PROJECTIVE SOLUTION OF DIFFERENTIAL EQUATIONS

ΒY

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> > Остовек, 1972.

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Zoltan Csendes

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BY

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A thesis submitted to the Faculty of Graduate Studies and Research in partial fulfillment of the requirements for the degree of Doctor of Philosophy

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ABSTRACT

A projective method is developed for the solution of differential equations which uses different finite dimensional spaces to approximate the domain and the range spaces of the differential operator. It is shown that rectangular matrices are equivalent to analytic operators in these spaces. Using Lagrangean interpolation polynomials, simple numerical matrix equivalents are evaluated for the two operators $\frac{d}{dx}$ and f(x). From these two types of matrices and a third projection matrix, the discretization of any differential equation is performed by ordinary matrix addition and multiplication. The result of such discretization is a rectangular matrix equation that is solved using the generalized inverse of the matrix to obtain an approximate general solution of the differential equation. Two computer programs using this procedure are presented, one of which solves arbitrary linear ordinary differential equations and the other for the solution of arbitrary linear two-dimensional partial differential equations. The input to these programs is a new computer language for the representation of differential equations called DECL, defined in this thesis. Several applications of the programs to engineering problems are demonstrated. In addition, numerical experiments are performed using the ordinary differential equation program in conjunction with Newton's method to solve nonlinear ordinary differential equations.

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CHAPTER I

INTRODUCTION

In recent years, considerable progress has been made in developing direct numerical methods for the solution of differential equations. To a large extent, this advance is the result of improvements in computer systems, enabling computationally lengthy algorithms to be handled with comparative ease. However, in addition to the procedural refinements which naturally accompany increased usage, there has also been a major change in the type of numerical methods used. A decade ago, the most common method of solving partial differential equations in engineering applications was by means of finite difference formulae; today, for many engineering problems, use of the finite element method is standard.

The effect of the new procedures has been to enlarge the role of projective methods in the numerical solution of differential equations. These methods, of which the Rayleigh-Ritz method, the Galerkin method and the least-squares method occupy a dominant role, are thoroughly treated in standard books by Mikhlin [1], Mikhlin and Smolitskiy [2] and Kantorovich and Krylov [3]. The basic idea in projective methods is to express the solution of a boundary value problem as a sum of known functions with undetermined coefficients. A matrix equation for the unknowns is then derived by requiring the approximate solution to satisfy certain minimality conditions. Provided that the functions and the approximation criterja are chosen judiciously, highly accurate approximate solutions result.

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As originally conceived, classical projective methods were designed to solve a single boundary value problem in a single geometric region. In recent times, however, it has become apparent that the geometric limitations of these methods can be eased by defining new sets of trial functions which are non-zero only over small, standard subregions. The equations for the solution of a boundary value problem in a complicated region can then be easily determined by evaluating the conditions on a projective solution in each of the subregions under specified continuity conditions. The technique of obtaining solutions by this procedure is commonly called the finite element method.

The use of the finite element method was first promoted as a distinct and generally applicable procedure for the solution of structural mechanics problems by Turner, Clough, Martin and Topp [4] in 1956. Its further application in structural mechanics and plate bending analysis was developed in the early 1960's by Clough [5], by Adini and Clough [6] and by Gallagher [7]. In the mid 1960" Argyris made extensive studies of aeronautical structures using the finite element method [8]. At the same time, Zienkiewicz and Cheung [9,10] developed finite element methods for problems occurring in civil engineering. To a large extent, the early popularity of the finite element method in mechanical and civil engineering is due to the considerable work of two groups,

one in aeronautical structures in which the work of Argyris predominates and the other in civil engineering, led by Zienkiewicz.

The application of the finite element method to other engineering disciples quickly followed. In the late 1960's, Oden [11,12] developed finite element methods for nonlinear elasticity and thermomechanical problems. In electrical engineering, the first use of the finite element method was the analysis of electromagnetic wave propagation in homogeneous waveguides by Silvester [13] in 1968 and by Ahmed and Daly [14] in 1969. This was followed in 1970 by the nonlinear finite **element analysis** of electric machine problems by Silvester and Chari [15] and the vector finite element formulation of dielectric loaded waveguides by Csendes and Silvester [16] and by Daly [17].

Applied mathematicians became interested in the finite element method in 1968 with the publication of theoretical papers by Zlamal [18] and by Birkhoff, Schultz and Varga [19]. Further theoretical treatment from the point of view of error bounds for the finite element method was given by Babuska [20] and by Schultz [21].

In the past two to three years, the number of papers appearing on the finite element method has been increasing at an explosive rate. A large part of the engineering applications of these methods is contained in two recent books on the finite element method, one by Zienkiewicz [22] and the other by Oden [23]. However, there is already so much

literature on the subject that, on many topics, these books are limited to mere surveys of the material or, in some cases, to just citing bibliographical references of the original papers.

A similar situation also exists in the theoretical analysis of the finite element method, as evidenced by a recent survey paper by Zlamal [24]. This paper, along with one by Whiteman [25], provides an extensive bibliography of finite element methods. A more reasonably self-contained treatment of the mathematical implications of the finite element method, including rigorous convergence proofs of the method, may be found in a monograph by Varga [26].

Another development of importance to the finite element method, but not directly related to it, is the study of optimal finite difference schemes made by Babuska, Prager and Vitasek [27]. They prove that the consistent Galerkin approximation of a boundary value problem solves the problem of minimizing the norm of the error operator for finite difference methods. Since this is the criterion by which finite element formula are established, it implies that finite element solutions are in general more accurate than the corresponding finite difference solutions.

As originally conceived, the term - the finite element method-had a rather narrow definition. To many early workers, such as Argyris and Zienkiewicz, it meant using basis functions that had local support, were interpolatory at the edges of the elements and resulted in piecewise continuous solutions. However, as in all

successful areas in which many people work, with every new development, the scope of the finite element method has become enlarged. Silvester and Hsieh [28] have used the term when working with an exterior problem in an infinite region to define finite elements of infinite size;Brauchli and Oden [29] employ conjugate approximation functions in the finite element method which are non-zero over the entire solution region; Hazel and Wexler [30] have formulated finite element solutions in terms of non-interpolatory basis functions; and Reichert and Vogt [31] provide such a general definition of the finite element method that it includes almost any method in which pieces are connected together, specifically including the five point finite difference operator. (On the other hand, Babuska, Prager, and Vitasek [27], and more recently Miranker [32], have adopted such a general definition of the finite difference method that it includes virtually all finite element methods!)

In view of this altering usage and common confusion over notation, it is necessary to examine those properties which specifically characterize the finite element method and to define the method accordingly. First of all, a finite element method must be projective. This property implies that there is some kind of correspondence between the analytic space of approximating functions and a finite dimensional vector space. It 5

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therefore distinguishes the finite element method from the more general definition of finite difference methods given by Babuska, Prager and Vitasek [27] in which no analytic function space is defined. Second, in a finite element method, one or more element shapes must be specified over which a standard set of basis functions is defined. Third, the finite element solution of a boundary value problem in a region that is a union of elements must be given as a combination of the solutions of the corresponding elemental boundary value problems. These three conditions correspond closely to the properties by which the finite element method was originally developed and yet provide a flexible definition well-suited to contemporary usage.

Within this general definition of finite element methods there coexist several widely different procedures of developing and using the method, its elements and their basis functions. Most of the early work [33] developed finite element procedures using linear approximation polynomials between the function values at the vertices of triangular or rectangular elements. With these simple approximating functions, the matrix elements of the discretized equation can be evaluated analytically for finite elements of a general shape. The resulting algebraic formulae are then used to generate the elements of a matrix equation for the solution in specific cases.

However, in quest of higher solution accuracies, more sophisticated finite elements have been introduced.

Among the large number of element shapes and basis functions that have been tried, two types of finite elements predominate. One is the isoparametric elements described by Zienkiewicz[22] in which the basis functions are Cartesian products of one dimensional interpolation polynomials in curvilinear coordinates. The others are the high-order symmetric triangular interpolation polynomial elements discovered independently and with entirely different approaches by Silvester [34] and by Irons [35].

An undesirable byproduct of the use of high-order finite elements is an increase in the level of complexity in the procedures used to evaluate the matrix elements. With isoparametric elements, the technique of evaluating general algebraic formulae for the elements of finite element matrices has been abandoned entirely and all calculations are performed from first principles using numerical methods [22]. With the high-order triangular finite elements, it is possible to retain the efficient algebraic approach for generating matrix elements, but the required analytic calculations are lengthy and difficult. Thus, at the present time, the main shortcoming of projective methods is the complicated and expensive calculations needed to discretize each separate type of differential equation.

The primary concern of this thesis is to derive an efficient and accurate algebraic procedure for the discretization of arbitrary differential equations. In order to do so, it will be necessary to devise discrete matrix operators which have analogous combinatorial behavior to the elementary components

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of analytic operators. Once such matrix operators are developed, the matrix equivalent of any differential equation can be determined by simply combining the elementary matrix operators algebraically.

The chief obstacle to forming matrix equivalents of the elementary components of analytic operators is the fact that classical projective methods treat differential equations and boundary conditions as an inseparable unit called a boundary value problem. In these formulations, it is not possible to develop an atomic approach to discretizing differential equations because the entire boundary value problem is fused together. Consequently, in this thesis, the extension of projective methods to solve differential equations which are independent of boundary conditions will be considered. It will be shown that the formulation obtained from this approach yields, for the first time by the direct application of a numerical method, general solutions of differential equations instead of particular ones.

There is, however, in addition to the theoretical limitations mentioned above, a major practical drawback in applying projective methods to solve many common engineering problems. The use of a projective method to solve a differential equation generally requires the development of an extremely complicated and expensive computer program. An indication of the extensive labor sometimes required is provided by the analysis of the Boeing model 747 aircraft wing-body intersection by finite element methods [36]. For this problem, approximately one hundred man-months of effort and about twenty-

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eight hours of CDC 6600 computer time were required to convert existing theory into a working computer program. Clearly, although many engineering problems would benefit from analysis by projective methods, very few of them generate sufficient interest to merit this kind of individual attention.

Consequently, in order for a discretization procedure for arbitrary differential equations to have maximum impact it must be coupled with the development of an efficient computer implementation of the method. By providing a general computer program which performs each of the discretization steps for differential equations automatically, both the difficulty of discretizing arbitrary differential equations and program development costs can be eliminated. For this reason, the ultimate objective of this thesis is to develop general computer procedures by which large classes of differential equations can be solved with a minimum of analysis and data preparation.

In pursuit of this objective, a special purpose computer language has been developed for the symbolic representatio.: of differential equations and two computer programs based on this language - one for the solution of arbitrary linear ordinary differential equations and the other for arbitrary linear two-dimensional partial **differential** equations - have been written. In addition, the linear ordinary differential equation program has been adapted to solve nonlinear ordinary differential equations by means of Newton's method.

Listings of these computer programs will be presented in the thesis along with numerous examples of their operation. Although many of the problems selected can be solved by using alternative methods, it will be demonstrated that none of these alternatives can furnish solutions as profitably as the present programs, in terms of accuracy, computing speed and, especially, units of an engineer's time.

In summary form, the major original contributions of this thesis are:

- (1) The development of a purely algebraic discretization procedure for arbitrary linear ordinary and partial differential equations.
- (2) The definition and evaluation of rectangular matrix equivalents for the three elementary components of differential equations
- (3) The application of matrix generalized inversion to determine approximate general solutions of differential equations.
- (4) The development of a computer language and computer programs for the automatic solution of arbitrary linear ordinary and twodimensional partial differential equations and for the solution of many nonlinear ordinary differential equations by Newton's method.

CHAPTER II

ELEMENTARY OPERATORS

The first step in developing a projective procedure capable of discretizing arbitrary differential equations is to devise a simple method for obtaining their numerical equivalents. In order to do this, it is important to observe that with analytic formulations, any differential equation is simply an algebraic combination of functions with the differential operator. It follows that the construction of a numerical equivalent of a differential equation is most easily accomplished by recognizing the primary numerical constituents of the method used and assembling complicated matrix equations from these components.

In order to proceed, it will be helpful to define the term <u>elementary operators</u> to refer only to those operators which cannot be decomposed further into sums or products of other operators. Although other possibilities exist, for differential equations the simplest and most generally applicable quantities to designate as elementary operators are $\frac{d}{dx}$ and f(x). Any operator that is composed of a product of elementary operators containing a sum of simple operators will be referred to as <u>compound operators</u>. As an example, the compound operator $D = \frac{d^2}{dx^2} + k^2$ consists of two simple operators $\frac{d^2}{dx^2}$ and k^2 . These two simple operators are composed of the elementary operators $\frac{d}{dx}$ and k.

The use of matrices to replace the process of differentiation in finite element discretization has been independently

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advocated by Boisserie [37] and by Silvester [38]. In the first case, the differentiation matrices are square and corresponded to one-dimensional differentiation, with matrix products genereting two and three dimensional finite elements. In the second, rectangular matrices are used to differentiate Newton-Cotes interpolation polynomials over triangular regions. However, neither paper properly develops a complete, consistent approach for obtaining discretized forms of compound operators from these matrix factors.

2.1 DIFFERENTIATION MATRICES

The elementary factor $\frac{d}{dx}$ is the most basic operator in the theory of differential equations. As is well known, its custom-ary definition is

$$Dy = \frac{dy}{dx} = \frac{\lim_{x_2 \to x_1} \frac{y_2 - y_1}{x_2 - x_1}}{(2.1.1)}$$

where y_1 and y_2 are the values of a function y(x) at the points x_1 and x_2 , respectively. Definition (2.1.1) is only valid if the limit exists and is unique for both $x_2 > x_1$ and $x_1 > x_2$. In this equation, the following convention, which will be pursued throughout this thesis, has been introduced: wherever practical, functions are denoted by lower case letters and operators by capital letters; italics are used to monote analytic expressions while Roman type denotes discrete representations.

The standard numerical equivalent of (2.1.1) is to write

a finite difference approximation [39]

$$Dy \simeq \frac{y_2 - y_1}{x_2 - x_1}$$
(2.1.2)

which converges to (2.1.1) as the length of the finite interval L = $x_2 - x_1$ tends to zero. It is apparent that this procedure is equivalent to approximating the function y on the interval I = $[x_1, x_2]$ by the linear polynomial which passes through y_1 and y_2 at the endpoints of the interval. The error introduced by using (2.1.2) instead of (2.1.1) in the interval I is therefore related to the accuracy with which the function y may be approximated by a linear polynomial in I.

It is well known that given a finite number of values, a function can be approximated much more accurately by a few high order polynomials than by many low order ones. Therefore, it is reasonable to seek more accurate finite difference representations of (2.1.1) by approximating y with high order polynomials and then determining the corresponding finite difference formula. Since these formulae are to be in terms of point values, the approximating polynomials should be expressed in terms of an interpolatory basis and may be written as follows

 $y \simeq \sum_{i=0}^{n} y_i l_i^{(n)}(x)$ (2.1.3)

where the { $\mathcal{I}_{\mathcal{L}}^{(n)}(x)$ } are the n'th order Lagrange interpolation polynomials

$$l_{i}^{(n)}(x) = \frac{p_{i}^{(n)}(x)}{p_{i}^{(n)}(x_{i})}$$
(2.1.4)

Here

$$p_{i}^{(n)}(x) = \prod_{\substack{m=0\\m\neq i}}^{n} (x - x_{i})$$
(2.1.5)

The key to forming high order discrete representations of the differential operator (2.1.1) is obtained by recognizing that if an n 'th order polynomial is used to approximate y, the function generated by the process of differentiation will be a polynomial of (n-1) 'st order. Therefore, whenever y is an n 'th order polynomial, the function Dy may be written as

$$Dy = z = \sum_{i=0}^{n-1} \sum_{i=0}^{(n-1)} (x)$$
 (2.1.6)

where again, since point values are desired, the expansion functions are Lagrange interpolation polynomials. Equations (2.1.1), (2.1.3), and (2.1.6) yield

$$\sum_{i=0}^{n-1} z_i \, l_i^{(n_i)} = \sum_{i=0}^n y_i \, \frac{d \, l_i^{(n)}(x)}{dx}$$
(2.1.7)

$$l_{i}^{(n)}\left(x_{j}^{(n)}\right) = \delta_{ij} \qquad j = 0, \dots, n \qquad (2.1.8)$$

So that

$$z_{j} = \sum_{i=0}^{n} \frac{d \, l_{i}^{(n)} \, (x_{j}^{(n-1)})}{dx} \, y_{i} \, j = 0, \dots, n-1 \, (2.1.9)$$

In matrix form this is

z = Dy (2.1.10)

where z is an *n*-component vector containing the numbers z_i , y is an (n+1) component vector containing the numbers y_i and D is the nx(n+1) matrix with elements

$$D_{i+1, j+1} = \frac{d l_i^{(n)} (x_j^{(n-1)})}{dx} \qquad (2.1.11)$$

$$= \begin{cases} \frac{n}{\sum} p \cdot \binom{(n)}{m \neq j} (x_i^{(n-1)}) \\ \frac{m \neq j}{m \neq j} \\ \frac{1}{(x_i^{(n-1)} - x_j^{(n)}) \cdot p_j^{(n)} (x_j^{(n)})} \\ \frac{1}{x_k^{(n)} - x_j^{(n)}} \\ \frac{p k^{(n)} (x_k^{(n)})}{p j^{(n)} (x_j^{(n)})} \\ \frac{1}{p k^{(n)} (x_j^{(n)})} \\ \frac{1}{p k^{($$

$$\begin{bmatrix} n & 1 \\ \Sigma & \frac{1}{x_j - x_m} \\ m \neq j & j & -x_m \end{bmatrix} \quad if \quad x_i^{(n-1)} = x_j^{(n)}$$

As introduced here, in this thesis, non-subscripted lower case letters will indicate vectors having as components the corresponding subscripted discrete values of a function. Similarly, non-subscripted capital letters will refer to matrices which give the discrete representation of an operator.

