(T\$/T)VHAX (OR 2.3\$3.AT) С T C - APPARENT INTEGRATED INTENSITY (M-1 CH-2)-14-4 B C RSAD(5,1)J,K,R,PL,VM,CP 1 FORMAT(212,257.4,2510.5) VRIT3(6,3) FORMAT(1:0,4x,'R', 10x, 'PL',9x,'WR',9x, 'CP') WRITS(6,4)R,PL,VM,CP 3 FORHAT(110,1-10.4,1-11.4,1-14.5,1-9.5) 4 WRITE(6,5) 5 FORNAT(11) WRITE(6,6) 6 FORMAT(110,4x, 'C', 10x, 'AT', 9x, 'T', 10x, 'AREA', 7x, 'B') DO 10 1-1, J WRITE(6,7) 7 FORMAT(1H#). DO 20 L=1, X RIAD(5,8)CC, V, AT 8 FORMAT(2F6,5,1F4,3) C=(CC*10**3)/Xt T-2.303*AT AREA=2.383+R+¥/CP R=AHEA/(2+PL) 20 WRITE(6,9)C, AT. T. AREA. B FORMAT(1+#, 1F1#.4, 1*1#.3, 1F12.4, 1F11.4, 1F12.4) 9 19 **ZUNITHOD** END

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Program II

PURPOS3 TO CONVERT FROM TRANSMITTANCE TO OPTICAL DENSITY VALUES (2) TO COMPUTE THE AREA UNDER THE APPARENT OPTICAL DESSITY CURVE BY SIMPSON'S RULS TO DETERMINE THE APPARENT INTEGRATED ASSORPTION INTENSITY (3) C DESCRIPTION OF FARAMETERS C - # POINTS MEASURED NDIN C C R CALIBRATION OF RECORDER (CM-1/INCH) - SAMPL ATHLENGTH (MM) - MOLECULAR WEIGHT SOLUTE (GM) Ĉ PL ٧M C CONCY GE SOLITE/LITER SOLUTION CC CONCY (MILLINOLES SOLUTE/LITER SOLUTION) C . CONCINS HEASURED ĸ # SECTIONS WITH A DIFFERENT INCREMENT OF ARGUMENT N VALUSS IN A BA"D С HI # INTERVALS/INCH IN A PARTICULAR SECTION H = INCREMENT OF ARCUMENT VALUES (CN-1) - BASELINE MEASURED FROM SPECTRUM C BL - HEIGHT OF PEAK AT MAXIMUM ABSORPTION C PH C P(I)- POINTS TAKEN FROM CURVE С ТØ - TRANSMITTANCE OF BASELINS (4) C TM - TRAVISNITTANCE OF PEAK AT MAXIMUM APSORPTION (4) - OPTICAL DEVENTY OF CURVE AT MAXIMUM ABSORPTION C TL C - LN(TØ/TM)VMAX C T(I) - TRANSMITTANCE OF POINTS ON CURVE (4) Y(I) - OPTICAL DENSITY OF POINTS ON CURVE C INPUT FUNCTION VALUES С 2(1) - RESULTING INTEGRAL VALUES (CM-1) C - SUN OF INTEGRALS (CM-1) Ĉ SUN C - AREA UNDER APPARENT OPTICAL DENSITY CURVE (CH-1) AREA C B - APPARENT INTEGRATED IFTE SITY (N-1 CH-2)*1#**4 C DIMENSION P(100), T(100), Y(100), Z(100) READ(5,1)K, ", R, PL, W 1 FORHAT(212,2F6.4,1-14.4) WRITE(6,3) 3 FORMAT(1H#,4X, 'R',1#X, 'PL',9X, 'WN') WRITE(6.4)R.PL. WN 4 FORMAT(1"\$, 1F1\$.4, 1F11.4, 1F13.4) DO 30 J=1,K WRITE(6,18) 18 FORHAT(1HØ) READ(5,11)CC, "L, PH 11 FORFAT(1=6.4,2=4.1) C=(CC+1Ø++3)/VA TØ-101.-PL TH=100.-PH TL-ALOG(TS)-ALOG(TK) WRITE(6.8) B FORMAT (180,4x, 'C', 10x, 'BL', 9x, 'PN', 9x, 'TP', 9x, 'TN', 9x, 'TL') WRITE(6,9)C, BL, PH, TØ, TH, TL 9 FORMAT(1H8, 1F18, 4, 188. 1, 1512. 1, 2511. 1, 1513.4) AREA-. DO 29 L=1, 4 READ(5, 13) VDIM, HI 13 FORMAT(113, 1F4.1) H=R/HI WRITE(6,14) 14 PORNAT(1H#, 4x, 'NDIR', 7x, 'HI', 9x, 'H') VRITE(6, 16) "DIN, HI, H 16 FORMAT(188, 116, 1512.1, 1514.4)

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E.

	READ(5.2)(P(I), I=1, VDIH)
2	FORMAT(1174.1)
-	WRITE(6.12)
12	PORNAT(1HM)
	TO IS THE YOUN I
	T(T)-161 -B(T)
	$\mathbf{v}(\mathbf{x}) = \mathbf{v}(\mathbf{v}) = \mathbf{v}(\mathbf{x})$
24	
18	
	CALL (SS(H,T,Z,NDIN)
•	SUE-Z(NUIA)
1 4	WRITE(0,7)L, SUN
<u> </u>	FORMAT(1HD,4X, 'AREA UNDER SECTION',112, ' - ',178
- 2,6	AREA=AREA+SUM
	B-AREA/(C°PL)
Зø	VRITE(6,15)AREA,3
- 15	FORMAT(1H\$,4X, 'AREA = ',1"8.4,1\$X, 'B = ',1"8.4)

.

END

Program III

C.	PURPOSE
C	(1) TO DETERMINE THE ABSOLUTE INTEGRATED ABSORPTION INTENSITY BY RANSAY'S
С	HETHOD FOR OVERLAPPING BAND SYSTEMS (R. N. JOYES, D. A. RANSAY, D. S.
Ċ	KEIR, AND K. DOBRINSR. J. AH. CHEN.SOC., 74, 80 (1952))
Č	
č	DESCRIPTION OF PARAMETERS
č	
č	
č	
č	
ž	
č	ANU - AND ALL DIALOS FOR ANY AND
č	NTG - FRA A TALOG FOR LACE DA. D. NTEP - DEVIATION OF SAUL OCCUPATION STON NEAR & VASIE
č	DIF:
	DIFF "Z " VARIA UL V" DALA UDDUTALIUN "RUN RDA V D VALUB
U O	an a statute astatio of a table for the state of a state
<u> </u>	
C .	= 15A + B + 14U + 2 + 0 + 16A + B + 14U + 16A + B + 14U + 16A +
U A	
G /	Contraction and a dia
	RSAD(5,2)(8(1),1=1,K)
	2 FORMAT(6F6.4)
	WRITS(6,3)
х.	$3 \text{ FORMAT}(1:0, 3X_{0}, 8, 9X_{0}, B - AVG^{+}, 4X_{0}, (B - AVG)^{+}, 2)$
•	VRITE(6,4)
•	4 FORMAT(1RD)
	SUX-P
	SUP-SUH+B(I)
	10 COVTINUE
	AVG-SUN/K
•	
	DIFF-B(1)-AVG
	DIFFSQ-DIFF**2
•	WRITE(6,9)3(1), DITE, DITESQ
1	9 FORMAT()~19.4)
•	26 DSQ=DSQ+DIFTSQ
1	SD-SQAT(DCQ/K)
	A=1.62*AYG
	₽=4/2-3∮3
,	

.4)

- VRITE(6,6) 6 FORMAT(1HØ, 3X, 'SUK', 7X, 'AVG', 7X, 'SD'/8X, 'A', 9X, 'E') VRITE(6,7)SUM, AYG, SD, A, £ 7 FORMAT(1HØ, 1F8,4,1:11.4, 3F1Ø,4) VRITE(6,8) 8 FORMAT(1HØ) 30 CONTIVUE CO TO 1ØØ 101 STOP EVD

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Program IV

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C	PURPOSE	• •
C	(1)	TO DETERMINE THE ABSOLUTE INTEGRATED ABSORPTION INTENSITY BY
Ċ	• • •	(A) THE VIESON-WILLS EXTRADOLATION NETHOD (* & VIESON IN AND
č		
6.		A. J. TSULS, J. GALA, PAIS, 14, 5/6 (1940)
C >	•	(B) THE RUSSELL-THOMPSO" EXTRAFOLATION METHOD (R. A. RUSSELL AND
C		H. W. THOMPSON, SPECTROCHIM, ACTA, 9, 133 (1957))
C		Y CONTRACTOR OF A CONTRACTOR OF
C	DESCRIPT	ION OF PARAMETERS
č		
ř	4.1	
	5	
5	K	
C	Т	- (A) CONCY(MILLIMOLIS SOLUTE/LITER SOLUTION) (WILSON-WELLS)
C .		- (B) LY(TØ/T)VHAX (RUSSELL-THOMFSON)
C	B	- APPAREVT INTEGRATED INTENSITY (M-1 CM-2)+1#++4
C	TERCEP	- INTERCEPT - A
Ċ	-	- ABSOLUTE INTEGRATED INTENSITY (M-1 CM-2)+16++4
č	•	= 1/2.363 (H-1 (**-2)+) #**
č	CAT CB	- RECOVER THE CALLY AND FLORE (N-1 CH-2) - Stand
~		- DUNTETON OF STON OF CONTRACT SUITE CLARAN
G,	D145	- DEVIATION OF SACH DESERVATION SCIATION
C	DIFF.+-S	- VARIA C.S OF EACH OPSENVATION EQUATION
C,	SDFD	- STANDARD DIVIATION OF PARENT DISTRIBUTION
ର୍	SDSL	- STANDARD DIVIATION OF SLOFE
C`	SDINT	- STANDARD DEVIATION OF A VÁLUE
C	SDS	- STANDARD DEVIATION OF E VALUE
č	PSDSL	- STANDARD DEVIATION OF SLOPE
ř	PSDINT	
Š v	1001.11	- A STANDARD DAVIATION OF A VALUE
U I	·	
	DINESIO	N T(10), n(10)
10	B READ(5, 1	, END=101)J,K
	1 FORMAT(2	13)
1	DO 30 L-	
•	READ(5.2	(T(T), T-1, K)
	BRAD(5.2	
	אודגייגרת ב	
	1011010	
•		J/ av int av int av tattont av themat fy thematal
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FURPOSE
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(1) TO COMPUTE THE VECTOR OF INTEGRAL VALUES FOR A GIVEN EQUIDISTANT
    TABLE OF FUTCTION VALUES.
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DESCRIPTION OF PARAMSTERS

Ħ	THE	INCREMENT OF	ARGUMETT	VALUES.
Y	THE	THEUT VECTOR	OF FUTCT	TON VALUE

- THE IMPUT VECTOR OF FUTCTION VALUES. z
 - = THE RESULTING VECTOR OF INTEGRAL VALUES. Z MAY BE IDENTICAL WITH Y.

1

NDIN THE DIMENSION OF VECTORS Y AND Z.

REMARKS

NO ACTION IN CASE NDIN LESS THAN 3.

SUBBOUTINES AND FUNCTION SUBPROCRAMS REQUIRED TOTE

NETHOD

BESING WITH Z(1)-0. EVALUATION OF VECTOR Z IS DONE BY KEAVS OF SIMPSOV'S RULE TOGETHER WITH NEWTOV'S 3/8 RULE OR A COMPLIATION OF THESE TWO RULES. TRUNCATION ZAROR IS OF ORDER H**5 (I.E., "OURTH ORDER HETHOD). OULT IN CASE NDIN-3 TRUVCATION EBBOR OF 2(2) IS OF ORDER H**4. FOR RETERENCE, SEE

- (1) F.B. HILLEBRAND, INTRODUCTION TO NUMERICAL AMALYSIS, MOCRAY-HILL, TEN YORK/TOROWTO/LONDON, 1956, PP., 71-76.
- (2) R. ZURHUEHL, FRAKTISCHE MATHEMATIK FUER INGEVIEURE UND PHYSIKER, SPRINCER, BERLIY/COETTINGEN/HEIDELBERG, 1963, PP. 214-221.

SUBROUTINE QST(H,Y,Z, YDIK)

DIMENSION Y(1), Z(1)