The simplest interpolation polynomials are the Lagrangian polynomials in which the modal points are equi-spaced in the interval I of length L. Taking

$$x_{i}^{(n-1)} = \frac{L}{n-1} i \qquad i=0,...,n-1$$

$$(2.1.12)$$

$$x_{j}^{(n)} = \frac{L}{n} j \qquad j=0,...,n$$

for $n \ge 2$ the components of the matrix D become

$$D_{i+1,j+1} = \begin{cases} \frac{1}{L} \frac{n(n-1)}{ni-(n-1)j} \sum_{\substack{m \neq 0 \\ m \neq 0 \\ m \neq 0}}^{n} (-1)^{j-m} \frac{(n-m)!m!}{(n-j)!j!} & \text{if } x_{i}^{(n-1)} \neq x_{j}^{(n)} \\ j=0,\ldots,n \\ \frac{n(-1)^{j-k}}{L(k-j)} \frac{(n-k)!k!}{(n-j)!j!} & \text{if } x_{i}^{(n-1)} = x_{k}^{(n)} \neq x_{j}^{(n)} \\ \frac{n}{L} \sum_{\substack{m \neq 0 \\ m \neq 0$$

The differentiation matrices (or D matrices) corresponding to equi-spaced Lagrangean polynomials have been evaluated exactly using Formac symbol manipulator language [40] for polynomial orders ranging from one to nine and are presented in Table 2.1.1. In this table, the following antisymmetry property

$$D_{ij} = -D_{n+2-i, n+1-j}$$
(2.1.14)

is evident. The lowest order case, with n=1, yields for equation (2.1.10).

$$\begin{bmatrix} z_1 \end{bmatrix} = \frac{1}{L} \begin{bmatrix} -1 & 1 \end{bmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix}$$
(2.1.15)

This is equation (2.1. 2) again, only in matrix form. The next higher approximation of (2.1.1) results when y is approximated by a quadratic (n=2) in which case equation (2.1.10) is

$$\begin{pmatrix} z_1 \\ z_2 \end{pmatrix} = \frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix}$$
(2.1.16)

The higher order approximations produce similar rectangular matrix equations.

The elements of the D matrices possess a number of interesting and potentially useful properties. First, it may be noted that the rank of each matrix is one less than the dimension of its domain and that, in fact, the row sums of each and every one of these matrices are zero. This is a necessary consequence of the fact that differentiating a constant, in Table 2.1.1

The first nine D matrices for equi-spaced tagrangean polynomials.

													-0.28350006 03 -0.27321946 02 -0.80104256 01 -0.10704056 02 -0.10701186 02 -0.250510586 01 -0.250510586 01 -0.25050646 02 -0.22040006 03
			-0.1000006 01 - 0.13580246 00 0.8641970E-01 0.83333335 01								-0.10000046 01 0.81907085-01 0.51824578-02 0.51928705-02 0.19282765-02 0.40382245 00 0.21742845 02		0.22870005 03 -0.28730035 02 0.17107285 02 0.17107285 02 -0.1011545 01 -0.111445 01 -0.10498935 02 0.1280005 03
0-10000005 01	-0.1250000E 00 0.550000E 01		0.5333333E 01 -0.9676543E 00 0.4543209E 01 -0.1600000E 02						0.1000000E 01 0.2239966-01 0.2238986-01 0.224600813E-02 0.3333074E 00 0.1814999E 02 0.1814999E 02		0.9142857E 01 -0.1847137E 00 0.18474265E-01 -0.6405771E-02 0.1217923E 02 0.1217923E 02 -0.6400000E 02		
-0-4100000E 01	0.33750006 01		-0.1200005 02 0.54814815 01 -0.54814815 01 0.1200005 02				0.1044140E 00 0.1044140E 00 0.2648806F01 0.2648806F01 0.2899840E 01 0.1470000E 02 0.1470000E 02		-0.0166666 01 0.7160556 00 0.20762076 00 0.6992166-01 0.17886756 00 0.9788646 01 0.4900006 02		-0.3733338 02 0.3754808 00 -0.48544808 00 0.93561386 00 -0.17730618 00 0.97341178 01 -0.22470066 03 0.11200006 03		0.81000000 02 -0.1452676 02 -0.4313146 00 0.4413726-01 0.4113726-01 0.13890946-01 -0.3899094610 00 -0.19010810 00 -0.19125000 02
0.00000000	-0.3375000E 01 0.450000E 01		0.1600006 02 -0.45432096 01 0.9876543E 00 -0.53333336 01		0.100000E 01 -0.118778E 00 0.2343750E-01 0.2343750E-01 0.2065430E 02 0.1141657E 02		-0.1200006 01 -0.80409606 00 0.24652806 00 -1.41512006-01 0.19464955 01 0.3600006 02		0.29399995 02 0.313335 01 0.3134935 01 0.35826825 00 0.858326825 01 0.16937055 02 0.73500005 02 0.73500005 02		0.89599995 02 -0.87526995 01 -0.23726995 01 -0.1393565 01 -0.13499255 02 -0.18641625 02 -0.18641625 03 -0.1893335 03		-0.254071E 02 -0.454859E 00 -0.4454859E 00 -0.41048115E-02 -0.404558E-02 -0.404558E-02 -0.1049117E-01 0.7539217E-01 0.7539217E-01
THIRD ORDER	-0.1250000E 01 -0.1250000E 01	FOURTH ORDER	-0.8641970E-01 -0.8641970E-01 -0.1358024E 00 0.1000000E 01		-0.6250000E 01 0.8422852E 00 0.842208E 00 0.8115728E 01 0.2500000E 02		-0.2230000E 02 0.280239E 01 -0.1228139E 01 -0.1228135E 01 0.688032.E 01 -0.1227024E 02 0.450000E 02		-0.6125000E 02 0.601412E 01 0.2765703E 01 0.8374023E 01 0.8374023E 02 0.1125256 02 0.1125256 02 0.1125256 02 0.8166666 02		-0.1400000 03 0.143455 02 -0.51817455 02 0.10334055 02 -0.10334055 02 -0.10334055 02 -0.14034455 02 0.14004050 03	·	0.10000000 01 -0.75392776-01 -0.14981776-01 0.149815886-02 0.41971846-02 0.41971846-02 0.44348596 00 0.44348596 00 0.2344071E 02
					0.1000000 02 -0.2856455 01 0.5859375 01 -0.8455435 01 -0.8455435 01 0.25000005 02		0.4030000 02 -0.61363195 01 0.78617995 01 -0.78617995 01 -0.78617995 01 -0.4000005 02		0.8166666 02 0.10415256 02 0.1041655 02 0.8374028 01 0.2957035 01 0.6970315 01 0.61250012 01 0.61250012 02		0.149333E 03 0.13404162E 02 0.13404162E 02 0.1340455E 02 0.1493546E 01 0.1493546 01 0.2372494 01 0.89526197 01 0.89520999E 02		-0.1012900E 02 0.789240E 00 0.1612081E 00 0.512681E-01 0.41212645-01 -0.1532645-01 0.411413725-01 0.411413725-01 0.411413725-01 0.411413725-01
			-0.1000001 01 0.3000001 01		-0.250000E 02 0.8455403E 01 -0.5559375E 01 0.2856445E 01 -0.166666E 02		-0.450000E 02 0.1227024E 02 0.1228319E 01 0.1228319E 01 0.225000E 02		-0.7390005 02 0.16937005 02 0.59264755 01 0.59264555 00 0.3919355 01 0.2939995 02		-0.1120006 03 0.27340066 02 0.27341176 01 0.117730616 00 0.18254480 00 0.18254480 00 0.18254480 00 0.18254480 00 0.182312996 01 0.181312995 01		0.4628571E 02 -0.3735512E 01 0.7833556 00 0.7833545 00 0.72538595 00 0.72538595 00 0.72538595 00 0.28852595 02 0.16200005 03 0.16200005 03
C ORDER	0-1000006 01	ID ORDER	0.400000E 01 -0.400000E 01	H ORDER	0.2500006 02 -0.61157236 01 0.32552066 00 -0.84226276 00 -0.84226276 00 0.62300006 01	H ORDER	0.300000E 02 -0.7946495E 01 -0.101200E-01 -0.446280E 00 -0.446280E 00 0.8040960E 01 -0.720000E 01	NTH ORDER	0.4907005 D2 -0.9973646 01 -0.1738575 00 -0.60992165-01 -0.40992165-01 0.279166955 00 -0.79166955 01 0.8166665 01	ITH CRDER	0.4400000E 02 -0.121792E 02 -0.3203715E 02 0.432057715E 02 0.434265E 00 -1804458E 00 -1804137E 00 -142045	TH ORDER	-0.12600005 03 0.12600005 02 0.123017355 01 0.10111445 01 0.10111445 10 0.10315959 02 0.1031595 02 0.1717285 02 0.27270005 03
FIRSI	-0.100000E 01	SECO	-0.30000005 01 0.10000005 01	FIFT	-0.1141667E 02 -0.20654590E 00 -0.2045750E 00 0.1189778E 00 -0.100000E 01	SIXT	-0.1470006 02 -0.289406 00 0.10912006-01 0.10912006-01 0.10941606 00 0.10041606 00 0.10041606 01	SEVE	-0.1814995 02 -0.35330745 00 0.2400875-01 0.8401315-02 -0.2239965-01 0.9293956-01 -0.1000005 01	EIG	-0.2174284E 02 -0.4038324E 00 0.3077522E-01 -0.492870E-02 -0.492870E-02 -0.8248470E-02 -0.8248470E-02 -0.8248470E-02 -0.1020300E 01	NIN	0.2289000E 03 -0.1994264E 02 0.202038E 01 -0.2714559F 01 -0.2714559 02 -0.12719455 02 0.401418E 02 -0.221959F 02 0.2833200E 03

which case all of the y are equal, results in zero. Thus the nullspace of a D matrix corresponds to the nullspace of the operator $\frac{d}{dx}$.

Similarly, by noting that the n'th order D matrix differentiates polynomials of degree less then or equal to n, one surmises from the equation

$$\frac{d}{dx} x^{p} = p x^{p-1}$$
 (2.1.17)

that

$$\sum_{j=0}^{n} \sum_{ij} (x_{j}^{(n)})^{p} = p(x_{i}^{(n-1)})^{p-1} \sum_{p=0,\ldots,n}^{i=0,\ldots,n-1} (2.1.18)$$

This result is valid for general differentiation matrix elements. In particular, with the Lagrangean matrix elements (2.1.13), the following formulae are obtained.

$$\sum_{j=1}^{n} \frac{(-1)^{j+1} n!}{(n-j)! j!} j^{p-1} = 0 \quad p=0,...,n \quad (2.1.19a)$$

$$\begin{array}{cccc} n & n & (-1)^{j-m} & (n-m)!m! \\ \Sigma & \Sigma & & \\ j=o & m\neq j & \hline (ni-(n-1)j) & (n-j)!j! & n & n & \frac{j}{(n-1)^p} & i=1,\ldots,n-2 \\ n & n & n & (n-1)^p & p=0,\ldots,n \end{array}$$

$$\sum_{j=0}^{n-1} \frac{1}{n-j} \left\{ \frac{(-1)^{j-n} n!}{(n-j)! j!} \left(\frac{j}{n} \right)^{p} + 1 \right\} = \frac{p}{n} \qquad p=0,\ldots,n \qquad (2.1.19c)$$

Here, equation (2.1.12) and the fact that $x_i \neq x_j$ unless $i \neq j=0$ or i=n-1 and j=n has been used. These formulae may also be proved by induction.

It is also interesting to note in passing that the differentiation matrices (2.1.10) may be used to provide an alternative definition of differentiation. First note that equation (2.1.1) may be written as

$$Dy = \lim_{L \to 0} \frac{1}{L} \begin{bmatrix} -1 & 1 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \lim_{L \to 0} D^{\binom{1}{1}} y^{\binom{1}{1}}$$
(2.1.20)

where $D^{\binom{1}{}}$ indicates the first order D matrix. One may also write

$$Dy = \lim_{L \to 0} \frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = \lim_{L \to 0} \frac{1}{V} \begin{pmatrix} y_2 \\ y_3 \end{pmatrix} (2.1.21)$$

which will converge faster, but under similar conditions as (2.1.20). Formulae of this type are essentially the same as (2.1.1) since differentiable functions may be approximated by straight lines in the limit as L+O. Consider, however, the sequence of functions

$$z^{(n)} = \sum_{i=0}^{n} z_{i}^{(n)} l_{i}^{(n)}(x)$$
(2.1.22)

defined by the equations

$$z^{(n-1)} = D^{(n)} v^{(n)}$$
(2.1.23)

An alternative definition of the derivative of y is provided by

$$y'(x) = \lim_{n \to \infty} z^{(n)}(x)$$
 (2.1.24)

whenever the sequence $\{z^{(n)}\}$ is uniformly convergent.

2.2 SIMPLE OPERATORS

In this section, methods of developing matrix representations of simple operators will be considered. These will be obtained by taking products of the elementary D matrices of the previous section and some new elementary matrices which correspond to functions.

Fortunately, as each elementary operator $\frac{d}{dx}$ of a product $\frac{d^p}{dx^p}$ acts on a polynomial expression, it generates a polynomial of one lower degree. Therefore, high order differentiation matrices can be obtained by simply multiplying together the succession of matrices which corresponds to the range of polynomials upon which the elementary operators act. For example, to obtain the quadratic equivalent of $\frac{d}{dx}^2$ one forms

$$\mathfrak{D}^{(1)} \mathfrak{D}^{(2)} \mathfrak{y}^{(2)} = \frac{1}{L} \begin{bmatrix} -1 & 1 \end{bmatrix} \frac{1}{L} \begin{bmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{bmatrix} \begin{bmatrix} \mathfrak{y}_1 \\ \mathfrak{y}_2 \\ \mathfrak{y}_3 \end{bmatrix} \quad (2.2.1)$$

$$= \frac{4}{L^2} \begin{bmatrix} 1 & -2 & 1 \end{bmatrix} \begin{bmatrix} \mathfrak{y}_1 \\ \mathfrak{y}_2 \\ \mathfrak{y}_3 \end{bmatrix}$$

This equation is identical to the well known finite difference formula for $\frac{d}{dx^2}$ [39].

In a similar manner, by using the matrix values given in Table 2.1.1, high-order polynomial representations of any differential operator of degree less than ten may be formed. These high-order D matrices have similar properties to low order ones. For example, the dimensionality of the nullspace of the D matrices corresponding to $\frac{d^2}{dx^2}$ is equal to two, with row sums and linearly weighted row sums equal to zero.

It remains to develop matrix representations for simple operators which contain functions. In order to do this, consider the expression

$$z(x) = f(x)y(x)$$
 (2.2.2)

where y(x) is given by an n'th order polynomial

$$y^{(n)}(x) = \sum_{i=0}^{n} y_{i}^{(n)} l_{i}^{(n)}(x)$$
 (2.2.3)

and assume that $\mathfrak{s}(x)$ is to be approximated by an m 'th order polynomial

$$z^{(m)}(x) = \sum_{i=0}^{m} z_{i}^{(m)} l_{i}^{(m)}(x) \qquad (2.2.4)$$

The problem is to determine the matrix operator F which produces the mapping

$$z = Fy$$
 (2.2.5)

corresponding to (2.2.1).

Taking

$$y_i^{(n)} = \delta_{ij}$$
 $j=0,...,n$ (2.2.6)

in (2.2.5) produces the *n* sets of *m* equations

$$z_{i}^{(m)} = F_{ij}$$
 $i=0,...,m$ (2.2.7)

Now if $z_i^{(m)}$ is taken to be the value of the function f(x)y(x) at the point $x_i^{(m)}$, then, using (2.2.6)

$$F_{ij} = z_{i}^{(m)} = f(x_{i}^{(m)}) \sum_{j=0}^{n} y_{j}^{(n)}(x_{i}^{(m)})$$

$$= f(x_{i}^{(m)}) l_{j}^{(n)}(x_{i}^{(m)})$$
(2.2.8)

Notice, first of all, that if f(x)=1 and $n \neq m$ then (2.2.7) defines an operator which maps a function into a higher or a lower order space. The raising operator with m>n is exact and provides a convenient method for imbedding a low order polynomial into z

higher order space. Table 2.2.1 contains the values of this matrix with m = n+1 for polynomial orders ranging from zero to eight.

The most useful form of (2.2.7) arises, however, when n=m since then $l_i^{(n)}(x_i^{(m)}) = \delta_{ij}$ and

$$F_{ij} = f(x_i^{(m)})\delta_{ij}$$
(2.2.9)

Thus, in order to find the product of a function $y^{(n)}$ expressed in terms of *n*'th order interpolation polynomials and an arbitrary function f(x), it is only necessary to multiply the vector of coefficients of y by a diagonal matrix containing the values of the function f(x) at the interpolation nodes. This procedure has the same effect as weigh ting the vector y by the nodal values of f(x).

2.3 COMPOUND OPERATORS

Before the matrix equivalent of an arbitrary differential operator can be determined, the effect of additions on the discretization procedure of the previous sections must be determined. That such operations may not be performed in a straightforward manner may be seen by considering the compound operator $D = \frac{dy}{dx} - y$. If y is approximated by a quadratic, then the first term in this expression is a linear polynomial having a two component vector representation while the second term is a quadratic having a three component representation. Clearly, it is not possible to add these two vectors.

Table 2.2.1

The first nine FU matrices for equi-spaced Lagrangean polynomials.