C

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HT-.3333333*H
IF( TDIM-5)7.8.1
```

NDIN IS GREATER THAN 5. PREPARATIONS OF INTEGRATION LOOP 1 SUM1=Y(2)+Y(2) SUR1-SUM1+SUM1

SUN1=HT*(Y(1)+SUP1+T(3)) AUX1=Y(4)+Y(4) AUX1=AUX1+AUX1 AUX1-SUY1+:T*(Y(3)+AUX1+Y(5)) AUX2=TT=(Y(1)+3.875*(Y(2)+T(5))+2.625*(Y(3)+T(6)) SU12=1(5)+1(5)SU12-SU12+SU12 SUP2=AUX2-RT • (Y(4)+SUK2+Y(6)) Z(1)-ø.

AUX=Y(3)+Y(3)

AUX=AUX+AUX 2(2)-GUM2-HT+(Y(2)+AUX+Y(4)) 2(3)-GUM1 2(4)-GUM2

17(NDIM-6)5,5,2

"SCIENTIFIC SUBBOUTINE PACKACE," IBN CO. LTD., pp 291-2.

C INTEGRATION LOOP 2 DO 4 1-7, NDIN,2 SUN1-AUX1 SUH2-AUX2 AUX1=Y(I-1)+Y(I-1) AUX1=AUX1+AUX1 AUX1=SUH1+HT*(Y(I-2)+AUX1+Y(I)) Z(1-2)-SUM1 IF(I-VDIM)3,6,6 3 AUX2-Y(I)+Y(I) AUX2-AUX2+AUX2 AUX2=SU}2+HT*(Y(I-1)+AUX2+Y(I+1)) 4 2(1-1)-SUH2 5 2(NDIH-1)-AUX1 Z(NDIN)-AUX2 RETURN 6 Z(NDIH-1)=SUM2 Z(NDIN)-AUXI RETURN C END OF INTEGRATION LOOP 7 IF(NDIN-3)12,11,8 C NDIN IS SQUAL TO 4 OR 5 8 SUN2=1,125+HT+(Y(1)+Y(2)+Y(2)+Y(2)+Y(3)+Y(3)+Y(3)+Y(4)) C SUM1=Y(2)+Y(2) SUM1=SUM1+SUM1 SUH1-HT*(Y(1)+SUH1+Y(3)) z(1)-#. AUX1-Y(3)+Y(3) AUX1=AUX1+AUX1 $Z(2) = SUM2 = MT \cdot (T(2) + AUX1 + T(4))$ IF(*DIM-5)10.9,9 9 AUXI-Y(4)+Y(4) AUX1=AUX1+AUX1 Z(5)=GIM1+HT+(Y(3)+AUX1+Y(5)) 10 Z(3)=GUN1 Z(4)=GUM2 RETURN C Ċ NDIM IS EQUAL TO 3 11 SU:11-HT+(1.25+Y(1)+Y(2)+Y(2)-.25+Y(3)) SUM2=Y(2)+Y(2) SUN2-GUK2+SUF2 2(3)-HT*(Y(1)+SUN2+Y(3)) 2(1)-#. 2(2)-SUN1 12 RETURN END

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INTERSITY HEASUREMENT PARAMETERS

Description of Terms: Tables 8-1 to 3-6

£.

VBAX	
Slope	

Appendix

denotes the frequency at maximum absorption.

represents: (1) the alope of the apparent intensity versus concentration plot (Wilson-Vells[®] method for determining the absolute integrated intensity of a band); or (2) the slope of the apparent intensity versus apparent peak optical density plot (Russell-Thompson^b method for determining the absolute integrated intensity). The errors listed are the standard deviations in the slopes from lesst-squares analyses of (1) the intensity versus concentration plots (Wilson-Wells); or (2) the intensity versus peak optical density plots (Russell-Thompson).

represents the spectral alit width at Vmax.

are the true and apparent half-intensity band widths.

is a ratio given in Ramsay's Table IF, and is expressed as a function of $dv_{\frac{1}{2}}^{\frac{1}{2}}$ and the ratio $s/\Delta v_{\frac{1}{2}}^{\frac{1}{2}}$.

signifies the frequency interval on either side of the band centre, u_{σ} (v_{max}), over which the integration was performed.

is the increment of the argument value used in the computation of band area by Simpson's Rule.

"Infrared spectra were recorded at at least four, usually five different concentrations for each sample. The values reported in the tables are the average values for the five or so different concentrations.

⁸E. B. Vilson, Jr., and A. J. Vells, <u>J. Chem Phys., 14</u>, 578 (1946).

^bR. A. Russell and H. V. Thompson, <u>Spectrochim, Acta</u>, <u>9</u>, 133 (1957).

^CD. A. Ransay, <u>J. AB. Chem. Soc.</u>, <u>74</u>, 72 (1952).

- ,	- - -	V	Slòpe,	H ⁻² cm ⁻²		۵۷#			Δν ^t	v-vo *	$\frac{\nu - \nu_0}{\Delta \nu_{\frac{1}{2}}^t}.$	Δν	$\frac{\Delta v_{\frac{1}{2}}^{t}}{\Delta v}$
Complex	Integration method	c=-1	Wilson- Vella	Russell- Thompson	ca ^{-1°}	ca ⁻¹	•	I	ca ⁻¹	ca ⁻¹	、	1	
α _r (α) ₆ α _r (α) ₆	Simpson's Veighing	1986.9 1987.2	- 4.83 [±] 1.72 - 5.14 [±] 1.18	- 3.50 ± 1.35 - 4.19 ± 1.00	2.81	3.75. 3.80	0.75	1.95 1.90	1.92 2.00	25 ×	13.0 12.0	0.42	4.6
	• • • • • • • • • • • • • • • • • • •	-		-		ł	ŗ		-	τ.	- • •	-	
•	, i , i	• • • • • • • • • • • • • • • • • • • •	•	· · ·			\$	· · ·	· . · .	• • •	· , · ,	``````````````````````````````````````	,• ,

Slope of Intensity versus Concentration Plot, Spectral Slit Width, Half-Intensity Band Width, Integration Interval and Increment of

TABLE B-1

Slopes of Intensity versus Concentration Plots, Spectral Slit Widths, Half-Intensity Band Widths, Integration Intervals and Increments of Argument Values (Simpson's Rule) for the CO Stretching Absorptions of the (Cycloheptatriene)X(CO), Complexes in Methylene Chloride Solution

	V. RAX	Slope,	≝ ^{−2} ca ^{−2}	• • •	۵۷		ANT at	∆v ^t	v-v0	<u>v-vo</u>	ر ک ۷	
, Complex	ca ⁻¹	Vilson- Vells	Russell- Thospson	cm ⁻¹	ca ⁻¹	7 3		ca ⁻¹	ca ⁻¹	- 1	a-1 ,	
 ۲,۴,8°2(۵۵)	1982.1	- 1.40 ± 0.30	- 1.79 ± 0.47	2.81	13.5	0.21	1.03	13.1	64	4.9	· ·	
	1916.0	- 1.33 + 0.26	- 2.38 - 0.61	2.70	19.5	0.14	1.02	19.1	- 72	3.8	•	
, i	1866.9	- 2.49 - 0.35	- 5.92 - 0.92	2.65	26.3	• 0.10	1.01	· 26.0	' 72	3.1		•{
G_H_Ko(00)_	1983.	- 1.15 - 0.58	~ - 0.96 ⁺ 0.56	2.81	12.4	• 0.23	1.04	, 11.9	30	2.5	2.5	4.8
78 - 3	~ 1912.	- 2.60 - 0.88	- 3.78 - 1.53	2.70	22.3	0.12	1.01'	22.0	50	2.3	2.5	8,8
	1880.2	- 1.33 - 0.12	- 2.37 - 0.16	2.66	27.3	0.10	1.01	27.0 3	50	1.9	2.5	10.8
ເ,ສ ,ສ(ໝ).	1981. <u>.</u>	- 2.28 + 0.32	<u>~</u> 2.50 ± 0.38	2.81	11.5	0.24	1.05	11.0	64	5.8		یہ ب
	1908.	- 1.22 - 0.38	- 2.17 - 0.68	2.70	21.0	0.13	1.02	20.6	72	3.5	-	
	1873. ₁	- 0.17 - 0.19	- 0.30 - 0.39	2.63	- 26.0	0.10	1.01	25.8	72	.2.8		
		· · ·					*****	······		.1	,	

2892 5 ...

TABLE B-2

TABLE 8-3

Slopes of Intensity versus Concentration Plots, Spectral Slit Widths, Half-Intensity Band Widths, Integration Intervals and Increments of Argument Values (Simpson's Rule) for the CO Stretching Absorptions of the (Bicycloheptadiene)K(CO), Complexes in n-Hexane Solution

, c	v	Slope,	n ^{−2} ca ^{−2}	۲ ۲	۵۷		$\frac{\Delta v_{\frac{1}{2}}^{a}}{\Delta v_{1}^{t}}$	Δv ^t	v-v ₀	$\frac{v-v_0}{\Delta v_1^t}$	Δv	Δν ^t Δν
Complex ·	ca ⁻¹	Wilson- Vells	Russell- Thompson	° c=-1	ca ⁻¹	<i>*</i> *•	7	c-1	c=-1		c= ⁻¹	