1 ST ORDER			4 TH ORDER MÁT	RIX				
0.1000000E 01 0.1000000E 01 2 ND DRDER			0.1000000E 01 0.1171875E 00 -0.6250000E-01 0.1049805E 00 0.5625000E 00	0+0 0+1054688E 01 0-5625000E 00 -0-2109375E 00 0+0	0.0 -0.2109375E 00 0.5625000E 00 0.1054688E 01 0.0	0.5625000E 00 0.1049805E 00 -0.625000E-01 0.1171875E 00 0.100000E 01		
0.10000005 01 C.50000005 00 D.0	0.0 0.5000000E 00		5 TH ORDER MAT	RIX				
3 TH DRDER			0.1000000E 01 0.7040000E-01 -0.3360000E-01 0.2756096E-01	0.0 0.1126400E 01 0.3584000E 00	0.0 -0.2816000E 00 0.8063999E 00	0.0 0.1023999E 00 -0.1536000E 00	-0.1536000E 00 -0.2841344E-01 0.2756096E-01	
0.1000000E 01 0.2222222E 00 0.111111E 00	0.0 0.88888888 c0 0.88888888 c0	0.0 -0.1111111E 00	-0.2841344E-01 -0.1536000E 00	0.1023999E 00 0.0	-0.2816000E 00 0.0	0.3584000E 00 0.1126400E 01 0.0	-0.3360000E-01 0.7040000E-01 0.1000000E 01	
5.0	0.0	9.1000000E 01	6 TH ORDER MAT	RIX				
7 TH ORDER HAT	AIX		0.1000000E 01 0.4632309E=01 -0.1920439E=01 0.1171875E=01 -0.3249779E=01 -0.9322089E=01 0.5733449E=01	0.0 0.1158076E 01 0.2400548E 00 0.9765625E-01 0.5733449E-01 -0.3380900E-01 0.5859375E 00	0.0 -0.3308792E 00 0.9602194E 00 0.5859375E 00 -0.2400548E 00 0.1781657E 00 0.0	0.0 0.1781657E 00 -0.2400548E 00 0.5859375E 00 0.9602194E 00 -0.3308792E 00 0.0	0.5859375E 00 -0.3380900E=01 0.5733449E=01 -0.9765625E=01 0.2400548E 00 0.1158076E 01 0.0	0.5733449E-01 -0.9322089E-01 -0.3249779E-01 0.1171875E-01 -0.1920439E-01 0.4632309E-01 0.1000000E 01
0.1000000E 01 0.3253746E=01 0.1172980E=01 0.6357893E=02 -0.7442292E=02 0.1173713E=01 0.3948175E=01 =0.3473582E=01	0.0 0.1171349E 01 0.1689092E 00 -0.6242297E-01 0.3929992E-01 -0.3470582E-01 0.3480585E-01 -0.1716632E 00	0.0 -0.3660464E 00 0.1055682E 01 0.4291579E 00 -0.1716632E 00 0.1319603E 00 -0.1331078E 00 0.0	0.0 0.2602997E 00 0.3127948E 00 0.7629474E 00 0.7629474E 00 -0.3127948E 00 0.2602997E 00 0.0	0.0 -0.1331078E 00 0.1319603E 00 -0.1716632E 00 0.4291579E 00 0.1055682E 01 -0.3660464E 00 0.0	-0.1716632E 00 0.3480585E-01 -0.3470582E-01 0.3929992E-01 -0.6242297E-01 0.1689092E 00 0.1171349E 01 0.0	-0.3470582E-01 0.3948175E-01 0.1173713E-01 -0.7442292E-02 0.6357893E-02 -0.1172980E-01 0.3253746E-01 0.1000000E 01		
8 TH CROER HAT	RIX .							
0.1000000000000 0.2399191E-01 -0.7587433E-02 0.3643781E-02 -0.24414006-02 0.9497318E-02 0.188-332E-01 -0.2165224E-01 0.6205067E-02	0.0 0.1175603E 01 0.1239281E 00 -0.4120275E-01 0.2392578E-01 -0.1453304E-01 0.6205067E-02 -0.1689447E 00 0.8075355E-01	0.0 -0.3918679E 00 0.1115353E 01 0.3213814E 00 -0.1196289E 00 0.8675355E-01 -0.9033471E-01 0.1212236E 00 0.5981445E 00	0.0 0.3457658E 00 -0.3717842E 00 0.8927263E 00 -0.2981445E 00 -0.2434708E 00 0.265468E 00 -0.2351207E 00 0.0	0.0 -0.2351207E 00 0.2065468E 00 -0.2434708E 00 0.5981445E 00 0.8927263E 00 -0.3717842E 00 0.3457658E 00 0.0	0.5981445E 00 0.1212236E 00 =0.9033471E=01 0.8675355E=01 =0.1190289E 00 0.3213814E 00 0.1115353E 01 =0.3918679E 00 0.0	0.8675355E-01 -0.1689447E 00 0.6205067E-02 -0.1453304E-01 0.2392578E-01 -0.4120275E-01 0.1239281E 00 0.1175603E 01 0.0	0.6205067E-02 -0.2166624E-01 0.1884082E-01 0.9497818E-02 -0.2441406E-02 0.3643781E-02 -0.7587433E-02 0.2399191E-01 0.1000000E 01	
9 TH DROER MAT	RIX							
0.10000000F 01 0.1836341E=01 =0.5147040E=02 0.2201561E=02 0.2255E=02 C.2242255E=02 =0.3848400E=02 =0.9161562E=02 0.1249007E=01 =0.3695211E=02	0.0 0.1175258E 01 0.9411728E-01 0.2817998E-01 0.1464620E-01 0.921258E-02 0.921258E-02 0.8695211F-02 0.6372003E-01 -0.4971430E-01	0.0 -0.4113403E 00 0.1152937E 01 0.2465748E 00 -0.3421564E=01 0.5383047E=01 0.4971430E=01 0.4161451E=01 -0.9275264E=01 =0.1813876E 00	0.0 0.4329898E 00 0.492498E 00 0.9862995F 00 0.4716078E 00 -0.1813376E 00 0.1408999F 00 -0.1590258E 00 0.2223461E 00 0.0	0.0 -0.3672681E 00 0.2882342E 00 -0.3082185E 00 0.7368872E 00 -0.3682185E 00 0.2882342E 00 -0.3672681E 00 0.0	0.0 0,2223461E 00 -0,1590258E 00 -0,1813876E 00 -0,1813876E 00 0.4716078E 00 0.9862995E 00 -0.4192498E 00 0.4329898E 00 0.0	-0.1813876E n0 -0.9275264E-n1 0.6161451E-n1 -0.4971430E-n1 -0.5383947E-n1 -0.6421564E-n1 0.2465748E n0 0.1152937E n1 -0.4113403E n0 0.0	=0.4971430E=01 0.8372003E=01 =0.8695211E=02 0.9212058E=02 =0.1001270E=01 0.1464620E=01 =0.2817998E=01 0.9411728E=01 0.1172558E 01 0.0	-0.8695211E-02 0.1249007F-01 -0.9161662F-02 -0.38424005-02 0.2042245F-02 -0.1315270E-02 C.22C1561F-02 0.1836342E-01 0.1836342E-01

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A way out of this dilemma is to form two separate approximations of y, one in the quadratic space and another in the linear space

$$y \simeq y^{\binom{2}{2}} = \sum_{i=0}^{2} y_{i}^{\binom{2}{2}} l_{i}^{\binom{2}{2}} (x)$$
(2.3.1a)

$$y \simeq y^{\binom{1}{1}} = \sum_{i=0}^{1} y_{i}^{\binom{1}{1}} l_{i}^{\binom{1}{1}} (x)$$
 (2.3.1b)

Then, using (2.3.1a) for the first term of D and (2.3.1b) for the second, this yields, just as before

$$D = \frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_1^{\binom{2}{1}} \\ y_2^{\binom{2}{2}} \\ y_3^{\binom{2}{2}} \end{pmatrix} - \begin{pmatrix} y_1^{\binom{1}{1}} \\ y_2^{\binom{1}{2}} \\ y_2^{\binom{2}{2}} \end{pmatrix}$$
(2.3.2)

Now, although the terms of the matrix expression (2.3.2) are consistent in that they are both two component vectors, the expression can be further simplified by defining a relation-ship between the two approximations (2.3.1a) and (2.3.1b). One such relationship is provided by (2.2.5) with m=2 and n=0

$$\begin{pmatrix} y_1^{(1)} \\ y_2^{(1)} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} y_1^{(2)} \\ y_2^{(2)} \\ y_3^{(2)} \end{pmatrix}$$
(2.3.3)

However, since this relationship is obtained by matching the function values at the endpoints of the interval, it provides a very crude method of approximating a quadratic by a linear polynomial.

A much better approximation is obtained by determining the best linear approximation to a quadratic in a least squares (L_2) sense in the interval I. Under this condition, the norm of the difference of the two functions $y^{(1)}$ and $y^{(2)}$

$$Q = \int_{I} |y^{(2)} - y^{(1)}|^2 dx \qquad (2.3.4)$$

must be a minimum. Using (2.3.1), this means that

$$\frac{\partial Q}{\partial y_{i}^{(1)}} = 0 = -\sum_{j=0}^{2} y_{j}^{(2)} \int l_{i}^{(1)} l_{j}^{(2)} dx + \sum_{j=0}^{1} y_{j}^{(1)} \int l_{i}^{(1)} l_{j}^{(1)} dx$$

$$i=0,1$$
(2.3.5)

Written in matrix form, this equation is

$$S y^{(2)} = T y^{(1)}$$
 (2.3.6)

where the elements of S and T are

$$S_{ij} = \int l_{i}^{(1)} l_{j}^{(2)} dx \qquad (2.3.7)$$

$$T_{ij} = \int l_{i}^{(1)} l_{j}^{(1)} dx$$

$$I = \int l_{i}^{(1)} l_{j}^{(1)} dx$$

Since the $\{l_i^{(1)}\}$ are linearly independent, the metric T is non-singular and may be inverted to yield

$$y^{(1)} = Py^{(2)}$$
 (2.3.8)

polynomials.	
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	-0.1523812E 0 0.18694895-0 -0.43033575-0 0.4470199E 00					0.2195871E- 0.5165751- 0.3193801E- 0.333801E- 0.473080E- 0.4730801312E- 0.4031312E- 0.9780403E
	0.6093240E 00 -0.1241621E 00 0.6659608E 00 0.6095246E 00				-0.3131356-01 -0.72741666-02 0.6301545-02 -0.6301545-02 -0.438974-02 0.4389746-02 0.105430166-01 0.96766416 00	-0.1976236E 00 -0.4695141E-01 0.4537535E-01 -0.6181378E-02 -0.87937231E-01 -0.877928E-01 -0.877928E-01 -0.877928E-01 0.1976281E 00 0.1976281E 00
-0.3333331E 00 0.666663E 00	-0.9142890E 00 0.4825394E 00 0.4822394E 00 0.4822394E 00			0.4761118E-01 0.8355518E-02 0.83555518E-02 0.25457053E-02 0.123733E-02 0.123733E-01 0.12373367E-01 0.3523977E 00	0.2586472E 00 0.2586472E 00 0.4833462E-01 0.4833462E-01 0.400313952E 01 0.40031355E 01 0.8503652E 00 0.2586508E 00	0.7905053E 00 0.17905053E 00 0.1111424E 00 0.1111424E 00 0.1111424E 00 0.20050518E 01 0.2005051E 00 0.7905123E 00
SECOND ORDER 0.6666658E 00 0.6666638E 00	FOURTH ORDER 0.60995246E 00 0.60595266E 00 -0.1541622E 00 0.60995240E 00		-0.7012907E-01 -0.6661860F-02 0.918529E-03 0.1261948E-01 -0.1261948E-01 0.3298682E 00	-0.332774E 00 -0.540358EE-01 0.3540195E-01 0.35384E-01 0.332584E-01 0.438251E 00 0.438251E 00 0.3332787E 00	-0.9032674E 00 -0.15016876 00 -0.1501689E 00 -0.131689E 00 -0.131944E 00 -0.131944E 00 -0.1924124E 00 0.1924124E 00 -0.9052740E 00	-0.1844520E 01 -0.39583945 00 0.39583945 00 -0.23452095 00 -0.25452095 00 -0.2545205 00 0.47370625-01 0.80372335-01 0.80372335-01
0.666663E 00 -0.3333331E 00	0.8476199E 00 0.430397E=01 0.1809687E=01 -0.1523812E 00	0.1033394E 00 0.11149625E-03 0.1114972E-01 0.3740272E-01 0.4906596E 00	0.4207786E N0 0.3370728E-01 0.4802225E-02 0.407788E 00 0.8078514E 00 0.4207785E 00 0.4207785E 00	0.9998300E 00 0.19998300E 00 0.191915E-01 -0.1910222E 00 0.2143154 00 0.2143154 00 0.243154E 00	0.18105396 01 0.3133156 00 0.3133156 00 0.89337565 00 0.89337565 00 0.12213555 00 0.12213555 00 0.12213555 00 0.121056775 00 0.18105446 01	0.2765784E 01 0.2014911E 00 0.26719346F 00 0.2671935E 00 0.8050532E 00 0.33310955E 01 0.337276784E 01 -0.3772766784E 01
		-0.5166997E 00 0.12664957E 00 0.126595E-01 0.5165998E 00 0.5166998E 00	-0.1051945 01 -0.49707805-01 -0.49707805-01 0.7248596 00 0.27861206 -0.1051947E 01	-0.1666386E 01 -0.179187E 00 -0.179167E 00 0.5781441E 00 0.2710918E 00 0.1666388E 01 0.1666388E 01	-0.2263177E 01 -0.3111444E 00 0.370431335E 00 0.3704381E 00 0.241335E 00 0.3111444E 00 -0.3111444E 00 -0.2263177E 01	-0.2765784E 01 -0.3772148E 00 0.393055E 00 0.8805598E-01 0.6055232E 00 0.501591155E 00 0.319636E 00 0.319636E 00 0.319634E 01 0.2765784E 01
	0.2230000E 00 -0.8290000E-01 0.7730000E 00	0.103397E 01 -0.100574E 00 0.509294FE 00 0.5271229E 00 -0.1033399E 01	0.1402993E 01 0.1402595E 01 0.4400875E 00 0.4400875E 00 0.1402593E 01 0.1402593E 01	0.1666388E 01 0.520188E-01 0.5210915E 00 0.5381441E 00 0.5381441E 00 0.1779187E 00 -0.1666386E 01	0.1810544E 01 0.1026579E 00 0.1221551E 00 0.172551E 00 0.8833556E 00 0.218105395E 00 0.18105395E 01 0.18105395E 01	0.1844526E 01 0.1844526E 01 0.457628E-01 0.4976048E 00 0.1345847E 00 -0.134581946E 00 0.3558394E 00 -0.3831695E 00 -0.1844520E 01
	-0.6749998E 00 0.5624994E 00 0.5749994E 00 0.6749994E 00	-0.1033399E-01 0.33771259E 00 0.5859367E 00 -0.1033347E 01 0.1033397E 01	-0.1051947E 01 0.2796120E 00 0.2796120E 00 0.2740395E-01 -0.49707802-01 -0.1091945E 01	-0.9998332E 00 0.2143134E 00 0.621431563E 00 0.62131563E 00 0.1196292E 00 0.14164392 00 0.14164393E 00 0.9998300E 00	-0.993740E 00 0.1924134E 00 0.8890578E 00 0.387941E 00 -0.1351689E 00 -0.1581689E 00 -0.1581689E 00 -0.853761E 00 -0.9952674E 00	-0.7905123E 00 0.2005551E 00 0.9080182E 00 -0.21574E 00 0.3460805F-01 0.113424E 00 0.1780190E 00 0.7905033E 00
FIRST ORDER 0.5000006 00	THIRD ORDER 0.674994E 00 0.5624996E 00 -0.6749998E 00	F1FTH ORDER 0.5166998E 00 0.7315500 0.7315501 0.176453F01 0.176453F01	SIXTH ORDER 0.42077855 00 0.80785165 00 0.8078516 00 0.4625256-01 0.42077865 00 0.42077865 00	SEVENTH ORDER 0.3332787E 00 0.8392511E 00 0.130268E 00 0.39401958E-01 0.39401958-01 0.39401958-01 0.39401958-01 0.39401951-01 -0.49401958-01	EIGHTH ORDER 0.2396508E 00 0.8303492E 00 0.1317535-01 0.1315335-01 0.4833452-01 0.5543250E-01 0.5543250E-01 0.2540672E 00	NINTH ORDER 0.1741E 00 0.1741E 00 0.8771196 00 0.8779639E-01 0.4873959E-01 0.487375950 0.4893141E-01 -0.4895141E-01 -0.4895141E-01
0.500000£ 00	0,7730000E 00 -0.6230300E 00 0.2233000E 00	0.8960596E 00 0.37652601 0.1118785601 0.11495026603 0.1033398E 00	0.9294382E 00 -0.3247972E-01 0.125345E-01 0.9153245=03 -0.9011400E-02 -0.9011400E-02	0.952873F 00 -0.2572467E-01 0.1273372F-01 0.1231E-02 -0.2647335F-02 0.839951E-02 0.4761118E-01	0.9676681E 00 -0.2142316-01 -0.105-3398-01 -0.4438974E-02 -0.4438974E-02 -0.330154311E-02 -0.727.1666-02	0.9783403E 00 -0.107706E-01 0.4013125-02 -0.4734592-02 0.3373481-03 0.373481-02 0.374201 E-02 0.2730487502 0.2730487502 0.2730471

1000

2.1

where

 $P = T^{-1}S$ (2.3.9)

Equation (2.3.8) gives the coefficients of the approximation (2.3.1a) and has the effect of projecting the second order polynomial $y^{(2)}$ into the best linear polynomial in the I_{2} sense. Using (2.3.8), equation (2.3.2) may be written as

$$Dy = \frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_{1}^{(2)} \\ y_{2}^{(2)} \\ y_{3}^{(2)} \end{pmatrix} - \frac{1}{3} \begin{pmatrix} 2 & 2 & -1 \\ -1 & 2 & 2 \end{pmatrix} \begin{pmatrix} y_{1}^{(2)} \\ y_{2}^{(2)} \\ y_{3}^{(2)} \end{pmatrix} (2.3.10)$$

The projection matrix (or P matrix) given in equation (2.3.9) may also be defined between polynomial spaces of differing degree and have been evaluated using the Formac language for all polynomial orders ranging from one to nine. These values are presented in Table 2.3.1. The P matrices have an analogous topological behavior to the differentiation matrices, although their physical interpretation is different. They are symmetric about their centroids

$$P_{ij} = P_{n+2-i, n+1-j}$$
(2.3.11)

and not anti-symmetric, as were the D matrices. Their rank is again one less than the dimension of their domain space, however, in this case, the nullvectors of P are symmetric about y = 0 in an L_2 sense.

As was the case with differentiation matrices, projection matrices between spaces differing by more than one order can be determined by ordinary matrix multiplication of the corresponding
elementary P matrices. For example, the best possible zeroth order polynomial approximation of the quadratic $y^{(2)}$ is given by

$$\begin{pmatrix} y_{1}^{(o)} \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 \end{pmatrix} \frac{1}{3} \begin{pmatrix} 2 & 2 & -1 \\ -1 & 2 & 2 \end{pmatrix} \begin{pmatrix} y_{1}^{(2)} \\ y_{2}^{(2)} \\ y_{3}^{(2)} \end{pmatrix}$$

$$= \frac{1}{6} \begin{pmatrix} 1 & 4 & 1 \end{pmatrix} \begin{pmatrix} y_{1}^{(2)} \\ y_{2}^{(2)} \\ y_{3}^{(2)} \end{pmatrix}$$

$$(2.3.12)$$

In a point matching sense, the equation corresponding to (2.3.12) would be from (2.2.5)

$$\begin{pmatrix} y_{1}^{(2)} \end{pmatrix} = \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} \begin{pmatrix} y_{1}^{(2)} \\ y_{2}^{(2)} \\ y_{3}^{(2)} \end{pmatrix}$$
 (2.3.13)

where the interpolation "node" of the zeroth order polynomial is taken in the midpoint of the interval. Therefore, for the differential operator

$$D = \frac{d^2 y}{dx^2} + y$$
 (2.3.14)

the following two finite difference representations are obtained

$$D = \frac{4}{L^2} \begin{bmatrix} y_{i-1} - 2y_i + y_{i+1} \end{bmatrix} + \frac{1}{6} \begin{bmatrix} y_{i-1} + 4y_i + y_{i+1} \end{bmatrix}$$
(2.3.15a)

and

$$D = \frac{4}{L^2} \left[y_{i-1} - 2y_i + y_{i+1} \right] + y_i \qquad (2.3.15b)$$

Equation (2.3.15b) is the standard finite difference formula for (2.3.14) while (2.3.15a) is the "asymptotically optimal"

finite difference formula for (2.3.14) derived by Miranker [32] using a more circuitous theory. The only difference between the two is in the type of approximation made between the polynomial spaces. Miranker compared the formulae (2.3.15) computationally [32] and found that the L_2 projective approximation was as much as 10^4 times more accurate in the Euclidean error norm of point values than the point matching approximation.

CHAPTER III

PROJECTIVE APROXIMATION

In the preceding chapter, the elementary matrix concept was introduced by using heuristic arguments. In this chapter, elementary matrices will be developed more formally, using the theory of projective approximation. The procedures outlined in Chapter 2 will thus be established rigorously here, resulting in a foolproof algorithm for the solution of differential equations.

Somewhat surprisingly, very little work has been done in the area of projective approximations of pure differential equations, considering the wealth of similar material on boundary value problems. The only reference available in the literature on this subject appears to be a paper by Locker [41], who considers the problem of determining the least-squares solution of an n'th order differential equation to which k boundary.conditions have been added, with k not necessarily equal to n. Unfortunately, the theory developed in this paper has several limitations to practical application of the method. It assumes that the homogeneous solutions of the differential equation are known, a condition very rarely occurring in practice, and concentrates on determining only the inhomogeneous solution. Moreover, the procedure is theoretically proved to converge to a unique inhomogeneous solution, but the paper does not specify how the com-

putation should be constructed numerically.

3.1 GALERKIN'S METHOD

In this section, the basic mathematical equations for the projective solution of differential equations will be given. These equations are related to the classically used Galerkin's method, except that the procedure will be applied to differential equations independently of boundary conditions. The main result of this section will be to determine a unique range space for differential operators which result in approximate general solutions of differential equations.

In a projective method, the objective of the analysis is to determine an approximate solution of the equation

$$Dy = f \tag{3.1.1}$$

by replacing the function y, which belongs to a real-valued Hilbert space Σ with norm $|| f || = \langle f | f \rangle^{\frac{1}{2}}$, with a function $y^{(f)}$ in a finite dimensional subspace $\Sigma^{(f)}$ of Σ . In the present case, D is taken to be a linear differential operator defined over a geometric region I and f is a given function of Σ which lies in the range of D. In the following analysis, it will be helpful to decompose the differential operator Dinto a sum of simple operators D_{c}

$$D = \sum_{s=1}^{S} D_{s}$$

(3.1.2)

and to consider the range spaces $\Sigma^{(s)}$ generated by each of these operators separately. These range spaces are defined by the equations

$$\Sigma^{(s)} = \{ z : z = D_{sf} y^{(f)}, y^{(f)} \in \Sigma^{(f)} \}$$
(3.1.3)

Here and in the following, subscripts on an operator indicate that

Range
$$(A_{ij}) = \Sigma^{(i)}$$
 (3.1.4a)
Domain $(A_{ij}) = \Sigma^{(j)}$ (3.1.4b)

$$Domain(A_i) = \Sigma$$
(3.1.4c)

Now, it is necessary to provide a general definition of the projective mapping introducted in section 2.3. This is done as follows. A projection operator P_{ij} is defined to be an operator which maps any element of a space $\Sigma^{(j)}$ into the closest (in the norm sense) element of $\Sigma^{(i)}$

$$P_{ij} = \{A_{ij} : y^{(i)} = A_{ij} y^{(j)}, y^{(j)} \in \Sigma^{(j)}$$

$$(3.1.5)$$

$$=> || y^{(i)} - y^{(j)} || = \min[||x||^{(i)} - y^{(j)} ||]$$

$$= x^{(i)} \in \Sigma^{(i)}$$

Projection operators are well known in the theory of linear spaces [42,43] and have a number of interesting properties. The most important of these are that they are linear

$$P_{ij}(y^{(j)} + z^{(j)}) = P_{ij}y^{(j)} + P_{ij}z^{(j)}$$
 (3.1.6)

and that for three nested subspaces $\Sigma^{(i)}$, $\Sigma^{(j)}$, $\Sigma^{(k)}$ with $\Sigma^{(i)} \subseteq \Sigma^{(j)} \subseteq \Sigma^{(k)}$

$$P_{ik} = P_{ij} P_{jk}$$
(3.1.7)

Note, however, that in the event that the space $\Sigma^{(i)}$ is larger than the space $\Sigma^{(j)}$, P_{ij} is not equal to the identity operator, as is commonly assumed, since the projective embedding of a smaller space into a larger one is not unique.

The underlying idea of many projective methods is to replace the problem of solving equation (3.1.1) by the task of solving the equation

$$\sum_{s=1}^{S} P_{0s} D_{sf} y^{(f)} = f^{(0)}$$
(3.1.8)

where $f^{(0)} = P_0 f$ and $\Sigma^{(0)}$ is some finite-dimensional subspace of Σ that is called the range space of the approximation. If $\Sigma^{(0)}$ is taken to be equal to $\Sigma^{(f)}$ then (3.1.8) yields the Bubnov-Galerkin method, while the Galerkin-Petrov method results if $\Sigma^{(0)}$ does not coincide with $\Sigma^{(f)}$ [2].

Unfortunately, however, the usual choices for $\Sigma^{(0)}$ that appear in the literature have a major flaw: they do not preserve the dimension of the nullspace of the differential operator *D*. As a result, with these projective methods, it is not possible to obtain the general solution of the differential equation (3.1.1), but only particular solutions to specified boundary value problems. In order to produce finite-dimensional equations of the form (3.1.8) which behave in a manner analogous to that of (3.1.1), a space $\Sigma^{(0)}$ must be found such that the operator $\sum_{s=1}^{S} P_{0s} D_{sf}$ has approximately the same nullspace as the original operator *D*. For ordinary differential equations, this implies that the range space $\Sigma^{(0)}$

should be of dimension

$$Dim(\Sigma^{(0)}) = Dim(\Sigma^{(f)}) - Nullity(D) \qquad (3.1.9)$$

where the nullity of a compound ordinary differential operator is given by the highest power of differentiation among its simple components. Extension of this concept to partial differential equations will be considered in section 3.5.

A comparison can be made between the method contained in equation (3.1.8) and the asymptotically optimal finite difference formulations of Babuska, Prager and Vitasek [27] and of Miranker [32]. In their theory, the criterion for establishing the asymptotic optimality of a finite difference formula is to determine whether or not it minimizes the norm of the residual of the Galerkin process:

$$y_{op}^{(f)} = \{ y^{(f)} : || Dy^{(f)} - f || = \min \| || Dz^{(f)} - f || \}$$

$$z^{(f)}_{\epsilon \Sigma}(f) = \{ (3.1.10) \}$$

In the method developed in this thesis, solutions are obtained by saying that functions in the space $\Sigma^{(0)}$ are the best possible approximations of functions in the spaces $\Sigma^{(s)}$. When the differential equation (3.1.1) has only two terms, as in equation (2.3.14), the two approaches yield identical results; if the equation has more than two terms, they do not. The advantage of the present formulation is the ease with which numerical matrix equivalents can be generated for arbitrary differential equations, compared to the difficulty of evaluating asymptotically optimal finite difference formula.

3.2 INTRODUCING BASIS FUNCTIONS

The first step in discretizing equation (3.1.8) is to introduce a set of basis functions $\{b_i^{(f)}\}\$ in the space $\Sigma^{(f)}$ and a basis $\{b_i^{(o)}\}\$ in $\Sigma^{(o)}$. Then

$$y^{(f)} = \sum_{i=1}^{n} y_i b_i^{(f)}$$
(3.2.1)

$$f^{(o)} = \sum_{i=1}^{m} f_i b_i^{(o)}$$
(3.2.2)

and (3.1.8) becomes

$$\sum_{i=1}^{n} y_{i} \sum_{s=1}^{s} p_{os} D_{sf} \frac{b_{i}^{(f)}}{i} = \sum_{i=1}^{m} f_{i} \frac{b_{i}^{(o)}}{i}$$
(3.2.3)

Taking the scalar product of this equation with each function in a conjugate basis set { $b_{(0)}^{i}$ } in $\Sigma^{(0)}$ results in the equation

$$S_{\Sigma} A_{S} y = Bf \qquad (3.2.4)$$

Here, as before, y and f are vectors containing the coefficients of the expansion (3.2.1) and (3.2.2), respectively, the A_s are $m \times n$ matrices with elements

$$A_{ij} = \left\langle b_{(o)}^{i} \middle| P_{\text{os}} b_{sf}^{(f)} \right\rangle \qquad (3.2.5)$$

and B is an $m \times m$ matrix with elements

$$B_{ij} = \left\langle b_{(o)} \middle| b_{j}^{(o)} \right\rangle$$
(3.2.6)

Hence, in a finite dimensional subspace, a differential equation is equivalent to a rectangular matrix equation with analytic operators replaced by rectangular matrices and functions by vectors. The correspondence between the vectors and the continuous functions is provided by equation (3.2.1)

Consider now the problem of determining the function $f^{(o)}$, the best approximant to the function f in the space $\Sigma^{(o)}$. This problem is treated in linear analysis and the presentation below draws from references [42,43,23]. Using (3.1.5) the definition of $f^{(0)}$ given by (3.2.2) implies that

$$||f-f^{(o)}||^{2} = \langle f - f^{(o)} | f - f^{(o)} \rangle$$

$$= \langle f - \sum_{i=1}^{m} f_{i} b_{i}^{(o)} | f - \sum_{i=1}^{m} f_{i} b_{i}^{(o)} \rangle$$

$$= \langle f | f \rangle - 2 \sum_{i=1}^{m} f_{i} \langle b_{i}^{(o)} | f \rangle + \sum_{i,j=1}^{m} f_{i} f_{j} \langle b_{i}^{(o)} | b_{j}^{(o)} \rangle$$

$$= \min_{f^{(o)} \in \mathbf{E}^{(o)}}$$

$$(3.2.7)$$

Differentiating this with respect to f_k yields

$$\langle b_k^{(o)} | f \rangle = \sum_{i=1}^m \mathsf{T}_{ki} \mathsf{f}_i \qquad k=1,\ldots,m \qquad (3.2.8)$$

where

$$T_{ij} = \langle b_i^{(0)} | b_j^{(0)} \rangle$$
 (3.2.9)

Therefore

$$\mathbf{f}_{i} = \sum_{j=1}^{m'} \mathsf{T}^{ij} \left\langle b_{j}^{(o)} | f \right\rangle = \left\langle h_{(o)}^{i} | f \right\rangle \qquad (3.2.10)$$

where

$$h_{(o)}^{i} = \sum_{j=1}^{m} T^{ij} b_{j}^{(o)}$$
(3.2.11)

and T ij are the elements of the inverse of T

$$\sum_{j=1}^{m} \mathsf{T}^{ij} \mathsf{T}_{jk} = \delta_{ik}$$
(3.2.12)

Using this result, the function $f^{(o)}$ becomes

)

$$f^{(o)} = P_{o}f = \sum_{ij} b_{i}^{(o)} T^{ij} \langle b_{j}^{(o)} | f \rangle$$
 (3.2.13)

Consequently, the operator P_0 may be written as

$$P_{o} = \sum_{ij} b_{i}^{(o)} T^{ij} \langle b_{j}^{(o)} |$$
(3.2.14)

where the open inner product is understood to be completed with whatever function \mathcal{P}_0 acts on. Although (3.2.14) has been derived specifically between the spaces Σ and $\Sigma^{(o)}$, it is clear that the projection operator P_{ij} between any two spaces $\Sigma^{(i)}$ and $\Sigma^{(j)}$ will have the same form, with the basis $\{b^{(o)}\}$ replaced by $\{b^{(i)}\}$.

The projective property of the operator P_{o} is demonstrated in the following manner. If a function f is decomposed into two functions, one in $\Sigma^{(o)}$ and the other in $\Sigma^{(x)} = \Sigma \setminus \Sigma^{(o)}$

$$f = f^{(0)} + f^{(x)}, f^{(x)} \in \Sigma^{(x)}$$
 (3.2.15)

then, taking the scalar product of this with $b_i^{(o)}$ and using (3.2.2) gives

$$\langle b_i^{(o)} | f \rangle = \sum_j \mathsf{T}_{ij} \mathsf{f}_j + \langle b_i^{(o)} | f^{(x)} \rangle$$
 (3.2.16)

Together with (3.2.10), this implies that

$$\langle b_i^{(o)} | f^{(x)} \rangle = 0$$
 $i=1,...,m$ (3.2.17)

As a result, when the f_i are chosen as in (3.2.10), the operator P separates any function in Σ into a function $f^{(o)}$ in $\Sigma^{(o)}$ and a function $f^{(x)}$ which is orthogonal to $f^{(o)}$.

The functions $\{h_{(o)}^{i}\}$ in (3.2.11) are biorthogonal to the basis set $\{b_{i}^{(o)}\}$ in $\Sigma^{(o)}$ since

$$\langle h_{(o)}^{i} | b_{j}^{(o)} \rangle = \sum_{k} \mathsf{T}^{ik} \mathsf{T}_{kj} = \delta_{ij}$$
 (3.2.18)

In terms of finite element analysis, their use has been introduced by Brauchli and Oden [29] and cultivated by Oden [23]. Since (3.2.10) states that the best approximation of a function f in the space $\Sigma^{(o)}$ occurs when the expansion coefficients f_i equal the inner products of f with the biorthogonal basis functions in $\Sigma^{(o)}$, their application in approximating functions cannot be avoided. However, in determining the optimal solution of differential equations, it is fortunately not necessary to restrict the conjugate functions $\{b_{(o)}^{i}\}$ to be biorthogonal. In order to see why, note that a typical entry of the term Bf in equation (3.2.4) is

$$t_{i} = \sum_{j} B_{ij} f_{j} = \sum_{jk} B_{ij} T^{jk} \left\langle b_{k}^{(o)} | f \right\rangle \qquad (3.2.19)$$

Since both sets $\{b_i^{(o)}\}$ and $\{b_{(o)}^i\}$ form a basis in $\Sigma^{(o)}$, they are related by a linear transformation G

$$b_{(o)}^{i} = \sum_{j} G^{ij} b_{j}^{(o)}$$
(3.2.20)

The elements of the matrix B may then be written

$$B_{ij} = \langle b_{(0)}^{i} | b_{j}^{(0)} \rangle = \sum_{l} G^{ll} T_{lj}$$
(3.2.21)

and equation (3.2.19) becomes

$$t_{i} = \sum_{ljk} G^{il} T_{lj} T^{jk} \langle b_{k}^{(o)} | f \rangle$$

= $\sum_{k} G^{ik} \langle b_{k}^{(o)} | f \rangle$
= $\langle b_{(o)}^{i} | f \rangle$ (3.2.22)

In this equation, the relationship (3.2.20) between the bases $\{b_i^{(o)}\}$ and $\{b_{(o)}^{i}\}$ is immaterial. Therefore, in order to obtain the best approximation of f in (3.2.4), the components of the term Bf must simply equal the scalar products of the conjugate basis functions $\{b_{(o)}^{i}\}$ with the function f. Provided that all terms of (3.2.4) use it consistently, any conjugate basis set pair will generate an optimal approximation of f.

The evaluation of the remaining terms $\sum_{s=1}^{\infty} A_s y$ of equation (3.2.4) will now be examined in terms of the elementary matrix concept. First of all, each simple differential operator $P_{1s}^{D}sf$ may be factored into a sequence of elementary operators

$$P_{1S} D_{Sf} = E_{12} E_{23} \circ \cdots E_{f-1,f}$$
 (3.2.23)

and a sequence of spaces $\Sigma^{(i)}$ with bases $\{\mathcal{B}^{(i)}\}$ can be defined by

$$\Sigma^{(i-1)} = \{z : z = E_{i-1, i} \ y , y \in \Sigma^{(i)}\} \ i=1, \dots, f$$

The matrix elements in (3.2.5) may be written

$$A_{sij} = \langle b_{(o)}^{i} | E_{u1}^{i} E_{12}^{oo} E_{f-1f}^{b} j \rangle \qquad (3.2.25)$$

$$= \langle b_{(o)}^{i} | E_{01}^{i} e_{f-2}^{i} f_{f-1}^{p} f_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} \rangle$$

$$= \Sigma \quad \langle b_{(o)}^{i} | E_{01}^{i} e_{f-2}^{o} f_{f-2}^{e} f_{f-1}^{e} \rangle T^{pq} G_{qr} \langle b_{(f-1)}^{r} | E_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} \rangle$$

$$= \Sigma \quad \langle b_{(o)}^{i} | E_{01}^{i} e_{f-2}^{e} f_{f-2}^{e} f_{f-1}^{e} \rangle T^{pq} G_{qr} \langle b_{(f-1)}^{r} | E_{f-1}^{e} f_{f-1}^{e} f_{f-1}^{e} \rangle$$

$$= \Sigma \quad \left\langle b_{(o)} \stackrel{i}{\models} \stackrel{E}{=} \begin{array}{c} b_{p}^{(1)} \right\rangle T^{pq} G_{qr} \left\langle b_{(1)} \stackrel{|E_{12}}{=} \begin{array}{c} b^{(2)} \right\rangle \\ & \\ & \\ & \\ & \\ & \\ \end{array} \right\rangle T^{tu} G_{uv} \left\langle b_{(f-1)}^{v} \stackrel{|E_{f-1}, f^{b}}{=} \begin{array}{c} f^{(f)} \\ & \\ \end{array} \right\rangle$$

In the event that the basis functions $\{b^{(p)}\}$ and $\{b_{(q)}\}$ are biorthogonal

$$\langle b_{(p)}^{i} \rangle b_{j}^{(q)} \rangle = \sum_{q} G^{pq} T_{qr} = \delta_{pr}$$
 (3.2.26)

and the calculation of A_s may be performed by multiplying together the sequence of elementary matrices formed by the elementary operators

$$A_{s} = \prod_{k=0}^{f} E_{k}$$
(3.2.27)

where

$$E_{k_{ij}} = \langle b_{(k)} | E_{k,k+1} b^{(k+1)} \rangle \qquad (3.2.28)$$

Otherwise, the matrix products $T^{-1}G^{-1}$ must be inserted between

the elementary matrices as shown in section (3.2.2).

3.3 RELATION TO THE ENERGY NORM APPROXIMATION

It is well known that an energy norm can be defined for every boundary value problem such that the extremization of this energy norm in a projective subspace results in a least residual approximation [43]. It is therefore instructive to compare the conditions in equation (3.2.4) with the conditions obtained from energy norm extremization.