<u>ποτ</u> -μ ₃ ατ(ω)	2033.6	- 0.14 ± 0.06	- 1.05 ± 0.51	2.90	3.73	0.78	2.30	1.62	20	12.4	-0.42	3.9
(•	1958.7	- 0.12 - 0.07	- 1.50 + 1.02	2.77	4.51 -	0.61	1.58	2.86	15	5.2	0.83	3.4
· · · ·	1944. 4	- 0.25 + 0.33	- 0.58 + 1.09	2.74	4.66	0.59	1.38	3.3 ⁸	20	5.9	0.83	4.2
•	1914.8	- 0.18 - 0.06	- 0.74 - 0.23	2.70	4.98	0.54	1.33	3.74	25	6.7	0.83	4.5
 <u>همت</u> -۳۵۲ ^۹ ۶۲۵	2044. h	+ 0.00 ± 0.07	+ 0.08 = 0.52	2.92	3.74	0.78	2.50	1.50	10	, 6.7	0.42	3.6
	1959.	- 1.28 + 0.16	- 3.19 ± 0.32	5 2.73	4.98	0.56	1.33	3.74	20	5.4	0.83.	4.5
	1913.6	- 0.55 - 0.14	- 2.01 - 0.59	2.70	5.14	0.53	1.33	3,86	25	6.5	° 0.83	4.8
nor-C,HgV(CO)4	2043.9	- 0.01 - 0.01	- '0.10 + 0.11	2.91	4.00	9.73	2.08	1.95	, 32	16.4	, 	•
	¹⁹⁵⁷ ·3	- 1.74 - 0.35	- 4.32 - 0.83	2.77	4,75	0.58	1,36	3.50	32 ·	9.2	ı	
	1910.1	- 0.24 + 0.12	- 9.85 - 0.41	2.70	4.75	,0.57	1.40	3.40	24	7.1		·

• •	WRAX		° Slope,	× ^{−2} ca ^{−2} -	• • • •	۵۷ ^a			۵۷ţ	¥-V0	<u>v-vo</u> <u>Av</u>	۵۷	۵۷ ۲۷
Complex	ca ⁻¹	,	Vilson- Vells	Russell- Thompson	a-1	ca ⁻¹		2	ca ⁻¹	, ca ²¹		ca ⁻¹	
e-Ko(0)3(PC13)3	2040.2		- 6.12 - 1.56	- 2.10 - 0.42	2.91	3.58	0,81	2.50	1,43	20	14.0	0.42	3.
	1991.1		- 8.25 - 2.35	- 2.45 - 0.78	2.81	5.75	0.49	1.29	4,46	25	5.6	0,83	5.
	2010.	4.	- 0.2h + 0.h0	* *	2 01	•		1 06	30.6	b 0	~ - •	1 9	
······································	1936.6	ø	- 0.15 + 1.30	+ 0.08 + 1.56	° 2,81	27.0	0,10	1.01	26.8	60	2.2	2.5	_10.
و-۲۰۰(۵۵) ر ا ۲(۲ ₆ ۲ ₄) ۲۵٫	~ 2006. ₀		- 0.36 - 0.19	- 0.44 = 0.26	, 2 .8 4	13.3	đ.21	1.03	12.9	45	3.7	•	
	1940.8		- 0.01 + 0.32	+ 0.01 = 0.53	2.73	33.5	0.08	1.01	33.2	96	2.9	ť	ć
	[°] 1970		= 0.17 + 0.84	<+ 0.36 ⁺ 0.80	2.79	14.0	0.20	1.03	13.6	48.	3.5	ŝ	84. · 1
E	1889.	\$	+ 0.97 = 2.56	+ 3.18 - 3.47	2.68	37.0	0.07	1.00	37.0	96	2.6	•	$\hat{}$
n-Mexane solution		, ,					·····		,		-	<u></u>	
· · ·				٠		<u> </u>		•	~	I	!		•
4- 9			\										۵

TABLE B-4

. ~	VEAX	Slope,	H ⁻² ca ⁻²		Δν å ,		$\frac{\Delta v_{\frac{1}{2}}^{2}}{\Delta v_{\frac{1}{2}}^{2}}$	∆v}	V-V0	<u>v-vo</u> <u>Av</u> t	۵۷	Sut -
Complex	c=-1	Wilson- Wells	Russell- Thospson	_ cs ⁻¹	ca ⁻¹	,	T	ca ⁻¹	cm ⁻¹	•	. ca ⁻¹	
<u>eis</u> -Xn(@) ₄ [P(C ₆ H ₅) ₃]er	2089.2	- 0.09 + 0.10	- 0.15 - 0.16	3.00	4.53	0.66	1.70	2.66	20	7.5	0.42	6.;
· · · · · · · · · · · · · · · · · · ·	2024.6	- 0.49 = 0.24	- 0.57 - 0.28	2.88	6.88	0.42	1.18	5.83	20	3.4	0.83	7.(
,	2006. ₄	- 1.37 ± 0.61	- 1.12 ± 0.49	2.84	7.50	0.38	1.14	6.58	25	3.8	0.83	• 8.(
3	1960.8	- 0.59 ± 0.10	- 0.89 - 0.15	2.78	11.6	0.24	1.05	11.0	25	2.3	0.83	13.3
10-Ma(00) [As(C6H5)]Br	2091.8	+ 0.11 = 0.07	+ 0.17 + 0.10	3.00	4.53 -	0.66	1.74	2.60	20	7.7	0,42	6.:
	2023.7	- 0.92 + 0.47	- 0.82 - 0.50	2.88	7.03	0.41	1.17	6.00	20	3.3	0.83	7.1
	2011.4	- 0.12 - 0.48	- 0.05 = 0.36	2.84	7.50	0.38	1.14	6.48	25	3.9	0.83	.7.1
	1962-3	+ 0.41 ± 0.40	+ 0.59 + 0.52	2.78	10.9	0.25	1.04	10.5	20	1.9	0.83	12.7
<u>не-на(</u> СО) ₆ [Sb(C ₆ H ₅) ₃]вт	2087.5	- 0.44 + 0.12	- 0.63 ± 0.18	3.00	4.8 4	0.62	1.58	3.06	20	6.5	0.42	7.:
	2021.6	- 0.89 ± 0.10	- 0.93 ± 0.12	2.88	6.56	0.44	1,20	[•] 5.48	20 '	. 3.6	0.83	, 6.(
ı	2009.2	- 1.72 = 0.37	² 1.37 [±] 0.29	2.84	7481	0.36	1.13	6.91	25	3.6	0.83	6.
•	1960.	- 0.36 ± 0.30	- 0.45 - 0.39	2.78	10.9	0.25	1.04	10.5	20	1.9	0.83	12.

TABLE B-5

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The state of the state of the

Slopes of Intensity versus Concentration Plots, Spectral Slit Widths, Half-Intensity Band Widths, Integration Intervals and Increments of Argument Values (Simpson's Rule) for the CO Stretching Absorptions of the (Π-Cyclopentadienyl)Kn(CO)₃ Complexes in <u>n</u>-Hexane Solution and for the CO and CS Stretching Absorptions of the (Π-Cyclopentadienyl)Hn(CO)₂CS Complex in Carbon Disulphide Solution

	v _{nax} .	Slope,	M ⁻² cm ⁻²	8	Δv		$\frac{\Delta v_{j}^{a}}{\Delta v_{1}^{t}}$	۵vi	۷-۷ 0	<u>v-vo</u> <u>Av</u> t	۵۷	
Complex	ca ⁻¹	Vilson- Vella	Russell- Thompson	a-1	a-1		- 1	cm ⁻¹	cz ⁻¹	•	cm ⁻¹	
π-C ₅ H ₅ ×n(ω) ₃	2030.8	- 0.10 - 0.06	- 0.69 ± 0.45	2.90	3.91	0.74	2.10	1.86	20	10.8	0.42	4.4
	1948.6	- 0.73 ± 0.10	- 2.31 + 0.30	2.76	5.31	0.52	1.29	4.11	30	7.3	0.83	5.0
(m-CH3C5H4)Hn(00)3	2026.7	+ 0.06 + 0.07	+ 0.37 + 0.42	2.89	3.91	0.74	2.05	1.91	20	10.5	0.42	4.5
	1944.5	- 0.95 - 0.35	- 3.75 [±] 1.41	2.74	6.88	0:40	1.15	6.60	.30	4.6	0.83	8.0
12-C5H5Hn(00)208	2009.7	- 0.08 ± 0.01	- 0.91 - 0.15	2.84	6,96	0.41	1.15	5.65	32	5.7		
	1956.9	- 0.11 - 0.02	- 1.28 - 0.23	2.77	8.72	0.32	1.09	8,00	72	9.0		
·	1263.8	- 0.11 ± 0.01	- 1.62 + 0.22	2.14	9.52	0.23	1.04	9.15	56 -	6.1		

TABLE B-6

Appendix C

INTEGRATED INTENSITY RESULTS

Description of Terrs: Tables C-1 to C-19

V.RAX

signifies the frequency at maximum absorption.

Conca

represents the concentration of solute.

are the apparent intensities of incident and transmitted rediation at maximum absorption.

 $\log_{10}(\frac{T_0}{T})_{v_{max}}$, $\ln(\frac{T_0}{T})_{v_{max}}$ are the apparent absorbance at maximum absorption and the conversion

Band weight*

Band area*

R •

denotes the weight of the absorbance curve profile cut-outs.

to the Maperian base, the apparent peak optical density.

denotes the area under the apparent optical density curve and is expressed by $\int_{band} \ln(T_0/T)_v dv$.

represents the apparent integrated absorption intensity defined as $(1/cl)\int_{band}\ln(T_0/T)_{y}dv$, where c is the concentration of solute in moles per liter, 1 is the cell path length in cs, and T_{0y} and T_{y} are the apparent intensities of the incident and transmitted radiation when the spectrophotometer is set at frequency v.

equals the absolute integrated absorption intensity defined as $(1/cl)\int_{band}\ln(I_0/I)_v dv$, where I_{O_V} and I_V are the incident and transmitted intensities of monochromatic radiation of frequency v.