In order to do this, the differential equation (3.1.1) must be replaced by a variational functional. Of the commonly used variational principles, the most generally applicable is the one contained in the following theorem. <u>THEOREM 3.3.1</u> [43]. Let *D* be an arbitrary differential operator and D^* its adjoint

$$D^* = \{A: \langle Aw | y \rangle = \langle w | Dy \rangle, y \in \text{Domain}(D), w \in \text{Range}(D)\}$$
(3.3.1)

Then the solutions of the pair of equations

$$by = f'$$
 (3.3.2)

and

Π.,

$$D^* u = f \tag{3.3.3}$$

are the one and only pair of functions in a Hilbet space Σ which makes the functional

$$F(v,z) = \langle v | Dz \rangle - \langle v + z | f \rangle \qquad (3.3.4)$$

stationary.

The proof of this theorem may be found in reference [43]. Note that the functional

$$F(z) = \langle z \mid Dz \rangle - 2 \langle z \mid f \rangle \qquad (3.3.5)$$

often used with symmetric operators $(D^* = D)$ is a special case of (3.3.4) with the arbitrary function v set equal to z.

Approximate solutions of the differential equation (3.3.1) are obtained from the functional (3.3.4) by determining those elements $y^{(f)}$ and $u^{(1)}$ of the finite dimensional subspaces $\Sigma^{(f)}$ and $\Sigma^{(1)}$ which make F stationary. These approximate solutions will converge to the exact solution of (3.3.2) as the space $\Sigma^{(f)}$ is enlarged to Σ . As before, let $\{b_{i}^{(f)}\}$ be the basis in $\Sigma^{(f)}$ and $\{b_{i}^{(i)}\}$ be the conjugate basis of $\{b_{i}^{(1)}\}$ in $\Sigma^{(1)}$. Then

$$\dot{z}_{i=1}^{(f)} = \sum_{i=1}^{n} z_i b_i^{(f)}$$
 (3.3.6a)

$$\binom{1}{i=1} = \sum_{i=1}^{m} v_i b_{(1)}^{i}$$
 (3.3.6b)

and

υ

$$F(v^{(1)}, \mathbf{z}^{(f)}) = \sum_{i=1}^{n} \sum_{j=1}^{m} v_{i} \left\langle b_{(1)}^{j} | Db_{i}^{(f)} \right\rangle$$
$$- \sum_{i=1}^{m} v_{i} \left\langle b_{(1)}^{i} | f \right\rangle - \sum_{i=1}^{n} z_{i} \left\langle b_{i}^{(f)} | f \right\rangle$$
(3.3.7)

Differentiating this with respect to the coefficients \mathbf{z}_k and \mathbf{v}_k gives

44

$$\frac{\partial F}{\partial z_{k}} = 0 = \sum_{j=1}^{m} v_{j} \langle b_{k,1}^{j} | Db_{k}^{(f)} \rangle - \langle b_{k}^{(f)} | f \rangle$$

$$= \sum_{j=1}^{m} v_{j} \langle b_{k}^{(f)} | D^{*} b_{j}^{j} \rangle - \langle b_{k}^{(f)} | f \rangle$$

$$\frac{\partial F}{\partial v_{k}} = 0 = \sum_{i=1}^{n} z_{i} \langle b_{(1)}^{i} | Db_{i}^{(f)} \rangle - \langle b_{(1)}^{i} | f \rangle \quad (3.3.9)$$

Equation (3.3.9) differs from (3.2.4) only in that (3.3.9) does not contain the projection operators P_{Os}. Provided that these projection operators are added to the operator D in the above equations, (3.3.9) will result in the same rectangular matrix equation for (3.3.2) as is obtained from Galerkin's method. Equation (3.3.8) needs to be solved only if a solution to the adjoint problem (3.3.3) is desired.

3.4 ELEMENTARY MATRICES REVISITED

θV 7.

In this section, the elementary matrices of Chapter 2 will be derived from the expressions contained in section 3.2. The chief impediment to doing this is to determine functions $\{b_{(k)}\}$ which possess the sifting property in (2.1.9). Fortunately, Brauchli and Oden give finite dimensional delta functions with this property in reference [29]. These finite dimensional delta functions may also be produced by projecting

the Dirac delta function into a finite dimensional subspace by using (3.2.14). Let $\delta(x - x_0)$ be the Dirac delta function

$$\delta(x-x_0) = 0 \qquad \text{If } x \neq x_0 \qquad (3.4.3)$$

$$\langle 1 | \delta(x - x_0) \rangle = 1$$

and define $\Delta(x - x_0)$ to be the function obtained by acting with P_1 on $\delta(x - x_0)$

$$\Delta(x-x_{0}) = \sum_{ij} b_{i}^{(1)} T^{ij} \langle b_{j}^{(1)} | \delta(x - x_{0}) \rangle$$

=
$$\sum_{ij} b_{i}^{(1)} T^{ij} b_{j}^{(1)} (x_{0})$$

=
$$\sum_{ij} \Delta_{i} b_{i}^{(1)}$$

(3.4.4)

where

$$\Delta_{i} = \sum_{j} T^{ij} b^{(1)}_{j} (x_{0})$$
 (3.4.5)

The remarkable property of the function Δ is that while it is a smooth, well-defined function in the finite dimensional subspace $\Sigma^{(1)}$, it has preserved the important sifting property of the Dirac delta function for the space $\Sigma^{(1)}$. For example. let $y^{(1)}$ be an arbitrary function in $\Sigma^{(1)}$. Then

$$\langle \Delta^{(1)} (x - x_0) | y^{(1)} \rangle = \sum_{ijk} T^{ij} b_{j}^{(1)} (x_0) \langle b_{i}^{(1)} | y_k b_k^{(1)} \rangle$$

$$= \sum_{ijk} y_k b_{j}^{(1)} (x_0) T^{ij} T_{ik} \quad (3.4.6)$$

$$= \sum_k y_k b_k^{(1)} (x_0) = y^{(1)} (x_0)$$

Now let $y^{\binom{1}{i}}$ be equal to the basis function $\mathcal{I}_i^{\binom{1}{i}}$ which

interpolates on the set of points $\{x_{j}\}$. Then

$$\langle \Delta^{(1)}(x - x_j) | l_i^{(1)} \rangle = l_i^{(1)}(x_j) = \delta_{ij}$$
 (3.4.7)

Consequently, the finite dimensional delta functions $\{\Delta^{(1)}(x - x_i)\}$ form a biorthogonal set to the interpolatory basis functions.

Using the delta functions for the conjugate basis set $\{b_{(k)}\}$ in (3.2.2), the elements of the matrix E_k are

$$E_{k_{ij}} = \left\langle \sum_{k=1}^{k} (x - x_{i}) \mid E_{k,k+1} \mid z_{j}^{(k+1)} \right\rangle$$
$$= E_{k,k+1} \mid z_{j}^{(k+1)} \mid x = x_{i}$$

If $E_{k,k+1}$ is equal to $\frac{d}{dx}$ or to f(x), this expression reduces to (2.1.11) or to (2.2.8), in that order. Note that this result is independent of the norm used to define the space. If $E_{k,k+1}$ is equal to $P_{k,k+1}$ and an L_2 norm is taken

$$E_{k_{ij}} = \sum_{ml} b_{m} (x_{i}) T^{ml} \langle b_{l}^{(k)} | b_{j}^{(k+1)} \rangle$$

$$= \sum_{l} T^{il} S_{lj} \qquad (3.4.9)$$

This result is the same as (2.3.9). Therefore, the elementary matrices of Chapter 2 are identical to those obtained from (3.2.2).

The general form of the elementary matrix in equation (3.2.2) permits, however, many other basis functions to be used, besides the interpolatory basis of section 2.2, and some of these will be considered here. The most obvious such set is, of course, the monomials

$$b_i^{(k)} = x^i$$

(3.4.10)

With the conjugate functions being the biorthogonal set, the following elementary D and P matrices are obtained

where the p_i depend on the order of the approximation and all blank entries are zero. The function matrix in this case requires the expansion of the function in a Taylor series.

Another set of basis functions for which the D and P matrices generate simple forms are the Legendre polynomials. In this case

Again, however, the function matrix must be evaluated analytically to obtain the Legendre polynomial expression of the function. With a non-interpolatory basis, it is not possible to obtain a purely numerical method to generate function matrices.

3.5 <u>MULTI-DIMENSIONAL REGIONS</u>

In section 3.2, the technique for obtaining the approximate solution of a differential equation in a projective subspace was restricted to ordinary differential equations. Here, the extension of this method to partial differential equations will be made in a simple and natural way.

In order to retain the usage of the ordinary elementary matrices in the discretization of partial differential equations, it is necessary to use multi-dimensional basis functions which are separable into a product of one-dimensional Newton-Cotes interpolation polynomials. There are two known sets of multi-dimensional functions in two different geometric regions having this property. In N-dimensional cubes, the product separable multi-dimensional functions are

$${}^{b}i_{1}\cdots i_{N} = \prod_{k=1}^{N} {}^{l}i_{k}^{(n}k^{)}(x_{k}) , i_{k}=0, \dots, n_{k} (3.5.1)$$

where the $\{x_k\}$ are N independent variables and the $\{l_i^{(n)}\}$ are, as before, the equi-spaced Lagrangean interpolation polynomials

$$l_{i}^{(n)}(x) = \frac{(-1)^{n-1}}{i!(n-i)!} \prod_{\substack{m \neq 2 \\ m \neq 2}}^{n} (\frac{n}{L}x - m)$$
(3.5.2)
$$l_{i}^{(o)}(x) = 1$$

In N-dimensional simplexes, the required functions are given by [45]

$${}^{b}i_{1}\cdots i_{N} = \prod_{k=1}^{N+1} \iota_{i_{k}}^{(i_{k})}(w_{k}) \qquad i_{k}=0,\ldots,n \qquad (3.5.3)$$

with

$$\sum_{k=1}^{N+1} i_k = n \tag{3.5.4}$$

Here $\{w_k\}$ are homogeneous coordinates [45]. The spaces generated by the functions (3.5.1) will be called $\Phi^{(n_1 \cdots n)}$ and those generated by (3.5.3) $\psi^{(n)}$. If N=1, then the two spaces coincide

$$l_{i}^{(n)}(x) = l_{i}^{(i)}(x) \ l_{n-i}^{(n-i)}(L-x)$$
(3.5.5)

since the one-dimensional form of both the cube and of the simplex is a line segment.

Of the two spaces ϕ and ψ , only ψ contains polynomials of degree *n* which are complete in every monomial of m_k of degree less or equal to *n*. (This is not to be confused with completeness in the sense of a Hilbert space; both (3.5.1) and (3.5.3) form a complete set of basis functions since any polynomial in m_k may be obtained by taking n_k large enough). However, in the functions (3.5.3), products of polynomials of different degree are umavoidable; with (3.5.1) all the n_k may be taken to be equal. This, plus the fact that the independent variables in a differential equation must be transformed into homogeneous coordinates when using (3.5.3), make the cubic basis functions the easier of the two to work with.

The ability of ordinary elementary matrices to generate derivatives, projections and function multiplications on the functions (3.5.1) and (3.5.3) is a result of both of these function sets defining an equi-spaced, right-angled grid of interpolation nodes in their respective coordinate systems. On this grid of points may be superimposed the coefficients of the expansion of a function in these bases as shown in Figure 3.5.1 for a three-dimensional space. The effect of an elementary operator acting on this function will be equivalent to a matrix multiplication of this cubic block of coefficients by the elementary matrices in the proper direction. Thus, if a function $y(x_1, \ldots, x_N)$ is expanded in either the functions (3.5.1) or (3.5.3)

$$y(x_{1},...,x_{N}) = \prod_{k=1}^{N} \sum_{i_{k}=1}^{n_{k}} y_{i_{1}}..., i_{N}^{b} i_{1}... i_{N}^{b} (3.5.6)$$

then the elements of the matrix equivalent of the equation

$$z(x_{1},...,x_{N}) = E(x_{p}) y(x_{1},...,x_{N})$$
(3.5.7)

are

$$Z_{i_{1}} \cdots i_{p} \cdots i_{N} = \sum_{q=1}^{n_{p}} E_{i_{p}q} y_{i_{1}} \cdots i_{p-1} q_{p+1} \cdots i_{N}$$
(3.5.8)

where E is the n_p^{*} th order elementary matrix of $E(x_p)$. Here $E(x_p)$ may act only on the variable x_p^{*} .

For the case N=2, equation (3.5.8) provides a particularly





Nodal points where $b_{i_1i_2i_3} = 0$ if $n_1 = 1$, $n_2 = 1$ and $n_3 = 2$. The numbers in the circles designate the values of i_1 , i_2 and i_3 .

useful result. In this case for an elementary operator acting purely in the direction x_1 , the matrix elements Z become

$$Z_{ij} = \sum_{q=1}^{n_{1}} E_{iq} y_{qj}$$
(3.5.9)

and for an elementary operator acting in the direction x_2

$$Z_{ij} = \sum_{q=1}^{n_2} E_{jq} y_{iq} = \sum_{q=1}^{n_2} y_{iq} E_{qj}^{\mathsf{T}}$$
(3.5.10)

where E^{T} is the transpose of the matrix E. Therefore, in discretizing equations in two dimensional spaces, operators acting in the direction x_{2} result in multiplication of the coefficient matrix from the left by the matrix equivalent of the operator, and operators acting in the direction x_{2} result in multiplication of the transpose of the matrix equivalent of the operator.

Now, any partial differential equation can be reduced to an ordinary differential equation by holding all but one of the independent variables constant. Thus, along the line

$$x_k = a_k$$
, $k = 1, \dots, i-1, i+1, \dots, N$ (3.5.11)

the partial differential equation

$$D(x_{1},...,x_{N}) \quad y(x_{1},...,x_{N}) = f(x_{1},...,x_{N}) \quad (3.5.12)$$

reduces to the ordinary differential equation

$$D(a_{1},...,x_{i},...,a_{N}) \quad y(a_{1},...,x_{i},...,a_{N}) = f(a_{1},...,x_{i},...,a_{N}) \quad (3.5.13)$$

The approximate range space $\Sigma^{(0)}$ can be easily determined for this equation and will be called $\Sigma_i^{(0)}$. An N-dimensional partial differential equation will generate N such spaces $\Sigma_i^{(0)}$ and, since the partial differential equation (3.5.12) is a combination of the N equations (5.3.13), the approximation range space $\Sigma^{(0)}$ of a partial differential equation must coincide with the intersection of the range spaces $\Sigma_i^{(0)}$

$$\Sigma^{(0)} = \bigwedge_{i=1}^{N} \Sigma^{(0)}_{i}$$
(5.3.14)

An example will help to fix these ideas. Consider the two dimensional Helmholtz equation in Cartesian coordinates

$$\frac{\partial^2 g}{\partial x^2} + \frac{\partial^2 g}{\partial y^2} = k^2 g \qquad (3.5.15)$$

and let the domain of the operator $(\nabla^2 - k^2)$ be approximated by the spaces $\Phi^{\binom{2}{2}}$. Then β may be expanded in the functions (3.5.1)

$$z(x,y) = \sum_{i,j=0}^{2} Z_{ij} l_{i}^{\binom{2}{2}}(x) l_{j}^{\binom{2}{2}}(y) \qquad (3.5.16)$$

The first simple operator $\frac{\partial^2}{\partial x^2}$ in (3.5.15) acting on this space gives

$$\frac{\partial^{2} x}{\partial x^{2}} = \sum_{i,j=0}^{2} Z_{ij} I_{j}^{(2)}(y) \frac{\partial^{2}}{\partial x^{2}} I_{i}^{(2)}(x)$$

$$= \sum_{k=0}^{0} \sum_{i=0}^{2} W_{kj} I_{j}^{(2)}(y) I_{k}^{(0)}(x) = w(x,y)$$
(3.5.17)

Since $w(x,y) \in \Phi^{(u,2)}$, the space $\Sigma^{(1)}$ of equation (3.1.3) is $\Phi^{(0,2)}$. Similarly, if $\Sigma^{(f)} = \Phi^{(2,2)}$, the range spaces of the second and third simple operators, $\frac{\partial^2}{\partial y^2}$ and $-k^2$, in (3.5.15) are $\Sigma^{(2)} = \Phi^{(2,0)}$ and $\Sigma^{(3)} = \Phi^{(2,2)}$

$$\Sigma^{(0)} = \Sigma^{(1)} \cap \Sigma^{(2)} \cap \Sigma^{(3)} = \Phi^{(0,0)}$$
(3.5.18)

Therefore, according to equation (5.3.14), the solution of (3.5.15) in the subspace $\Phi^{(2,2)}$ is given by the solution of the equation

$$(3.5.19)$$

$$P(o,o), (o,2) \frac{\partial^{2}z}{\partial x^{2}} + P(o,o), (2,o) \frac{\partial^{2}z}{\partial y^{2}} = k^{2} P(o,o), (2,2)^{2}$$

Now $P_{(o,o)}$, (o,2) is a function of y only since it must leave functions in x intact. Similarly, $P_{(o,o)}$, (2,o) is only a function of x. Consequently, according to the methods of section 3.2 and equations (3.5.9) and (3.5.10), the discretized form of (3.5.15) is

$$DZP^{T} + PZD^{T} \Rightarrow k^{2}PZP^{T}$$
 (3.5.20)

where

$$D = \frac{4}{L^2} [1 -2 1] \qquad (3.5.21a)$$

$$P = \frac{1}{6} [1 4 1] \qquad (3.5.21b)$$

$$\int \frac{1}{\sqrt{2}} \left(\begin{array}{c} z_{11} & z_{12} & z_{13} \\ z_{11} & z_{12} & z_{13} \\ z_{21} & z_{22} & z_{23} \\ z_{31} & z_{32} & z_{33} \end{array} \right) \qquad (3.5.21c)$$

The x and y coordinate axes have been drawn on equation (3.5.21c) to indicate how the coefficients are related to the coordinates of the corresponding interpolation nodes.