is the absolute integrated absorption intensity expressed in practical units (i.e., $\Lambda/2.303$).

Four individual recordings of the infrared spectra were taken at each concentration; the values reported in the tables are the average values for the four recordings.

Ransay

A is computed by taking the mean of the 3 values obtained at a number of different concentrations and increasing it by 2%. The errors listed are the standard deviations from the mean.

Wilson-Wells

A is evaluated by extrapolating the B values taken at a number of different concentrations to zero concentration. The errors listed are the standard deviations in the intercepts, A. from least-squares analyses of the B versus c plots.

Russell-Thompson

A is determined by extrapolating the B values taken at several different concentrations to zero apparent peak optical density. The errors listed are the standard deviations in the intercepts, A, from least-equares analyses of the B versus $\ln(7_D/T)_{\rm V_{BBX}}$ plots.

^aR. N. Jones, D. A. Ramsay, D. S. Keir, and K. Dobriner, <u>J. Am. Chem. Soc.</u>, <u>74</u>, 80 (1952). ^bE. B. Wilson, Jr., and A. J. Nells, <u>J. Chem. Phys.</u>, <u>578</u> (1946). ^cR. A. Russell and H. W. Thompson, <u>Spectrochim. Acta</u>, <u>9</u>, 133 (1957).

wax Concn		T _{övrax}	T _{VRAX} .	lr:(To)	Bani area	10 ⁻⁴ B	1	0 ⁻⁴ A, R ⁻¹ c	- ²	10 ⁻⁴ 2, x ⁻¹ cs ⁻²			
ca ⁻¹	د He	*	۲		ca ⁻¹	x-1 cm-2	Ransay	Vilson- Vells	Russell- Thompson	Parsay	Vilson- Vella	Russel)- Thospao:	
1986.	0.654	94.7	34.7	ì.∞	4.90	66.8	65.1	68.7	68.7	28.3	29.5	29.9`	
,	0.832	95.5	29.9	1.16	5.65	63.4	± 1.4	* 4. 3	* 4. 7	± 0.6	± 1,9	- 2.0	
	0.929	97.1	27.4	1.27	6.37	64.0		٠					
	1.09	· 97.2	21.5	1.51	7.27	62.7			,	۰,		-	
	1.19	96.4	18.4	1.66	5.00	62.9					**		
	1.35	97.7	15.2	1.86	9.18	63.1		,					
Colecular .	weight =	220.1 g mo	1-1			Chart scale = 4	.96 cm ⁻¹ in	ch ⁻¹		Ce11	path length	- 0.107 m	
<u>ىرى بەتەرىپ بويۇلۇشۇر.</u>		÷,									~		

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				1.0 1.1 1.1	•	- 0.307 =		
	· • •	₩ ₩3, ₩-1 G	allees Kellts	29.9 1 1 - 1		tth leat	,	
		iching Zetho 10	frank		•	611 7	· ·	
) (K		Russell- Thompson	6 .2		lach ⁻¹	• • •	•, 3 -#
•		kare Solstie b, x ⁻¹ a ⁻²	111500- . Kells	6°.5 - 2.9		• - 7.95 ca ⁻¹		,
ير ب ۵		10 ⁻¹ 10 ⁻¹	A	2° 38 € 1		Chart seal		
• • •	Thù ce	intee of C	x-1 3-5	62 57 57 56 56 57 5	(. (. (.) (.) (.) (.) (.) (.) (.	- 	,	
		aling Abso 2a ch	78	2 2 3 2 2 3 8 9	7.25 8.17 9.93	ei 3 252 0		q
	,	and Strett		0.071 0.069 0.123	0.113 0.128 0.140	- HOL	~ /~	
· · · · · · · · · · · · · · · · · · ·	· · ·	eastly of the $la(\frac{2}{T})_{hart}$		0.917 1.12 1.23	1.66	T: C	p	,
. , , , , , , , , , , , , , , , , , , ,		Integrated Int log ₁₀ (^T 0) _{wax}		5, 0 5, 0 5, 0	0. 65 26 26 26 26 26 26	220.1 6 ml ⁻¹	,	
) , , , , , , , , , , , , , , , , , , ,	ii Conce	5	0.694 0.27	1 2 8	wight -	\$ *	-
	·	-	78	1927-2		Kalecular	ı	
	، ،		,			da - Pilda	a a a strategy and the second strategy and a strategy and the second strategy and the second strategy and the second strategy and the second strategy and	

VRAX	Concn	locito (To) VEAX	$\ln(\frac{T_0}{T})_{v_{\text{DAX}}}$	Chart weight	Pand weight	Band area	10 ⁻⁴ B	10	4. H ⁻¹	-2	10	- ⁴ £, :: ⁻¹	ся ⁻²
ca ⁻¹	a it			g inch ⁻¹	6	cm ⁻¹	H ⁻¹ cm ⁻²	Ransay	Wilson- Wells	Russell- Thompson	Ransay	Wilson- Wells	Russell- Thompson
1992.1	1.05	0.460	1.06	0.235	0.258	16.2	14.7	14.0	16.1	16.4	6.10	6.98	7.14
*	1.58	0.628	1.45	0.295	0.361	22.6	13.7	÷ 0.5	± 1.2	± 1.7	± 0.22	± 0.53	± 0.75
	1.70	0.649	1,50	0.296	0.391	23.6	13.3					-	
	1.79	0.697	1.61	0.290	0.417	23.6	13.7		•				~
	1.85	0.710	1.64	0.296	0.432	26.Q	13.5	-		3			
	1.93	0.751	1.73	0.296	0.454	27.4	13.6						
1916.	1.05	0.262	0.603	0.285	0.279	17.5	15.9	15.4	17.3	17,6	6.68	7.50	7.64
v	1.53	0.362 -	0.834	0.2 3 5	0.396	24.8	15.1	± 0.4	± 1.0	± 1.3	± 0.17	± 0,45	± 0.56
	1.70	0.333	_0,882	0.296	0.433	26.1	14.7			-			
	1.79	0.406	0.935	0.290	0.462	28.4	15.2						-
	1.85	0.418	0.963	0,296	0.476	29.7	14.9						
	1.93	。 0.437	1.01	0.296	0.490	29.5	14.7						
1886.	1.05	0.219	0.504	0.285	0.237	18.0	16.4	15.0	18.8	19.1	6.51	8,17	8.31
7	1.59	0.311	0.716	0.285	0.333	24.0	,14.6	\$ 0.8	+ 1.4	* 1.7	+ 0.35	2 0.61	± 0.73
	1.70	0.331	0.762	0.295	0.418	25.2	14.2						
	1.79	0.351	0.808	0.290	0.437	26.9	14.4				,	4	
	1.85	0.361	0.831	0.296	0.465	28.0	14.5			•			
	1.93	02379	0.873	0.296	0.472	29.5	14.2						

V	Conen	T _{OVIIX}	T _{VREX}	ls(To)	Band area	10 ⁻⁴ B	1	0 ⁻⁴ A, M ⁻¹ c	a ⁻²	10) ⁻⁴ Е, К ⁻¹ с	a ⁻²
m -1 [°]	s.".	\$	۶		ca ⁻¹	н ⁻¹ ся ⁻²	Bansay	Vilson- Vells	Russell- Thompson	Ransay	Vilson- Vells	Russell- Thompson
1983.	0.652	97.3	41.5	0.852	12.6	16.6	16.6	17.5	17.6	7.23	7.62	7.63
0	0.982	97.5	31.6	1.13	17.2	17.0	± 0.5	+ 1.4	+ 1.6	* 0.22	+ 0.60	+ 0.71
	1.10	97.2	25.5	1.34	20.1	15.8		-	-	-	-	-
	1.29	97.4	19.9	1.59	24.4	15.5		• •				•
	1.40	93.8	19.0	1.65	25.3	15.7						
1912.	0.592	96.5	49.9	0.660	18.9	18.6	18.3	- 21.1	21.3	7.93	9.17	9.24
~	1.10	96.Z	43.3	0.795	- 22.9	. 18.0	± 0.8	± 2.5	* 3.1	= 0.35	± 1.07	± 1.33
	1.29	%.2	37.3	0.947	27.5	18.6						
	1.40	97.3	36.2	0.999	22.3	17.6			-			
-	1.54	97.0	33.1	1.09	29.5	16.6		۔ د	-	,		
1980.	0.882	96.2 -	53.2	0.592	18.4	15.2	18.0	19.3	• • 19.5	7.81	8.37	5.48
4	1.10	96.5	46.5	0.724	22.5	17.8	± 0.3	= 0.3	± 0.3	± 0.13	: 0.Y	± 0.13
	1.29	95.3	41.2	0.839	25.9	17.5	-	4				
	1,40	96.3	40.1	0,881	29.1	17.5						
	1.5	97.2	37.0	0.996	30.6 *	17.2						

V _{RAX}	Conca	$\log_{10}\left(\frac{10}{T}\right)_{v_{\text{max}}}$	$lr.(\frac{T_0}{T})_{v_{RBX}}$	Sand velght	2and area	10-43	1	0 ⁻⁴ , y ⁻¹ c	-2	<u>ر</u> اد) ⁻⁴ 2, ² ⁻¹ cz	- ²
c ⁻¹	27			6	,cn ⁻¹	^{n−1} cs ^{−2}	Ransay	Vilson- Vells	Russell- Thompson	Parsey	Tilson- Zells	Russell- Thompson
19?1.,	1.02	0.479	1.10	0,275	17.0	16.1	15.5	13.2	15.6	6.72	7.90	3.07
-	1.16	0.542	1.25	0.304	19.7	15.5	= 0.5	± 1.0	± 1.3	± 0.22	- 0.45	± 0.36
	1.27	0.574	1.32	0.321	19.8	15.9		•			,	
	1.37	0.621	1.43	0.351 .	21.6	15.1				•		•
	1.50	C.679	1.56	0.372	23.3	14.9						
	1.63	0.716	1.65	0.400	z4.7	14.5						
1908.	1.02	0.290	0.645	0.302	12.6	17.6	17.3	18.6	18.7	7.50	8,06	. 8.13
,	1.16	0.317	0.730	0.332	20.5	16.9	± 0.3	÷ 1.2	± 1,4	* 0.13	÷ 0.53	= 0.69
	1.27	0.341	0.735	0.359	22.1	16.8						
	1.37	0.367	0.945	0.395	24.3	17.0				1		
	1.50	0.402	0.925	0,429	26.4	16.9						
٠	1.63 /	0,429	0.936	0.456	29.1	16.5			-	×		٨
1873.	1.02	0.243	0.560	0.233	17.5	16.5	17.0	16.9	16.9	7.37	7.33	4.2
1	1.16	0.274	0.631	0.325	29.0	16.5	+ 0.1	* 0. 6	± 0.7	÷ 0.04	+ 0.27	- 0.29
	1.27	. 0.296	0.692	0,357	22.0	16.7						-
	1.37	0.322	0.742	0.397	23.9	16.7						
	1.50	0. 353	0.913	0.425	26.2	16.8						
ء ر	1.63	0.372	0.857	0.456	29.0	16.5					x	````
7.elecula	x weight	- 360.0 g mol ⁻¹	Chart	wight = 0	.297 6 1	nch ⁻¹	Chart so	ale = 7.96 c	inch ⁻¹	Cell	path longth	- 1.04 m

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TABLE C	-6
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Y _{REX}	Concn	TOYEEX	T _{vnax}	$\ln(\frac{T_0}{T})_{V_{BBX}}$	Band Area	10 ⁻⁴ b	10	0 ⁻⁴ A, M ⁻¹	a ^{−2}	ľ	0 ⁻⁴ E, H ⁻¹	-2
ca ⁻¹	a.M	*	` *		ca ⁻¹	M ⁻¹ cm ⁻²	Bancay	Vilson- Vella	Russell- Thompson	Bensey	Vilson- Vells	Russell- Thompson
2033.	5.47	97.5	41.3	0.859	4.09	7.55	7.24	8.15	8.25	3.14	3.94	3.59
Ŭ	6.60	97.8	36.1	0.997	4.63	7.10	± 0.26	± 0.97	± 1.25	± 0.11	- 0.42	± 0.54
	7.08	97.7	34.1	1.05	4.73	6.25	•				**	
•	8.08	95.0	29.4	1.24	5.72	7.17					•	
-	9+43	97.6	25.8	1.33	6.36	6.83		>			n	
1958.	5.47	97.4	61.9	0.453	2.63	4.26	4.41	5.19	5.20	1.92	2.26	2.26
7	6.60	97.0	57.7	0.520	2.72	4.18	+ 0.27	+ 1.20	\$ 1.33	± 0.11	÷ 0, 92	\$ 0.58
	. 7.08 >	97.0	55.3	0.562	2.90	4.14						
	8.08 -	96.8	50.4	0.653	3.33	4.18					_	
	9.43	96.8	46.3	0.738	3.99	4.27					•	
1944.	3.17	98.4	31.1	1.15	6.67	21.3	21.4	22.0	21.8	9.29	9.57	9.49
•	3.94	99.8	24.8	1.39	8.35	21.5	± 0.6	\$ 2.9	± 3.2	\$ 0.26	÷ 1.24	- 1.42
	4.76	99.7	20.1	1.59	9.36	19.9			-			
	5.47	98.2	14.9	1.89	11,4	21.2						
1914.	3.17	98.4	40.5	0.658	5.61	17.9	17.9	18,4	18.5	7.76	· 7.99	8.02
	3.94	99.8	33.7	1.09	6.91	17.8	± 0.3	= 0.6	± 0.8 .	± 0.13	2 0.25	= 0.3+
	4.76	98.7	29.1	1.22	8,08	17.2	-	**	•	•	-	
	5.47	98.2	22.9	1,46	9.42	17.5					د	,
-	6.60	99.0	18.2	1.69	11.3	17.3			,			•
iolemier	weight =	256,1 g m	1-1 .		Che	rt scale = 4	.98 cm ⁻¹ in	ch ^{~1}	¹	Cell p	th length	· 0.0987 m

and the second of the second
v _{mex}	Conca	To _{vrax}	TVRAX	$\ln(\frac{T_0}{T})_{v_{BAX}}$	Band area	10 ⁻⁴ B	3	10 ⁻⁴ A, M ⁻¹ c	-2	. 1 0) ⁻⁴ E, 5 ⁻¹ a	÷2
ca ⁻¹	RM.	×	*		ca ⁻¹	H ⁻¹ cm ⁻²	Ransay	- Wilson- Wells	Russell- Thospson	Ransay	Vilson- Velís	Russell- Thompson
2044.	3.07	97.8	63.9	0.426	1.80	- 5.98	6.04	5.91	5.87	2.62	2.57	2.55
,	3.60	97.9	59.7 ·	0.495	2.14	6.04	± 0.15	± 0.57	* 0.59	± 0.06	± 0.25	± 0.25
	4.13	97.9	56.8	0.544	2.30	5.67						
	6.00	97.1	43.4	0.805	j.54	6.00						
1959.,	2.00	98.1	38.7	0.930	5.77	29.3	28,1	31.7	32.2	12.2	13.8	14.0
-	3.07	. 99. 5	24.3	1.40	8.25	27.4	÷ .1.0	± 1.0	± 0.9	± 0.4	= 0.4	± 0.4
	3.60	98.2	19.6	1.61	9.53	27.1						
	4.13	99.1	16.6	1.79	10.8	26.6			•		```	
1913.	2.00	99.2	53.2	0.623	4.22	² 21.5	21.2 [*]	22.8-	23.0	9.19	9.90	- 9.97
U	3.07	99.8	38.8	0.945	6,42	21.3	± 0,8	± 1.2	± 1.5	± 0.35	± 0.52	± 0.64
	3.60	99.8	33.7	1.09	7.60	21.5						د
	4.13	100.0	30.9	1.17	8.09	19.9				-		-
	6.00	99.2	18.3	1.69	11.6	19.5		A STANKER	د			
Molecular	r weight =	300.0 6 80	1-1			Chart scale -	4.99 cm ⁻¹ 11	nch ⁻¹		Cell ;	ath length	• 0.0983 m

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V _{RAX}	Concn	$\log_{10}(\frac{T_0}{T})_{v_{max}}$	$\ln(\frac{T_0}{T})_{v_{max}}$	Chart weight	-Eand weight	Band area	10 ⁻⁴ B	10	4, n ⁻¹ c	-2 m-2	10	-4 _{2, x} -1	ca ⁻²
ca ⁻¹	rM			g inch T	1. 6	هـ	M ⁻¹ cn ⁻²	Ransay	Wilson- Wells	Russell- Thompson	Ramsay	Vilson- Vells	Russell- Thompson
2043.0	3.13	0,159	0.435	0.292	0.034	2.17	6.47	6.48	6.43	6.43	2.81	-2.79	2.79
y	3.61	0.216	0.497	0.292	0.039	2.38	6.15	,= 0.13	± 0.19	+ 0.20	÷ 0.06	± 0.08	\$ 0.09
	3.90	0.230	0.530	0.292	0.043	2.70	6.45		4		-		•
	5.63	0.342	0.738	0.283	0,060	3.89	6.44		-				序
	13.3	0.735	1.69	0.2%3	0.138	9.92	6.25			•			
1957	2.56	0.520	1.20	0.292	0,127	7.96	29.0	28.0	33.2	34.0	12.2	14.4	14.8
5	3.13	0.623	1.44	0,292	0.147	9.24	27.5	± 1.0	÷ 2.3	± 2.5	± 0.4	± 1.0	÷ 1.1
	3.61	0.703	1.63	0.292	0.163	10.3	26.5	*	-	-	1		
	3.90°	0.756	1.74	0.292	0.179	11.3	26.9			· ·	,		
1910	2.56	0.356	0.220	0.292	0.083	5.21	19.0	18.7	19.3	19.4	8.13	5.37	9.41
· 1	3.13	0.427	0.933	0.292	0.099	6.19	18.4	20.4	± 1.0	± 1,1	± 0.17	± 0.43	+ 0.47
	3.61	0.469	·1.13	0,292	0.110	6.94	17.9						,
	3.90	0.518	1.19	0.292	0.123	7.69	18.4	٩					
	5.63	0.737	1.70	0.283	0.169	10.9	18.1						

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v _{fax}	Conen	Tovcax	Theax	$\ln(\frac{70}{T})_{v_{REX}}$	Band ' area	10 ⁻⁴ 3		10 ⁻⁴ л, н ⁻¹ с	a ⁻²	, 1	0 ⁻¹⁴ 2, M ⁻¹ c	a ⁻²
e ⁻¹	*	\$	E/	•	a ⁻¹	м ⁻¹ ст ⁻²	Pansay	Wilson- Vells	Ruscell- Thompsong	Ransay	Vilson- Vella	Russell- Thompson
2040.,	0.135	96.0	56.1	0.537	2.43	15.6	15.3	16.5	16.8	6.64	7.17	7.29
-	0,169	95.2	49.1	0.662	3.00	15.4	2 0.5	0.8	± 0.8	± 0.22	± · 0.34	± 0.35
	0.220	95.2	41.7	0.26	3.84	15.2				. ·		r
	. 0.270	96.0	34.7	1.02	4.42	14.2	۰ ۱			,		
-	0.338	95.8	29.0	1.29	5.63	14.5		,	-			
1991.,	0.135	97.0	59.2	0.495	3.60	23.2	23.2	24.6	24.7	10.1	10.7	10.7
•	0.169	96.2	51.8	0.619	4.61	23.7	± 0.7	± 1.2	± 1.4	± 0.3	·± 0.5	± 0.6