Of course, many multi-dimensional functions cannot be separated into a product of functions of one independent variable. For these functions multi-dimensional function matrices are defined by

$$F_{i_1...,i_N} = f(x_{i_1}^{(i)},...,x_{i_N}^{(N)})$$
 (3.5.22)

The function muliplication z = fy is then given by

$$Z_{i_1...i_N} = F_{i_1...i_N} y_{i_1...i_N}$$
 (3.5.23)

In order to facilitate the evaluation of the matrix product (3.5.23), and also the solution of equations such as (3.5.20), it is necessary to introduce the Kronecker product of matrices. This is done as follows. Let A be an $m \ge n$ matrix and B be an $s \ge t$ matrix. The Kronecker product of A and B is defined to be the $ms \ge nt$ matrix [46]

$$A \otimes B = \begin{pmatrix} A_{11}B \circ \circ \circ A_{1}nB \\ \circ & \circ \\ \circ & \circ \\ A_{m1}B \circ \circ \circ A_{mn}B \end{pmatrix} (3.5.24)$$

It is known to obey the following associative and distributive laws [46]

$$(A \otimes B) (C \otimes D) = (AC) \otimes (BD)$$
 (3.5.25a)
(A + B) $\otimes (C + D) = A \otimes C + A \otimes D + B \otimes C + B \otimes D$

In addition, let A_j denote the *j*'th column of A and define the matrix operation vec A to give the following *mn* component vector [47]

vec A =
$$\begin{pmatrix} A_1 \\ \circ \\ \circ \\ A_n \end{pmatrix}$$
 (3.5.26)

The application of Kronecker products to matrix equations of the form (3.5.20) is a result of the following [47]

vec (AB) = (
$$I_p \otimes A$$
) vec B = ($B^T \otimes I_m$) vec A (3.5.27)

By using this equation and property (3.5.25a), an equation of the form

$$\sum_{k} A_{k} Z_{k} B_{k}^{T} = F \qquad (3.5.28)$$

where the A_k are mxn matrices, Z is an nxt matrix of unknown coefficients, the B_k are ext matrices and F is a specified mxe matrix, yields

$$\sum_{k} (B_{k} \otimes A_{k}) \text{ vec } Z = \text{vec } F \qquad (3.5.29)$$

where the $B_k \otimes A_k$ are $ms \times nt$ matrices, vec Z is an ntcomponent vector and vec F is an ms component vector. Equation (3.5.29) is a standard rectangular matrix equation in terms

of the unknown coefficients Z_{ij} .

1

In terms of the Helmholtz equation (3.5.15), the above result means that the discretized form of the equation (3.5.20) may be written as

 $(P \otimes D + D \otimes P)$ vec Z = k^2 $(P \otimes P)$ vec Z (3.5.30) From (3.5.25b) and the identity

 $(A + B)^2 = A^2 + A \otimes B + B \otimes A + B^2$ (3.5.31) equation (3.5.30) may be written in the symmetric forms

$$[(D + P)^{2} - D^{2}]$$
 vec Z = $(1 + k^{2})$ P² vec Z (3.5.32a)
 $[(D + P)^{2} + (D - P)^{2}]$ vec Z = $2k^{2}P^{2}$ vec Z (3.5.32b)

These results are valid for all D and P matrices of second degree. If a biquadratic approximation is used for the unknown, equation (3.5.30) becomes

$$\frac{4}{3L^{2}} \begin{bmatrix} 1 & 1 & 1 & 1 & -8 & 1 & 1 & 1 \end{bmatrix} \text{ vec } Z$$

$$=\frac{k^{2}}{36} \begin{bmatrix} 1 & 4 & 1 & 4 & 16 & 4 & 1 & 4 & 1 \end{bmatrix} \text{ vec } Z$$
(3.5.33)

For two-dimensional matrix products of the form (3.5.23), equation (3.5.27) implies that

vec Z = vec (FY) =
$$(F^{T} \otimes I_{N})$$
 vec Y (3.5.34)

Therefore, as in the one-dimensional case, the function matrix for two-dimensional problems is a diagonal matrix of the two interpolation point values of the multiplying function.

CHAPTER IV

DISCRETIZATION OF ARBITRARY DIFFERENTIAL EQUATIONS

In the numerical solution of differential equations, the step that has always been the most laborious and time consuming has been the discretization process. Aside from a single known attempt at automation [48], the usual procedure has been to churn through either Taylor series expansions or tables of integration formulae by hand or in sophisticated cases with the aid of computational machinery until the coefficients in the discretization process were evaluated for the particular differential equation of interest.

The most automatic discretization procedure presently available is the Taylor series method of producing finite difference equations. In this method [49], the differential terms in the Taylor series expansion of the solution of a differential equation are matched to the components of the differential operator on an arbitrary set of points. Since a general expression for the Taylor series expansion of an arbitrary function is known, the procedure does not require algebraic manipulation once the point set is chosen.

Silvester has reported a fully automatic computer program for discretizing arbitrary differential operators by this method [48]. However, since its discretization accuracy depends upon the locations of the point set, and since there is no a priori way of determining good locations, this "automatic" procedure is a hit or miss approach.

Using the elementary matrix concept of Chapter 2, the automatic discretization of any differential equation can be accomplished in a direct manner. First, it is necessary to break the operator into its elementary analytic factors and to form the corresponding elementary matrices. The assembly of these matrices to form simple operators is performed according to (3.2.27) by multiplying the elementary matrices together, keeping in mind that the degree of the range space of each operator must coincide with the degree of the domain of the next elementary operator acting upon it. Provided that the simple matrices are augmented by the proper projection matrices, the compound operator is assembled according to (3.2.4) by adding the matrices corresponding to the simple operators. Finally, the forcing function in equation (3.2.4) is added by evaluating the analytic function on the interpolation point set for the space $\Sigma^{(o)}$ as given by (3.2.22) and (3.4.6). Thus, the method of assembling elementary matrix factors to form a discrete matrix equation is completely analogous to the mathematical operations performed on the parent operators in the analytic differential equation.

This chapter contains the steps required to convert this procedure into a direct numerical algorithm for the automatic discretization of arbitrary differential equations. The discussion is complemented by listings of Fortran subroutines for the generation of matrix equivalents of ordinary and two-dimensional partial differential equations. The aim of these programs is to provide a very simple, accurate, and efficient

method of discretizing differential equations for anyone having access to a Fortran compiler.

4.1 COORDINATE-INDEPENDENT FACTORS

Since discretized equations are often used repeatedly in different locations, a major practical consideration in developing a discretization algorithm is to separate the coordinate-dependent factors from the coordinate-independent ones. Of the three elementary matrices, only the function matrix is entirely dependent on coordinate information, although the differentiation matrix does depend on the length of the interval taken. This suggests that the discretization process should be pursued on two levels: the high-order differentiation and projection matrices should be evaluated and stored separately from the function matrices, which should be generated only when needed in a specific location.

In order to avoid the time-consuming and involved computations required for generating the elementary D, P and FU matrices, it is natural to store the numbers in computer programs which use them. Figure 4.1.1 contains one method for storing these matrices. It is a Fortran block data subprogram in which the matrix elements given in Chapter 2 are stored in hexadecimal units for a 32-bit word computer. Due to the symmetry properties of these matrices, only the top halves of even rowed matrices and the top halves plus one line of odd rowed matrices are stored. This requires 535 memory locations for all matrices corresponding to the polynomials of degrees one through nine.

Figure 4.1.1a

DECL9315 2 Z417DC9C4, ZC2168000, Z412CD5F9, ZC113A732, Z41733333, ZC CD243C, F1. 7 BLOCK DATA DECL9320 3 Z403F1C75, ZC1100000, Z401ABB01, ZBF605065/ DATA FU5 /Z41100000, Z3FBDBD47, Z8F4EA943, Z3F300000, Z0 000000 DECL9325 THIS SUBPROGRAM CONTAINS THE ELEMENTARY DUP AND FU MATRICES OF 1 Z41128778, Z4030743C, ZC0190000, Z0100000, ZC0548475, Z4 E5001,000 2 Z40960000, Z00000000, Z40209C44, ZC030743C, Z40960000, Z0 000000, DECL9330 FIRST TO MINTH OPDER. THE NUMBERS ARE STORED IN HEXADECIMAL FORM UNDER THE APPROPRIATE MATRIX NAME. THE ELEMENTARY MATRICES DECL9335 3 ZBFF9A535, Z40118EEC, ZC0190000, Z01700000, Z3F25=247, ZFF273+41/* THE PLACED IN A COMMON BLOCK CALLED (DPFU1) 544 WORDS LONG, THE DEC1.9340 DECL9345 LARAY MANA CONTAINS THE LOCATIONS MINUS DNE OF THE FIRST 4 Z3F300000/ DATA P7 /Z40F3CFA7, Z8F6E488F, Z3 BRAESE, Z8E9FF010, Z4 555111, 1 1 Z40D6979F, ZC01C2E95, Z3F620011, ZC (-F512, Z40360060, Z4 22625, F) DECI.9350 ELEVEN OF EACH OF THE P MATRICES. DECL9355 2 ZCO1E+005, Z41144986, Z3FE69CCD, Z4 +36640, Z40991FFB, Z011+4984, DECL9360 DIMENSION NMA(9) 3 ZCO2D3C14, Z3FE69616, Z40991FF5, Z4 -++40C, Z40244321, 20-15735E, 55. DI-ENSIGN P1(2)+p1(2)+ P2(3)+p2(3)+ P3(8)+p3(8)+ DECL9365 4 ZCO1E#002, ZCc555148, ZBF005440, Z3F910143, Z3F620000, 23,C3C3FC, *** P4(_);,:4(10); P5(18);D5(18); P6(21);D6(21); DECL9370 × P7(32), 7(32), P8(35), D8(36), P9(50), D9(50) DECL9375 5 Z3F223DB2, Z8F173656, Z8E9FFFE0/ DEC1.9380 UATA D7 /ZC2I22666, ZC0547254, Z3F6502CD, Z3F140000, Z4F310000; 1 ZC19F4587, ZC02DC443, ZC0112666, ZC7498000, Z4210EFDF, ZC18F41 () DIMENSION FU0(1), FU1(4), FU2(6), FU3(12), FU4(15), FU5(24), FU6(28), DEC1.9385 FU7(40), FU8(45) = 20-314-004 CO 1201 / JPFL1/ NTA, P1, D1, FUO, P2, D2, FU1, P3, D3, FU2, P4, D4, FU3, DECL9390 2 Z408EEAAA, Z4251AAAA, ZC1840946, Z41461265, ZC185F100, 3 Z416FAFC5, ZC12C9230, Z4185FCn0, Z47106666, ZC130F2CC, 24-E:4111 # P5,05,FU4, P6,06,FU5, P7,07,FU6, P8,08,FU7, P9,09,FU8 DECL9375 DECL9400 4 ZCOBEEAAA, ZC182A4AA, Z40CAAAPA, ZCC357642, Z40112666, 24.100000, 0474 \ 4/0,5,15,37,69,120,186,278,390/ DECL9405 5 ZC017CAE3, Z3F58957D, Z8F140000/ P1 /Z40800000, Z40800000/ 04-4 DATA _ FU6 /Z41100000, Z3F654600, 78F300897, Z3F140480, Z0100000, 1 D1 /ZC1100000, Z41100000/ DECL9410 JATA. DECL9415 1 Z41128008, Z402830A2, Z8FFFAF3A, Z00000000, Z0058:538, Z4:105413, F 2272 FUD /Z41100000/ 2 Z406DUD48, Z0000000, Z4042A300, 70551352, Z40035085, Z01000000, DECL9420 117A P2 /Z45A/A/A5, Z40AAAA9C, ZC0555551/ DECL9425 3 ZC0221354, Z46210326, ZC028F21E, Z00000000, Z3F457162, 1=:965713, T LATA CZ /ZC13000000 Z41400000, ZC1100000/ DECL9430 4 Z3F457162, Z0000000, Z8F163655, Z3F133706, Z8F138300/ 2474 FU1 /Z41100000, Z40800000, Z00000000, Z40800000/ UATA F8 /Z4nF78919, Z8:578CAB, Z3F232F38, Z8F122294, 24:4235=1; P3 /24-C66-67, 7C0100000, 240ACCCC3, 2408FFFF9, 2CCACCCCA, DECL9435 2274 1 Z40098101, ZC51A6538, Z3FA43953, ZC5E7CC09, Z4C3141F6, Z4 E393---* DECL9440 1 Z-03FFFF9, Z4-399994, ZC0100000/ 2 ZCO2DC573, Z411CF7FD, Z401A4871, Z401F4834, Z40C82685, ZC+2435F975F D3 /ZC1580000, Z40200000, Z41900000, ZC1360000, ZC1480000, DECL9445 2174 3 ZC04FA729, Z403E3F4E, Z405ED658, Z411CF7F8, Z40515F85, ZC 472277, 11 17 1 Z4135:000, Z41100000, ZC0200000/ DECL9450 LITA FUZ /Z41100000, Z4038E38E, Z00000000, Z40E38E38, Z00000000, DECL9455 4 Z3E3Au500, ZC0E78F9A, ZC02E2D21, Z40287DC1, Z8F891840, 74-42851-1 DECL9460 5 Z3FE30D40, ZBFC60F7E, Z3F382D3F, ZBF846DE0, ZBF100384, Z3F190F44, 17 1 ZC0107107/ P4 /Z40D2=09E, Z8F8043F7, Z409C09CE, Z40AA7C69, ZC0EA0E95,DECL9465 6 Z8E837C50/ C174 UATA D8 /ZC215BE28, ZC067618F, Z3F7E0E27, Z8E7C=40C, Z4740C010, E1 1 1 ZC1C2DE23, ZC05204F8, Z3F1A3CF0, ZC270C000, Z42167856, ZC198EE1, E1 11 1 Z-0780783, Z40900904, ZC01F0917, ZC0270275, Z3F409303/ DECL9470 D4 /ZC1855555, ZC0161F9A, Z42100000, ZC148B0FC, ZC1C00000, DECL9475 LATA 2 Z402Do3EE, Z42955555, ZC212A441, Z41D7=C40, ZC1951F51, ZC56C1101 1 2.157.425, Z41555555, ZCOFCD6E9, ZC1100000, Z4022C3F3/ DEC1.9450 04-4 FU3 /Z41100000, Z401E0000, ZC0100000, Z00000000, Z4110E000, DECL9485 3 Z41EA26B6, ZC152E882, Z41A56073, Z42599999, ZC188±365, Z4125F8T2+11 1 Z4095: 000, Z0000000, ZC0360000, Z40900000, Z00000000, Z3FA000000, DECL9490 4 ZC11785A5, ZC2255555, Z41359EEC, ZC0D35090, Z40518668, Z4192491-, ... 5 ZCOC9E924, Z402E3183, ZCO103DC7, ZC1100000, Z40157483, Z8:428804, 24 2 20010-000/ DECL9495 LATA P5 /Z4CE5887C, ZBF99877C, Z3F300007, Z40844670, Z4CC05442, DECL9500 6 Z3F17+9CF/ DATA - FU7 /Z41100000, Z3F624558, ZBF1F1400, Z3EEECC80, ZBF460000, C3F6245 1 ZCO18: FFB, ZC11088CD, Z40608853, Z4095FFF3, Z411088CB, ZC018CB55, DECL9505 1 Z00000000, Z4112CF45, Z401FB9C0, ZBFA8C438, Z3F620000, Z^*CCTCC**** 2 ZC0645174, Z4111B37C, Z4052460E, ZC01EA000, Z00000CC0, Z4*58E-1*, ** 2 24095-FF4, ZC084456F, Z3F485A99, ZC018FFFE, Z401A747A, ZBr788D00, DECL9510 DECL9515 3 Z3F3000107 D5 /ZC1B6414A, ZC034E000, ZBF600000, Z42190000, ZC161D400, DECL9520 3 ZL05F2040, Z40E48986, Z40992000, Z00000000, ZC03C30DF, Z4 34E1-1, 👘 uara -1 Z-3535555, ZC2193030, Z41874955, ZC15DC000, Z421044AA, ZC12DB400/DECL9525 4 ZL03E5414, Z40952000, Z0000000, Z40165008, Z00158600, Z41154625, Ellon 2 Z4150.000, ZC1640000, Z4007A000, ZC0535555, Z41100000, ZC01E7555, DECL9530 5 ZL01E4000, Z0n000000, Z8F757218, Z3F599400, Z8F514208, Z3F520000, E 111 6 Z00000000, Z3EE09080, ZBEASCOOD, Z3E8F4780, ZBEA00000 DECL9535 3 Z3F601000/ CATA _ FU4 /Z41100000, Z401205BC, ZBF89A027, Z00000000, Z411205BC, DECL9540 UATA P9 /Z40FA60DA, Z6F43E654, Z3F20E571, Z8F136070, Z3:22F240, 1 Z4055.01A, Z00000000, Z004816F0, Z40CE703A, Z00000000, Z401A36E2, DECL9545 1 Z403297C1, Z4cDA2401, ZC0167732, Z3FBB94F8, Z8F195145, ZC C15F54, 4 2 ZC0275254, Z00000000, ZBF4816F0, Z3F58C01A/ DECL9550 2 Z403420CB; Z40E873E2; ZC036F458; Z3F80C105; Z4110E32E; Z401493-E; EE 3 Z3FC20DFD, Z40F784F3, ZC0227424, ZC12C448F, ZC0609515, Z4-54112, TE1_1 DATA P6 /240EE03D8, ZBF651C8A, Z3F33C575, Z406BB824, Z40CECF5A, DECL9555 4 Z4016667E, Z409R0A1F, Z412C448F, Z43806189, ZC065510C, Z4 44419, ISS 5 Z409R09D5, ZC11D8327, ZC062171F, Z4062C65F, ZC03C098F, IC 127351, ISS 1 ZCO1814FC, ZC11004C6, Z4047531E, Z40868883, Z41167105, ZBF7EFC6E, DECL9560 2 Z-07232C3, ZC11034C4, ZBFC39A67, ZC0128E71, Z40688826, Z3F8A10A7, DECL9565 DECL9570 6 Z40C45E8E, Z402D92A8, ZC02C0562, Z401C80EF, Z3F800040, ZC0329775, -1 3 Z3F13_3D0, ZCC11F408, Z8F1C840A, Z3E40FC00/ 7 Z5FC05020, Z3FB9DB80, Z8F784F04, Z8F1950C0, Z3F59F160, Z3 103340, 101 8 Z8F15298C, Z3EDD1056, Z3E22FD60/ LIA 06 /ZC1E53333, ZC0443C64, Z3F2CB20F, Z42240000, ZC17F24D8, DECL9575 1 Z3F42.9F6F, ZC2200000, Z41C452E7, ZC16E15CA, Z42280000, ZC:622E5D, DECL9580

σ ÷ .