	0,220	96.2	44.3	0.775	5.67	.22.4			9		Ŷ	-
	0.270	96.9	.36. 7	0.971	7.09	22.8			, • -			•
	0.339	· 96.5	30.7	1.15	8.41	21.7						

v max	Concn	Tovzax	T. Vmax	$ln(\frac{T_0}{T})_{y_{\text{Bax}}}$	Band area	10 ⁻⁴ B	- 10	0 ⁻⁴ A, M ⁻¹ cm	-2	. 10	0 ⁻⁴ E, S ⁻¹ c	» ⁻²
c# ⁻¹	2M	*	*	۰ -	ca ⁻¹	N ⁻¹ cm ⁻²	Ransay	Wilson- Wells	Russell- Thompson	Ransay	Vilsen- Vells	Russell- Thompson
2039.0	0.510	95.8	92.7	0.598	8.16	13.9	14.0	14.0	13.9	6.10	6.07	6,04
,	0.679	95.5	43.3	0.791	11.0	14.1	\$ 0.3	± 0.9	± 1.0	± 0.13	1 0.40	± 0.43
	0.845	95.2	39.5	0.911	12.9	13.3			-			
	1,02 -	95.9	31.0	1.13	16.1	13.7			1		-	•
	1,18	95.9	25.4	1.33	18.9	13.9		•	ĩ		. 0	
1926.6	0.510	96.9	61.2	0.460	15.1	25.8	26.1	25.7	25.5	11.3	11,2	11.1
Ū	0.679	97.0	53.2	0,601	20.4	26.1	± 0.7	± 2.5	± 2.6	2 0.3	± 1.1	/= 1,1
	0.845	. 97.0	47.8	0.709	24.2	24.9	•					
	1.02	97.0	40.9	0.964 5	29.1	24.7				,	~	
	1.18	. 98.0	35.0	1.03	35.9	26.4					-	5

Vmax	Concn	$\log_{10}(\frac{T_0}{T})_{v_{\text{max}}}$	ln(To)	Band weight	Pand area	10 ⁻⁴ b		0 ⁻⁴ A, K ⁻¹ cm	-2 .	10	o ^{−4} £, X ^{−1} ლ	-2
ca ⁻¹	2 .	د ۲		ß	ca ⁻¹	M ⁻¹ cm ⁻²	Ransay	Vilson - Wells	Russell- Thompson	Ransay	Vilson- Yells	Russell- Thompson
2006.	1.12	0.432	0.995	0.251	15.5	13.2	13.3	13.6	- 13.6	5.77	5.90	5.90
, v	1.30,	0,500	1.15	0.298	17.7	13.1	± 0.2	± 0.7	± 0.8	± 0.09	= 0.32	ž 0.36
	1.43	0.538 -	1.24	0.313	19.3	12.9			•			
	1.62	0.614	1.41	0.364	22.4	13.3				u		
	1.79	0.655	1.51	0.399	24.0	· 12.9						
	1.95	0.717 .	1.65	0.42)	26.0	12.8	ه <u>ب</u> ه					
1940.2	1.12	0.319	0.735	0.474	29.2	24.9	25.3	* 24 . 8	24.8	11.0	10.e	10.8
0	1.30	0.369	0.850	0.542	33.4	24.6	± 0.2	± 1.2	± 1.3	± 0.1	± 0.5	÷ 0.6
	.1.43	0.402	0.926	0.596	36.7	24.7		·	\$	•		
	1.62	0.457	1.05 [°]	0.691	42.6	25.2				•	5	
_	1.79	0,496	1.14	0.746	46.0	. 24.7						~~
•	1.95	0.539	1.24	0.815	50.2	24.7	5					
Holecula	r weight	- 716.9 g mol ⁻¹	Chart	weight = (.298 8 1	nch ⁻¹	Chart sc	ale = 7.96 cm	-1 inch ⁻¹	Ce1:	1 path length	= 1.04 m

V _{RAX} Con	Conch	log10(To)VBAX	$ln(\frac{T_0}{T})_{v_{\text{max}}}$	Pand weight	Band area	10 ⁻⁴ B	× , 1	0 ⁻⁴⁴ A, Y ⁻¹ cm	-2	1	0 ⁻⁴ E, ^{p,-1} cm	-2
a ⁻¹ `	a #	•		6	ca ⁻¹	H ⁻¹ cm ⁻²	Ransay	Vilson- Vells	Russell- Thompson	Ransay	Yilson- Vells	Russell- Thompson
1970.	0.935	0.300	0.691	0.229	14.0	13.7	24.3	14.2	13.6	6.20	6.17	5.92
,	1.05	0.413	0.951	0.259	16.0	14.6	± 0.4	± 2.3	± 1.8	± 0.17	± 1,00	= 0.79
	1.21	0.442	1.02	0.291	17.3	13.7						
	1.36	0.497	1.15	0.324	20.0	14.1						
	1.43	0.546	1.26	0.350	° 21.6	14.0			,		•	
1889.1	0.935	0.261	0.601	0.426	25:3	25.6	· 27.9	26.2	24.8	12.1	11.4	10.9
•	1.05	0.323	0.744	0.509	31.4	29.7	± 1.I	± 7.0	± 6.2	= 0.5	· = 3.0	± 2.7
	1.21	0.348	0.901	0.563	34.7	27.5				۶		· ·
	1.36	0.337	0, 891	0.649	40.0	28.1						
•	1,49 1	0.425	0.979	0.673	41.5	26.9			L			
folecula	r weight -	941.0 g mol ⁻¹	Chart	weight = (.299 g 1	nch ⁻¹	Chart Sci	1e = 7.96 cm	-1 inch-1	Cel	1 path length	1 = 1,04 23

True West

TABLE C-12

* RAX	Coaca	Toyanx	T, YRX	In(TO) VERN	Band area	10	1	0 ⁻⁴ A, z ⁻¹ c	-2	1	0 ⁻⁴ 2, x ⁻¹ e	* -2
cs ⁻¹	Ma	× `	4		a ⁻¹	x ⁻¹ cm ⁻²	Resay	Filson- Vells	Passell- Thompson	Baraay	Vilson- Vella	Euroell- Theupeon
2099.	0.663	80.7	49.1	0.497	2.81	3.65	3.70	3.73	3.75	1,61	1.62	1.63
, 4	0.791	54,5	45.0	0.559	3.39	3.73	2 0.08	± 0.25	± 0.29	± 0.03	± 0.11	± 0.12
	1.09	80.0	36.6	0.782	4.45	3.53			-	,		
	່າ.ກ	80.2	32.1	0.916	5.38	3.53			v			
*	1.55 `	84.7	29.2	1.07	6.70	3.68						·
2024	0.663	79.6	39.8	0.693	5.76	7.56	7.48	7.87	7.96	3.25	3.42	3.45
	0.791	83.6	37.5	0,302	6.91	7.60	+ 0.25	+ 0,59	+ 0.66	= 0.11	÷ 0.26	+ 0.23
	1.09	75.7	26.7	1.00	9.17	7.28				•		
	1.33	78.1	22.0	1.27	10.5	6.91				•	•	-
	1.58	83.0	18.7	1.49	13.3	7.32						•
2006.	0. Sk7	79.8	15.3	0.516	7.87	32.5		13.2	13.4	5.13	\$75	5.82
	0.663	80.7	31.2	0.950	2 9.00	11.8	÷ 0.6	+ 1.2	± 1.4	÷ 0.26		: 0.77
	0.791	8.8	28.2	- 1.10	11.7	12.5						
,	1.09	80.0	17.7	1.51	14.8	11.8	•					
•	1.33	80.2	13.9	1.75	17.2	11.3						
1960.	0.663	80.7	48.3	0.513	6.67	8.74	8.76	9.23	9.30	3.80	4.01	4,04
8	0.791	84,5	46.5	0.601	8.09	8.89	= 0.21	± 0.24	± 0.27	= 0.09	: e.n	÷ 0.12
	1.09	80.0	35.2	0.821	10.8	8,58						
3	1.33	80.2	30.7	0.960	12.8	8.40			`			
	1.58	8.7	27.7	1.12	15.1	8.31		•				
Nelecular	: wight -	509.2 6 20	1-1			thart scale	- 4.98 ca ⁻¹ 1	nd ⁻¹		Cel	1 yeth longt	h = 1.15 m

V RAX	Conca	Toy	τ _ν	$\ln(\frac{T_0}{T})_{V_{\text{RAX}}}$	Band	10 ⁻⁴ B	1	0 ⁻⁴ A, f ⁺¹ c	-2	1	0 ^{-h} s, n ⁻¹ c	_ -2
ca ⁻¹	e)i	*	*		ca ⁻¹	K ⁻¹ cm ⁻²	Ransay	Wilson- Wells	Russell- Thompson	Ransay	Vilsen- Vells	Russell- Thompson
2091.	0.594	93.0	61.2	0.419	2.39	3.50	3.65	3.47	3.47	1.59	1.51	<u></u>
•	0.631	94.5	60.8	0.441	2.53	3.56	± 0.05	± 0.14	± 0.15	± 0.02	± 0.06	÷ 0.06
~	0.834	92.4	48.9	0.636	3.70	3.64	•				•	
	1.19	93.4	40.9	0.826	4.86	3.55				-		
	1.26	94.7	39.6	0.872	5.23	3.63		-				
2023.	0,594	92.0 ·	47.3	0.665	6.06	8,86	9.00	9.67	9.65	3.91	4.20	4.19
	0.632	93.5	45.B	0.714	6.59	9.07	± 0.39	\$ 0.96	± 1.11	± 0.17	÷ 0,42	1 0.48
_	0.884	91.4	32.8	1.03	9.54	9.39	-					
_	1.19	92.4	25.9	1.27	-11.4	8.32	-					
•	1.26	93.7	24.2	1.35	12.3	. 8,49		,				ũ
2011.,	0.442	92.2	49.0	0.632	5.88	11.6	12.0	11.9	11.8	5.21	5.15	5.13
•	0.594	93.0	40.1	0.841	7.92	11.6	± 0.3	± 0.8	2 0.8	\$.0.13	\$ 0.35	\$ 0.37
• `	0.632	94.5	37.9	0.914	8.90	12.1	-				•	
	0,882	92.4	25.9	1.27	12.3	12.1						
	1,19	93.4	18.1	1.64	15.7	11.5		\$				
1962	0, 594	93.0	53.2	0.469	5.47	8.00	8.69	8.14	8.09	3.77	3.9	3.51
3	0.632	94.5	56.2	0.520	6.33	8.72	+ 0.27	± 0.82	± 0,84	\$ 0.12	2 0.36	\$ 0.36
3	0.88	92,4	45.6	0.706	8,83	8.69			•			
	1.19	93.4	37.1	0.923	11.6	8.47		¢			· Sectores	•
	1.26	94.7	35.1	0.993	12.6	8.71		~				~
Helemie	r weight =	551.1 c =	1-1			(hert scale s	. A. 98 cm ⁻¹ 4	nch ⁻¹		Call	neth Janeti	- 1.15 =

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V _{RAX}	Conco	TOVRES	TVRAX	$\ln(\frac{T_0}{T})_{v_{\text{max}}}$	Band area	10 ⁻⁴ b	1	0 ⁻⁴ , H ⁻¹ c	n ⁻²	10	о ⁻⁴ я, я ⁻¹ с	n ⁻²
ca ^{-1 ·}	aM	×	×		ca ⁻¹	H ⁻¹ cm ⁻²	Rensay	Vilson- Vells	Russell- Thompson	Rantay	Wilso Wells	Russell- Thospyos
2037.	0.550	92.5	59.1	0.449	2.70	4.86	4.30	4.61	4.64	1.87	2,00	2.02
• >	0.633	92.2	56.1	0.497	3.21	4.40	± 0.14	± 0.25	+ 0.27	± 0.06	± 0,11	± 0.12
	0.933	94.9	46.2	0.720	4.61	4.29						
	1.10	92.5	40.8	0.819	5.16	4.05				•		
	1.27	91.5	35.4	0.950	5.86	4.02						
2021.	0.550	91.8	50.1	0.606	4.99	7.89	7.70	8.34	8,41	3.34	3.62	3.65
. 0	0.633	91.2	46.1	0.682	5.66	7.77	± 0.25	± 0.21	± 0.24	+ 0.11	+ 0.09	± 0.11
	0.933	.93.9	35.7	9.967	8.07	7.52						
	1.10	91.5	30.0	1.12	9.19	7.26					÷	-
	1.27	90.5	24.9	1.29	10.6	7.30						•
2009.2	0.467	94.7	46.9	0.703	7.15	13.3	12.8	13.8	14.0	5.57-	6.01	6.09
- •	0.550	92.5	41.0	0.814	8.14	12.9	\$ 0.5	\$ 0.6	\$ 0.7	± 0.22	\$ 0.27	\$ 0.30
	0.633	92.2	36.7	0.921	9.05	12.4					,	-
	0.933	94.9	26.1	1.29	13.1	12.2						4
	1.10	92.5	20.6	1.50	15.3	12.1						
1960.	0.550	92.5	57.9	0.469	5.95	9,40	9.23	9.37	9.37	4.01	N.07	4.07
•	0.633	92.2	54.8	0.520	6.43	8.83	± 0.20	± 0.59	± 0.64	\$ 0.09	.25	\$ 0.25
	0.933	94.9	44.1	0.766	9.90	9.13		•	~		\sim	
-	1.10	92.5	38.6	0.874	11.4	9.02			\sim	*		
	1.27	91.5	33.1	1.02	12.9	8,83	-					

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TABLE C-15

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V.RRX	Conen	Toveax	TVEAX	$ln(\frac{T_0}{T})_{V > ax}$	Eand area	l 10 ⁻⁴ 8	1	10^{-4} A, K^{-1} cm ⁻²			0 ⁴ 2, ^{x⁻¹} e	,-2
e -1	nt	×	*	0	cm ⁻¹	н ⁻¹ ся ⁻²	Bansay	Vilson- Sells	Russell- Thompson	Bansay	Vilson-	Eussell- Thompson
2030.8	3.45	99.0	54.3	0.601	2.85	8.27	8.41	8.68	8.75	3.65	3.77	3.90
•	3.92	93.9	50.6	0.670	3.28	ê. 39	± 0.12	- 0.64	± 0.75	= 0.05	- 0.28	- 0.32
	4,50	9º.7	46.8	0.746	3.76	8.34						
e	- 4.92	9º.3	43.7	0.811	3.96	8.04					,	
	5.42	9 ⁹ .7	40.6	0.898	4,44	8.19						•
10/19	2 66	0 6 0	35.0	. 0 013	6 61	25 R	25 6	76	, 27 Q	11 1	12 0	12 1
174-3-6	2.95	90.9 99.1	34.1	1.07	7.52	25.5	± 0.5	± 0.8	± 0.8	± 0.2	± 0.3	± 0,4
	3.45	100.1	29.3	1.23	8.56	24.8						<i>.</i>
/	5.92	100.0	25.6	1.36	9-77	24.9		-	, .			1
•	4.50	9 9.9	21.0	1.56	10.9	24.3		,				

$\begin{array}{cccccccccccccccccccccccccccccccccccc$	Russell- Thompson 8.17 ± 0.84	Ramsay 3.77 - 0.05	Vilson- Vells 3.55	Russell Thompson
20267 4.32 98.7 47.1 0.740 3.64 8.43 8.67 8.17 4.74 98.8 43.7 0.816 4.08 8.60 $= 0.12$ $= 0.82$ 5.24 99.4 40.5 0.898 4.36 8.31 5.81 99.6 37.1 0.988 4.99 8.59 6.21 100.0 34.7 1.06 5.33 8.58 1944 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0	8.17 ± 0.84	3.77 ± 0.05	3.55	3, 55
4.74 98.8 43.7 0.816 4.08 8.60 ± 0.12 ± 0.82 5.24 99.4 40.5 0.898 4.36 8.31 5.81 99.6 37.1 0.988 4.99 8.59 6.21 100.0 34.7 1.06 5.33 8.58 1944.2 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0	± 0.84	± 0.05	+	//
5.24 99.4 40.5 0.898 4.36 8.31 5.81 99.6 37.1 0.988 4.99 8.59 6.21 100.0 34.7 1.06 5.33 8.58 1944.2 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0			- 0.36	± 0.37
5.81 99.6 37.1 0.988 4.99 8.59 6.21 100.0 34.7 1.06 5.33 8.58 1944.2 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0				
6.21 100.0 34.7 1.06 5.33 8.58 1944 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0			υ	
1944. 2.78 99.9 42.5 0.855 7.50 27.0 25.9 29.0		,		
	29.5	11.2	12.6	12.8
3.33 99.9 37.1 0.991 8.29 24.9 ± 0.9 ± 2.9	÷ 3 <u>3</u> 5	± 0.4	± 1.3	± 1.5
· 3.81 100.5 32.3 , 1.14 9.66 , 25.4	¶î, 4			
4.32 100.0 29.2 1.23 10.7 24.8				•
4.74 100.3 26.0 1.35 11.8 24.8				
Nelecular weight = 218.1 g mol ⁻¹ Chart scale = 4.98 cm ⁻¹ inch ⁻¹		Cell	path length	= 0,100 m

TABLE C-1?

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TABLE C-18

V.RAX	Concn	log10(To)vFEX	ln(To)	Iand weight	^b and area	10 ⁻⁴ B	10	-4 _A , H ⁻¹ cr	-2	-	ο ⁻⁴ ε, ^{χ-1} α	- 2
ca ⁻¹	M.		~	8	<u>ca</u> -1	M ⁻¹ ca ⁻²	Ransay	Wilson- Wells	Russell- Thompson	Ransay	Vilsor- Vella	Ruczell- Thospaos
2009. ,	7.75	0. 391	0.901	0.126	8.95	10.8	10.6	11.3	11.5	4.62	4.92	5.01
•	9.25	0.447	1.03	0.147	10.5	10.6	± 0.2	± 0.4	· ± 0.5	± 0.09	± 0.17	± 0.20
	10.2	0.485	1.12	0.163	11.5	10.6		-	•			-
	12.1	0.562	1.30	0.189	13.3	10.2						
	13.7	0.624	1,44	0.211	14.9	10.2	<i></i>		~	-		
	15.5	0.692	1.57	0,240	17.0	10.2					t	
1956.0	7.75	0.330	0.975	0.169	12.0	14.4	14.1	15.1	15.4	6,14	6,56.	6.68
-	9.25	0.435	1.00	0.194	13.7	13.9	±• 0.3	± 0.5	\$ 0.7	± 0.13	± 0.23	+ 0.29
	10.2	0.473	1.09	0.215	15.2	14.0						
	12.1	0.549	1.26	0.254	19.0	13.8	٥					
	13.7	0.603	1.39	0.294	20.1	13.7						
	15.5	0.664	1.53	0.313	22.2	13.4						
Zolecula	r weight	- 220.1 g mol ⁻¹	Chart	: weight =	0.260 g	Inch ⁻¹	Chart sce	1e = 8.02 cm	⁻¹ inch ⁻¹	Ce11	path length	= 0.107 mm

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V. RAX	Conen	log10(To)VRAX	$ln(\frac{T_O}{T})_{V_{RAX}}$	Band weight	Band area	10 ⁻⁴ B	10	0 ⁻⁴ A, M ⁻¹ cm	-2	_	10) ⁻⁴ E, K ⁻¹ ci	-2
cn -1 °	R.			8	cn ⁻¹	H ⁻¹ cm ⁻²	Ransay	Yilson- Vells	Russell- Thompson	_	Ransay	Wilson- Wells	Russell- Thompson
1263.8	7.75	0.304	0.700	0.147	10.4	12.5	12.5	13.5	13.8	-	5.42	5.89	5.99
•	9.25	0.354	0.815	0,176	12.5	12.6	± 0.3	± 0.4	± 0.5	ž	0.13	± 0.18	± 0.22
	10.2	0.392	0.880	0,190	13.5	12.4						-	
	12.1	O. lalata	1.02	0.224	15.9	12.2					۹ مه		
	13.7	0.491	1.13	0.246	17.4	11.9						•	
•	15.5	0.542	1.25	0.275	19.5	11.8							
Molecula	r weight	= 220.1 g mol ⁻¹	Chart	weight = 0	.260 5 1	.nch ⁻¹	Chart sc	ale = 8.02 cm	⁻¹ inch ⁻¹		Ce11	path length	= 0.107 m
		į						0					1
		-		-	, ,		*						

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