Figure 4.1.1b

D9 /ZC21975F1, ZC0720B5E, Z3F893B32, ZBF112A5B, ZBF460000, DECL9855 DATA 1 Z42510000, ZC1E86924, ZC06E6D84, Z3FA9A16E, Z3F4116D8, ZC2A20000, DECL9860 2 Z421CE1EB, ZC1B6CAF8, ZC0145A50, ZC01FE4CC, Z42FC0000, ZC21CC029, DECL9865 2 Z42111877, ZC1A8417A, Z40CEB800, ZC3118800, Z4218526D, ZC169C44D, DECL9870 4 Z41CB43C7, ZC1AE6840, Z42E2CCCC, ZC213FB8E, Z41510A85, ZC12E826B, DECL9875 5 Z41AE6840, ZC27E0000, Z41A7FB9C, ZC1261896, Z41102DA5, ZC1CEP800, DECL9880 6 Z422E4924, ZC13=6673, Z40C98ECE, 7C04972A9, Z401FE4CC, ZC1A20000, DECL9885 7 Z40CA3876, ZCh2BFCh8, Z3FDA0439, ZBF4116DB, Z41100000, ZCh138E70,DECL9890 DECL9895 8 Z3F3D643B, Z8F133239, Z3E460000/ FU8 /Z41100000, Z3F48376D, ZBF151510, Z3E904812, ZBE563CA2, DECL9900 DATA 1 Z00000000, Z4112C0DB, Z40181812, ZBF736CDB, Z3F38FDA8, Z00000000, DECL9905 2 ZC0694099, Z4112726E, Z403F1F87, ZC0158F28, Z00000000, Z4r6ED868, DECL9910 3 ZC06853F5, Z40FC7E1F, Z40788B44, Z00000000, ZC05E0548, Z40490988, DECL9915 4 ZC04EE769, Z40BCA4A3, Z00000000, Z4038EBAC, ZC02885EA, Z40241204, DECL9920 ZCO2E6F6B, Z00000000, ZCO16E45B, Z3FF88C89, Z8FC9FE7F, Z3FD88312, DECL9925 5 Z0000000, Z3F578650, Z8F396A62, Z3F2C64F1, Z8F2C6264, Z00000000, DECL9930 DECL9935 7 ZBE966EDA, Z3E60604A, ZBE482409, Z3E44FD4E/ DECL9940 END

1111 1111

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A program for building up matrix operators of an arbitrary order from the block data subprogram is given in Figure 4.1.2. Written in Fortran, the program is in the form of a subroutine called OPRATR, the input to which is an NC by ND array in OP which may contain any elementary matrix, including the identity matrix. It returns an NR by ND array in OP which is equal to the product of the original matrix times (NC - NR) successive elementary matrices. The type of matrix used depends upon the value of KIND; if KIND = 0, differentiation matrices are used; KIND = 1 results in the application of projection matrices; KIND = 2, the FU matrices. Thus, if on input OP is a 6 \times 6 identity matrix, NR is equal to 3 and KIND equal to 0, on output OP would contain the 3 x 6 matrix which differentiates the coefficients of fifth order polynomials into second order ones.

The subroutine OPRATR uses the symmetry or antisymmetry properties of the elementary matrices to speed up the computation. This means that an additional index called IEO must be supplied to OPRATR to designate whether the input matrix is symmetric (IEO =1) or antisymmetric (IEO = -1) about the matrix centroid.

In order to use the subroutine OPRATR to generate the coordinate-independent matrices for ordinary differential equations automatically, a method of separating these operators in the differential equation must be devised. On a computer, one method of distinguishing the mathematical components of a differential equation is to assign each mathematical entity in the equation a numerical code. If the code is chosen cleverly,
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	SUBREUTINE EPRATR(SP/IEG/NR/NC/ND/KIND) UPRATR BUILDS UP ELEMENTARY DIFFFRENTIATION AND PROJECTION MATRICES FROM VALUES STORED IN THE COMMON BLOCK DPFUL. THE INPUT TO EPRATR MUST BE AN INCI BY INDI ARRAY IN IOPI WHICH IS CO.VERTED TO ATLINE! BY INDI ARRAY BY THE PRE-MULTIPLICATION DF IDDI WITH THE APPROPRIATE ELEMENTARY MATRIX. (KIND = 0: PROJECTION; KIND = 1: DIFFERENTIATION; KIND = 2: RAISING)	DECL7525 DECL7530 DECL7535 DECL7540 DECL7545 DECL7555 DECL7555 DECL7555 DECL7560 DECL7565
	DEJSLE PRECISION VECTOR(5) DIMENSION OP(1) DI ENSION NMA(9),0(535) EDMON / OPFUL / NMAxO IF(NR .EQ. NC) RETURN NSION = 1 IF(KIND .EQ. 1) MSIGN = T 1 M = IA4S(UC - NR) NR = NC DD 5 N = 1,M AC = NK	DECL7570 DECL7575 DECL7580 DECL7585 DECL7590 DECL7595 DECL7600 DECL7600 DECL7605 DECL7610 DECL7615 DECL7620 DECL7625
	IADD IS THE LOCATION OF THE ELEMENT BEFORE THE FIRST ELEMENT OF THE REQUIRED ELEMENTARY MATRIX IN DPFUL	DECL7630 DECL7635 DECL7640
3	$iF(KIND = EQ_{+} 2) GD TD 3$ $iR = NR - 1$ $lLIMIT = (:R + 1)/2$ $iDD = NMA(NR) + ILIMIT \neq NC \neq KIND$ $GC TO =$ $iR = NR + 1$ $ILIMIT = (iR + 1)/2$ $iADD = NMA(NC) + (NR/2) \neq NR \neq 2$ $JI'C = \Rightarrow NC$ $JI:R = = NR$	DECL7645 DECL7655 DECL7655 DECL7660 DECL7665 DECL7675 DECL7675 DECL7680 DECL7680 DECL7685 DECL7695 DECL7695 DECL7700
•	THE AREAY TEPT IS MULTIPLIED BY THE HALT OF THE DECREMENT MATRIX DE 2 J = 1,ND J1::C = J1NC + NC J1::R = J1NR + RR DE 1 = 1.1 MIT	DECL7705 DECL7710 DECL7715 DECL7720 DECL7725 DECL7730
1 2	$v_{ECTOR(I)} = 0.00$ $i_{K} = I + IDD - ILIMIT$ $DC = 1 < IDD - ILIMIT$ $VECTOR(I) = VECTOR(I) + DBLE(O(IK) \neq OP(K + J1NC))$ $UC = 1.1LIMIT$ $OP(I + J1NR) = SNGL(VECTOR(I))$ $i_{F}(ILIMIT - GE, NR) = 0$ THE FIRST	DECL7735 DECL7745 DECL7750 DECL7755 DECL7755 DECL7755 DECL7765 DECL7770 DECL7775 DECL7775
	HALF	DECL7785 DECL7790

and the second secon

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	IEC = NSIGNAIEC
	SIGN = FLOAT(IED)
	JINR = - NR
	J1NC = NR*(ND + 1) + 1
	DD 5 J = 1, ND
	J1NR = J1NP + NR
	JINC = JINC - NR
	DD 5 I = 1,ILIHIT
5	$UP(JINC - J) = SIGN \pm OP(I + JINR)$
6	CONTINUE
	RETURN
	END

محجب المتحدة والمتحدث والمستحج

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والمستريق والمستحدث والمروط المتركب والمحاصر المراسية أنفست المستحر والحراص والمحاصر بسر مستسب حرارت تقسم

DECLTRSS DECLTRSS

DECL7805 DECL7810 DECL7815

then certain ranges of values will indicate the proper algebraic 67 steps to be followed. Consider, for example, the following coding procedure

NUMBER	OPERATION
3000	end
2000	minus sign
1000	plus sign
200	y (the unknown function)
-1	$\frac{d}{dx}$
o	o
o	o
o	0
-9	$\frac{d^9}{dx^9}$

With this code, a number greater than or equal to 1000 designates the separation locations for simple operators, the presence of the number 200 in a simple operator indicates that the term contains the unknown while its absence signifies that the quantity is a known function, and the absolute values of the numbers from -1 to -9 indicate the corresponding orders of differentiation. Furthermore, the number of derivatives in a simple operator may be obtained by adding the magnitudes of the numbers between -1 and -9 for that operator. The order of the projection matrix to be used in each simple operator is then found by subtracting each of these numbers from their largest value.

A subroutine OPGEN1 has been designed to operate on the

SUBROUTINE CPGENI(INSTR,ISTART,NP,ND,LOC,LO,LFCN,LF,	DECL6655			$DD 6 I = 1_{\text{MTYMSN}}MN$	DECTRAS
# OPS+1 ROW+LPOWER+OP+IPROJ)	DECL6660		5	UP(I) = 1.*SIGN	DEC16930
	DECL6665			1EO = 1	DECL6935
THES REPORTING EVICTIMES THE COORDINATE-INDEPENDENT FACTORS	DEC15670			ISN = N	050169-0
	DECL6675			IF(N SEE HD) GD TD 9	860194-3
TUR 195 DISCRETIZATION DI CINCAR ORDINARI DI COLONIALI	DEC1 6680			IE(IPROL : EC. 1) M = NR	5261 49-3
EQUALINAS, THE EQUALITY DISCRETIZED IS CODED IN THE WARAC	02010000	r		The set of	72714345
TINSTRY STARTING IN THE "ISTART COCATION. THE MATRIX FACTORS	DECLOSOD	ž		THE REALIZED DOBLECTION MATRIX IS SUALLYTED	55514313
EVALUATED ARE RETURNED IN THE ARRAY 'OPS' WITH THE NUMBER OF	DECLOSIV	5		THE REQUIRED PROJECTION MAINIA IS EVALUATED	02010-00
RDUS IN EACH MATRIX BEING GIVEN BY THE ARRAY "LROW".	DECL6695	۲.			UELLEVEJ
	DECL6700			I = LGCS	DECLERIQ
$C_1 = C_1 = C_1 + C_2 $	DECL6705		7	I = I - 1	DECLARTS
	DECL6710			IF(INSTR(I) = 0.300) IDP = -1	050.6910
	DECI 6715			1F(1PR03 .ME. 1) GO TO 8	1=115772
	05016720			TECTARS(1)STR(T)) .GE. 10) GD TD 8	F201 6995
LUCL = ISIARI	00014725				727:4014
LUCS = ISTART				0 = 0 = 1.050017	05515030
	DECLOTO		9		
FIRST THE LECATION OF THE + OR - SIGNS IS DETERMINED	DECT0/33			IF(IPR0J 40m 17 90 10 9	Carry and
	DECL6740			CALL DPRATR(DP, IEU, M, M, N, O)	0501-010
1 1055 - 1055 + 1	DECL6745			MN = M	DECL7015
1 E/THS (2015)	DECL6750		9	L = L - 1	DECLIFCEO
	DECL6755			IF(IA35(INSTR(1)) .GE. 10) GO TO 10	2201_7025
2134 = 1	DECI 6760	c			1=107-32
1F(185) R(LOUL) +E0. 2000/ SIGN = - 1.	DEC1 6765	ē		THE RECUIRED DIFFERENTIATION MATRIX IS EVALUATED	02017435
LOOL = LOOS	DECL 5770	ř		The Readings Strictering Control is Eveloped	2201 72.2
L = LOOS - 1		C			D
IF(INSTR(L - 1) .EC. 200) L = L - 1	DECLOTIS			CPUMER(LB) = CPUMER(CB) = INSTR(L)	52627642
IF, INSTR(L) .EQ. 200) GD TD 3	DECL6780			H = MN + IIISTR(L)	0201/150
	DECL6765			CALL DPRATE(DP, IED, M, M), N, 1)	CEC17055
THE CHARENT TERM IS A KNOWN FUNCTION - ITS LOCATION AND SIGN	DECL6790			MN = M	DEC17047
ARE STREED IN ILECNI	DEC1.6795			GO TO 9	DECL7Ca5
	DECL6800		10	IF(INSTR(L) .GE. 1000 .OR. IABS(INSTR(L)) .LT. 10) GO 70 11	DECL7070
	DFC1 6805			L = L - 1	DECL7075
LF - LF T 1	DECI 6810		1		
Lr(x(Lr)) = Lm(r(x))	05016815		11		0=0:70=5
	0-016920	c	••		0201 7020
IF,INSTR(L) .LT. 1000) GU TO 2	02010820	č		THE MATRIX FORNITED IS FIGED IN 100.	DEC1 7035
LF = LF + 1	DECLORZO	ž		THE MATRIX COMPOSED IS STORED IN THE	
LFCN(LF) = L + 1	DECLORSO	C			02067100
IF(INSTR(LOCS) .GE, 3000) RETURN	DECL6832		1	$L_{C}(L_{O}) = IDP*LUC(L_{O})$	CECL/112
66 70 1	DECL6840			LKOW(LD) = M	DECL7110
	DECL6845		1	MTYMSN = M*N	5ECL7113
THE STUDIE PREDATED ENCOUNTERED CONTAINS THE UNKNOWN Y	DECL6850			IF(LO - EQ - 1) LOPS = - MTYMSN	DEC17120
145 DIMPES DEFALCY CHOROMETERS CONTAINS THE CONTAINS	DFCL6355		1	LOPS = LOPS + MTYMSN	DEC17125
	DECI 6850		i	DD 12 I = $1 \neq MTYMSN$	12017132
3 A = NS	DEC16865		12 1	PS(T + 10S) = 0P(T)	55617135
IDP = 1	DECL0800				
4 L2 = L3 + 1	02020870				
LGC(LC) = L - 1	DECLORID				9864/142
LPOWER(LD) = 0	DECTORED			51GN = 1.	25017150
N = M	DECLARES			1 (1N5/R(L = 1) .LT. 1000) GO TO 4	DEC17155
NTVNSN - MAN	DECL6890		t	LOC(LO) = - LOC(LO)	061710G
	DECL6895		1	IF: INSTR(LUCS) .GE. 3000) RETURN	D=C17145
THE ADVIDUDETATE UNIT HATSIN IS COEATED	DECL6900		(6 TD :	
THE APPRUPRIATE UNIT MATRIX IS CREATED	0=016905		;	END .	
	DECLA010				
DO 5 I # 1,MTYM5H	02010710				_
$5 \text{ DP}(\mathbf{I}) = 0$	DECL0912				5
NN = N + 1	DECE0420				œ

above principles and is presented in Figure 4.1.3. The instructions for the differential equation to be discretized are contained in the array INSTR with the first instruction in location ISTART. The numbers NR and ND give the number of interpolation points (degree of the interpolation polynomial plus one) in the range and domain spaces of the operator, respectively. During execution, the program checks successive values of INSTR, noting the starting and ending locations of simple operators in LOC and of functions in LFCN. As differential and projection operators are detected, their matrix equivalents are generated by calling OPRATR and stored consecutively in OPS. The arrays LROW and LPOWER give the number of rows in these matrices and the power to which the quantity (1/L) in the differentiation matrix is to be raised.

4.2 COORDINATE - DEPENDENT FACTORS

OPGEN1 completes the first level of the discretization process by evaluating all of the numbers in the matrix equivalent of a differential equation that are independent of coordinates. The second stage of the process, adding the coordinate dependent factors, is accomplished by the companion subroutine OPGEN2 shown in Figure 4.2.1. This subroutine takes each simple operator separately and evaluates the function matrices at the interpolation nodal coordinates given in XD. The previously calculated differentiation and projection matrices are then combined

<pre>SUBROUTINE_EPGEN2(INSTR;CENST;LEC;LPGWER;LRGW;L;ND; * XD;XT;N;NPT;GPS;LGPS;GP;TEMP)</pre>	DECL7180 DECL7185	DD 8 J = 1,ND IJ = 1; + 14	DECL7450 DECL7459 DECL7460
THE SUBROUTINE OPGEN2 TAKES THE COORDINATE-INDEPENDENT FACTORS	DECL7190 C DECL7195 C DECL7200 C	THE COMPUTED MATRIX IS STORED IN THE ARRAY 'OP'	DECL7455 DECL7470
APPROPRIATE CEORDINATE DEPENDENT PARAMETERS. THE POINT	DECL7205	IF(MN .EQ. ND) GD TO 7	DEC174-9 DEC1749D
COURDINATES ARE CONTAINED IN THE ARRAY 'XD' WITH THE QUANTITY	DECL 7210 DECL 7215	Ge To a	DEC17435
(1/L) STURED 10 (4)(+	DECL7220	$7 \text{ UP}(IJ) = FUNCT \neq \text{UPS}(IJ + \text{LOPS})$	DEC17490
DIMENSION INSTR(I);CONST(I);LOC(I);ÉPOWER(I);LROW(I);	DEC17225	8 CONTINUE	0202/473
<pre>xD(1),WT(1),DPS(1),DP(1),TEMP(1)</pre>	DECL 7230	IF(LUL(L + 1) LI + 0) REPORN $MN = M$	DECL7505
DOUBLE PRECISION DOUBLE	DECL7240	L = L + 2	DECL7510
1 H = LRGW(L)	DECL7245	GO TO 1	CEUL/319 75017530
LEPS = LEPS + M*MH	DECL7250	END	
1F(MN .EQ. ND) GD TD 4	DECL7250		
IF A WATRIX IS ALREADY STORED IN THE ARRAY 'OP', IT IS	DECL7265		
MULTIPLIED BY THE NEXT MATRIX FACTOR	DECL7270		
	DECL 7275		
IJ = C	DECL7285		
LON = LOPS = M JIMN = = MN	DECL7290		
DD 3 J = 1, ND	DECL7295		
J10N = J14N + MN	DECL 7300		
00 3 I = 1,M	DECL7310		
IJ = IC + ICM	DECL7315		
KJ = J135	DECL7320		
DBUBLE = 0.00	DECL7320		
UE Z K = 19/N	DECL7335		
$K_{J} = K_{J} + 1$	DECL7340		
2 DOUBLE = DOUBLE + DBLE(OPS(IK) +OP(KJ))	DECL7345		
3 FEMP(IJ) = SNGL(DOUBLE)	DECI 7355		
4 LBCL = IABS(LBC(L)) + 1))	DECL7360		
CHOCI - TANGTESOTE - IV	DECL7365		
THE NODAL POINTS ARE STEPPED OFF	DECL7370		
	DECL7380		
IE(M .FO. 1) GO TO 5	DECL7385		
STEP = $1.0/(FLDAT(n - 1)*WT(N))$	DECL7390		
$\lambda = XD(NPT) - STEP$	DECL 7399		
68 TB 6 - STCB - A F(ST(S))	DECL7405		
S = S = S = S = S = S = S = S = S = S =	DECL7410		
6 00 8 1 m 1 M	DECL7415		
X = X + STEP	DEC17420		
THE MARKE SE THE ENDETION AT THE NODAL DOINT IS EVALHATED	DECL7430		
THE VALUE OF THE FUNCTION AT THE MIDAL FORMETS EVALUATED	DECL7435		
FUNCT = FON(X, INSTR, CONST, LOCL, LOCL1) #WEIGHT	DECL7440		7(
ij = I - 8	DECL7445		0

.....

with the function matrices in the specified order to produce, in the array OP, the matrix corresponding to the simple operator.

The function values used by OPGEN2.are computed by the subprogram FCN given in Figure 4.2.2. This program requires the following extension of the instruction code

NUMBER	OPERATION
600	- sign within a function
500	+ sign within a function
400)
300	(
100	X (the independent variable)
19 _~ 15	auxiliary functions
. 14	logarithmic function
13	exponential function
12	cosine function
11	sine function
-101 ~ -199	constants
-201 _~ -299	exponentiation

With this code, FCN can generate the point values of any function configuration. The numerical values of constants and the powers to which an argument are to be raised are contained in the array CONST.

In order to permit the use of functions different from the four standard functions sine, cosine, exponential and log71

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Figure 4.2.2

ENVETTON SCALV INSTRACTINST DELLI DELL	DEC1.7855		GD TD 22	CECU3125
	DECL7860		9 IF(INSTR4 .GT 200) GD TD 10	DECL6130
THIS FUNCTION SUBPROGRAM RETURNS THE VALUE OF A FUNCTION STORED	DECL7865 C	C	THE ADDITION TO TO DE BALCED TO COME DOWER	02010102
IN LOCATIONS (LOC2) TO (LOC1) OF THE ARRAY (INSTR) AT THE	DECL7870 C	C r	THE ARGUMENT IS TO BE RAISED TO SOME POWER	DECLS: 45
LOCATION (X)	DECL/8/2 C	C	F(1) = F(1) ++CONST(- 200 - INSTR4)	DECLBISO
TNTERED+2 TNCTD+1 EVEL TOP	DECL7885		GD TD 22	DECL8155
DIMENSION INSTR(1), CONST(1), F(20), EVEL(20), IDP(20)	DECL7890	1	O IF(INSTR4 .GT 100) GD TD 12	DECLBISC
i = 0	DECL7895	_	IF(INSTR4 .LT 150) GO TO 22	01000102 05519177
1PARA = 0	DECL7900 C	C C	CONSTRUTE (LCD DDICTNATE & NEW TERM	DECERT
L = LGC1 + 1	DECL7905 C	c r	CONSTANTS AFON ORIGINATE & NEW (ENH	CECLBIEC
THE WE DE THE CHARTER AS DETERMINED BY READING THE INSTRUCT	DECL7910 C	C	I = I + 1	DECLEIES
THE VALUE OF THE FUNCTION IS DETERMINED BY READING THE INDIKOUM	DECL7920		F(I) = CDNST(- 100 - INSTR4)	DECL5190
112N3 FRU, F19F1 - 0 CE. 1	DECL7925		IOP(I) = 0	DECLE195
1 L = 1 - 1	DECL7930		GD TO 22	DECLOZOD
INSTRA = INSTR(L)	DECL7935	1	2 IN = INSTR4 = 10	DEC132.2
IN = NJD(INSTR4J100)	DECL7940	r	CD ID (13)14)10)10)101191392092109 IM	DEC13215
IF(IN .NE, C) 60 TO 9	DECL 7945 C	ř	ONE DE THE EDUIDAING EUNCTIONS HAS BEEN ENCOUNTERED	CECL3220
IN = INSTR4 / 100	DEC17955 C	č	and of the runchowing restrictions when being the second to	DEC18225
	DECL7960	- 1	3 F(I) = SIN(F(I))	DECLEZIO
	DECL7965		GO TO 22	DECLEZES
THE FUNCTION 'X' ORIGINATES A NEW TERM	DECL7970	14	4 F(I) = COS(F(I))	DECLEZAC DECLEZAC
	DECL7975		GO TO 22 E E(T) EXPLETIN	DECTORES
F(I) = X	DEC17980	1:	5 F(1/ = EAF(F(1)) Fo To 22	CECL0255
10P(1) = 0	DECL 7982	1.	6 F(T) = A(GG(F(T)))	DECLEZED
	DECL7995	•	GO TO 22	DEC15255
3 10 = LEVEL(IPARA) + 1	DECLBOCO	1	7 F(I) = FN1(F(I))	DECL8273
A CLOSING PARENTHESIS MEANS THAT THE FUNCTION WITHIN NEEDS TO	DECL8005		GO TO 22	DECL5275
BE EVALUATED	DECL8010	1	8 F(I) = FN2(F(I))	SECL3210 SECL3235
	DECL8015	1	00 10 22 0 E(I) - EN3(E(I))	DFC15290
IF(IN .EQ. I) GD TO 22	DECL8020	1	GETTE 22	DECLEZ75
	DEC18020	20	O(F(I) = FN4(F(I))	DECLABCO
00 2 K = 1Ny1 15/120/00 E0 00 CO TO 4	DECL8035	_	GD TD 22	02018309
F(TN) = F(X + 1) + TOP(K) + F(TN)	DECL5040	2	1 F(1) = FN5(F(1))	DECLABID
	DECL8045	22	2 IF(L .GT. LOC2) GO TO 1	DECLE315
$4 F(IN) = F(K + 1) \neq F(IN)$	DECL8050 C	C C	TO THE FOR ON FOUND TO TOCOMPETION THE VALUE OF THE SUNCTION	DEC16325
5 CONTINUE	DECLBOSS C	r r	IP L IS LESS OR EQUAL TO LOCZYRETORIA THE VALUE OF THE FORCITON	DECLE330
I = IN	DECLOUOU C	•	FCH = F(1)	DEC16375
10-11) = 0 10-24 - 10-24 - 1	DECI 8070		IF(I EQ. 1) RETURN	DECLE340
тедед = тедед = 1 Бр ТП 22	DECL9075		DD 23 L = 2 I	DECL83-5
6 IPARA = IPARA + 1	DECLADEO	23	3 FCN = F(L) * FCN	DECLASSO
	DECLBOBS		RETURN	DECL6355
OPENING PARENTESIS ARE NOTED	DECL8090		CNU	05063300
	0ECL5092			
LEVEL(IPARA) = I	DECL0100			
60 (8 <u>22</u> T 100/11 - 1	DEC18110			
F 10211 - 1 F9 T0 22	DECE8115			
3 IBP(I) = - 1	DECL8120			N
• · · · · • •				

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arithm included in the program, the numbers 15 to 19 are reserved for the auxiliary functions FN1 to FN5. These auxiliary functions allow users of FCN to specify any desired function by supplying separate Fortran function subprograms called FN1 - FN5 to generate them. In this way, the statements in FCN never need to be altered.

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Finally, the full matrix equivalent of an ordinary differential equation can be obtained by adding together the set of simple matrices generated by OPGEN2 and setting the result equal to the forcing function values at the interpolation nodes. This operation is performed by the subroutine DIFFEQ given in Figure 4.2.3. DIFFEQ is a master subroutine that requires as input only the correct values in the instruction code, plus a few other constants, and generates by calling OPGEN1 and OPGEN2 the matrix equivalent of any differential equation. Some additional features of DIFFEQ, which have not been theoretically explained at this point, will be described in the next chapter.

The use and operation of the automatic discretization subroutine package will now be illustrated by a simple example. Consider the differential equation

$$\frac{d^3y}{dx^3} - \sin(x) y = \cos(1.2 + \sin(x)) \quad (4.2.1)$$

where y is to be approximated by a sixth order polynomial in the interval 0 to 1. The instruction code corresponding to (4.2.1) is

Figure 4.2.3

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SUBROW INC. DISCEDION MONANT, NELMIS, INSTR. CONST. YD.NR.ND.	DFCL2810			IF(LOC(L), LT, 0) $IJ = MSIZE$	DECL3080
SUBRIDE THE STREETEN MONTH AND THE FILST THE THE DATE TO ADD THE THE STREETEN AND THE STREE	DEC1 2815			n 4 7 = 1 MeIZE	0=0L3035
# NULLD, NULLV, LIMDA)	02022012				- 100: 1121
	DECL2020	-			
DIFFED DETERMINES THE MATRIX EQUIVALENT OF AN ARBITRARY	DECL2825	Ç			
DESTINANT STREETENTIAL EQUATION AND RETURNS INELMTS! SECTIONAL	DECL2830	C		ADD THE SIMPLE OPERATORS IN A COMPOUND OPERATOR TUDETHER	
THE DISCHARTER THE DISCEPTION IS DEFINED IN THE	DECL2835	С			08113115
CELEVAL S 10 1 1354 DE DIFFERENTIAL ENDITION IS DE TRED IN THE	DEC1 2840		4	$\partial(TJ) = D(TJ) + \partial P(T)$	DECL3110
PYYAA (IU2 K, 743 (HE SOCOTION)2 KETOKWED IN (DA).	05012845		-		0=013115
					F2013173
DIMENSION (1) / WT(1) / INSTR(1) / CONST(1) / XD(1)	DECESSO			IP(LAMDA) GD IN 990	
DITERSION (300), IROW(10), JCCL(30), LCC(100), LPOWER(100), LROW(100)),DEC1.2852			1F(NR .EQ. 1) GU TO D	
	DECL2860			STEP = $1 \cdot 0/(FLDAT(NR - 1) \neq WT(N))$	DECE3130
	DECI 2865			X = XD(NPT) - STEP	05013135
E LA AGRA DE LA POLOCE COLE LA ORCERE LA CAPERE DE LA CAP	02012020				n=r. =
LCCICAL 14 74			-		
	JECE28/2		5		
DETERMINE THE CODRDINATE-INDEPENDENT FACTORS IN THE MATRIX	DECLERSO			X = XD(NPT)	32412.24
	DECL2885	С			CECT3125
	DECL2890	С		THE MATRIX IS AUGMENTED WITH THE SUM OF THE FORCING FUNCTION	DECL3160
	DEC1 2895	ē.		POINT VALUES	06013165
CALL DPGENI(INGING LOUPLOUPLOUPLOUPLOPSERDADE TO CONTIN	02012000	ř		- DIVE TATOLD	751212170
IFLE .LE. C .CR. LO .GT. 100 .DR. LF .GF. 100) GD (D 990	DECEZOU	C			
	DECL2905		5	$PU \neq L = 1 + L + 2$	
DEFERMINE THE MATRIX SIZE PARAMETERS	DEC1.2910			LGCL = IABS(LFCN(L))	
	DECL2915			LBCL1 = LFCN(L + 1)	05013163
	DEC1 2920		1	WEIGHT = FICAT/IDCI)/FIDAT/IFCN(I))	DEC13190
(ab) = (ab) + 1	00000000			Yo = V	
ASIZE = AF=ND					02013000
PBLANK = PSIZE + NR	DECE2930			D = 1 = 100 R	
	DECL2935			XR = XR + STEP	OFCENSIS
	DECL2940			IJ = I + MSIZE	DECLIZIO
(a) graduate and (a) of the second s Second second sec	DEC1.2945		7 1	D(TJ) = D(TJ) + FCN(XR)TNSTR,CDNS ^T ,LDCL/LDCL1)#WEIGHT	0ECL3215
	DEC1 2950	r	•		0FCL3222
1+(+W1++ L+*D+) 60 10 Z	05012055	ř		NETERMINE THE CENERAL SOLUTION IN EACH CLEMENT	05113225
$MB_ANK = 2 + MSIZE$		2		DETERMINE THE GENERAL SUCCION IN EACH ECCAENT	
50°. = 2480	DECEZANO	C			
NTEFAL = 2 = NO	DECL2965		8 1	LALL NULL(BANKANDADTUTALAIKAIKAIRUMAJCOLADNAMDGANULLVAMADDON)	D111224
	DECL2970			IF(IRANK .N.F. MR) GO TO 991	SEC132-C
	DEC1.2975		i	MADDDN = MADDDN + NDSIZE	DECL3345
	DFC1.2980	1	0.0	CONTINUE	DFC13250
	DECI 2085	-			DEC: 3255
MADDDN = C		r			020122240
	DECLZONO	č			32012222
TAKE EACH ELEMENT SEPARATELY	DECL2995	C	-	ERRORS ENCOUNTERED	12013222
	DEC1.3000	C			02013270
And the second	DECL3005	99	10 i	NRITE(5)200) LAMD4,LD,LF	DECL3275
	DEC1 3010		Č ;	0 = 0	07014730
$(12)^2 = (12)^2 + (12)^2$	5761 7015			SE THON	7261 3725
UD B I = 1,78LANK	02010010	0.0	. !		02020200
3 D(I) = 0.	DECESOZO	79	1	WRITE (07201) TRANKINK	551132-5
LEPS = + LRE#(1)#ND	DECL3025			4D = 0	0ECL3243
	DEC1.3030		F	KETURN	DECL3912
TAKE STOP STARS & DOCRATOR SEDARATELY	DECL3035	20	0 1	FORMAT(I-SORRY, THIS EQUATION CANNOT BE SOLVED BY THIS PROGRAY',	DECUERTS
TAKE SACH SIFFLE DESKITCK DEFANATELT	DEC1 3040	- •	*	315	52613210
	55013045	20	۰. ۲	EDAMAT, LETHE CALCHEATED RANKING & DOCE NOT COULD THE THE DANKE.	72012211
DO 4 L = 1,10,2	12013043	20	1 1	BRHAT (1-THE CALCULATED RANK 915) DUES NOT EVOLUTIE THE THEE RANK	
	DECL3050		*	13)	35013325
200 THE COORDINATE-DEPENDENT FACTORS	DECL3055		5	END	DEC13325
	DECL3060				
CALL PRESIDENTS, CONST. LOC. PRWER. LOCAL MON. N. N. WT. N.NPT.	DECI.3065				
URLE GEGENZIJASKJUMSKJEGUJE GMENJEKONJENNOVANJENNA V	05013070				
	DECL3075				7
IJ = 0	DECEDUID				4

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<u>I</u>	<u>INSTR (I)</u>	<u> </u>	<u>INSTR (I)</u>
1	1000	11	12
2	- 3	12	300
3	200	13	-101
4	2000	14	500
5	11	15	11
6	300	16	300
7	100	17	100
8	400	18	400
9	200	19	400
10	2000	20	3000

In the above instructions, the function on the right hand side of eqution (4.2.1) has been transferred to the left hand side. Also, the constant 1.2 of instruction 13 is to be stored in CONST (1).

Given these instructions, the subroutine OPGEN1 performs the following operations. It detects the four simple operator separation locations - numbers 1, 4, 10 and 20 - and determines, by looking for the number 200, that in the first and second group of instructions (numbers 1 to 3 and 4 to 9), the unknown y appears but that it is absent in the last group (numbers 10 to 19). This information is stored in the arrays LOC and LFCN as follows

LOC (1) = 1LOC (3) = 4LFCN (1) = 10LOC (2) = 3LOC (4) = 9LFCN (2) = 19

Consequently, coordinate-independent matrix operators need to be ⁷⁶ evaluated only with the first two groups of instructions. Since there is a -3 in the first group and no negative numbers whatsoever in the second, and since sixth order polynomials are to be used, the values of NR and ND become 4 and 7, respectively. Furthermore, the first simple operator stored in OPS is the 4 x 7 projection matrix.

Upon completion of these calculations, these numbers are returned to DIFFEQ which then calls OPGEN2 twice, first with instruction numbers 1 to 3 and then with instruction numbers 4 to 9. Since there are no functions in instructions 1 to 3, OPGEN2 only multiplies the elements of the first simple matrix by $(1/L)^3$ with L=1.0. However, in the next case, according to statements 5 to 8, the rows of the second simple matrix are multiplied by the value of sin(x) at the four points $x_i = i/3$, i=0,1,2,3.

As these two 4 x 7 matrices are produced, they are added together in the subroutine DIFFEQ and stored in the array D. Following this, the function in instructions 10 to 19 is augmented to the matrix D calling FCN at the aforementioned four points x_{i} .

This problem was run on an IBM 360/75 computer to determine an estimate of the computation times required by the programs. The generation of the 4 x 7 third order differentiation and projection matrices by the subroutine OPRATR required 16 Including this time, the complete determination of msec each. the coordinate-independent components by the subroutine OPGEN1 took 33 msec. In performing its computations for the single region considered above, OPGEN2 used 16 msec. Altogether

the complete discretization process, including all of the relevant operations performed by DIFFEQ, required 67 msec.

4.3 TWO-DIMENSIONAL EQUATIONS

Considering the physical differences between ordinary differential equations and two-dimensional partial differential equations, the formation of matrix equivalents for the two types of problems is remarkably similar. As explained in section 3.5, in a product separable space, the matrix equivalent of a twodimensional coordinate-independent differential operator factors into two one-dimensional matrix equivalents, one of which multiplies the matrix of unknown coefficients from the right and the other multiplies it transposed from the left. As a consequence of this reliance upon one-dimensional matrices, computer programs for the discretization of partial differential equations by the methods of Chapter 3 can be constructed by properly adapting the one-dimensional programs given in the preceding section.

In fact, the generation of the high-order one-dimensional differentiation and projection matrices needed for twodimensional partial differential equations can be performed without modification by the subroutine OPRATR in Figure 4.1.2. The essential difference in the two-dimensional case is that OPRATR needs to be called twice, once for the x-dependent operators and once for the y-dependent operators. A subroutine, OPG2D1, which has been designed to perform this operation for any linear two-dimensional partial differential equation, is given in Figure 4.3.1. OPG2D1 is very similar in form to OPGEN1 and, as a result, large blocks of statements which are identical to the original one-dimensional program have not been repeated. In their place, comment statements have been inserted to indicate the statement numbers deleted. Thus the comment between statement numbers 3175 and 3185 of OPG2D1 indicates that Fortran statements 6715 through 6860 of OPGEN1 should appear in that location. This method of shortening program listings will be used wherever possible in this thesis.

OPG2D1 requires the following modification of the instruction code given in sections 4.1 and 4.2

200	Z (the	unknown function)
101	Y (one	of the independent variables)
100	X (the	other independent variable)
9~1	D9/DY9	~ D/DY
-1~-9	D/DX	~ D9/DX9

Given a set of instructions according to this code for a partial differential equation, OPG2D1 determines the number of x-dependent differential operators and the number of y-dependent differential operators in each term of the equation. It then evaluates the corresponding matrix equivalents and applies the proper projection matrices to both the x-dependent and the y-dependent operators to keep the polynomial orders of their spaces consistent.

Figure 4.3.1

SUBKOUTINE DPG2D1(INSTR,ISTART,NRX,NRY,ND,LDC,LD,LFCN,LF, # DPS,LRCW,LPDMER,DPX,DPY,IPROJ)	DE2D3145 DE2D3150	CALL DPRATP(DPX)IEDX;MX;MX;NX;0) MXX = MX	DE2D3415 DE2D3420 DE2D3425
	DE2D3155 1	$2 IF(MY \cdot EQ \cdot NPY) GU TU 13$	05203420
THIS SUBROUTINE EVALUATES THE COURDINATE-INDEPENDENT FACTURS	05203165	CALL COPATE/COV. TERV. MY. MNV. NY. CA	DF 203435
FIR THE DISCRETIZATION OF LINEAR THUSDIMENSIONAL PARTIAL	05203130	MNV - MV	DE2D3440
DIFFERENTIAL EQUATIONS.	DE2D3175 1		DE203445
DECL 4715 6960	05203170 1	1 = 1 ABS(1) STP(1) SCE(10) GUTD 15	DE703450
DECE BUIDEDADU	DE203135 C		DE203455
	DE2D3190 C	ADD THE DIFFERENTIATION MATRICES	DE203460
	DE203195 C		DE2D3465
	DF2D3200	IF(INSTR(I), GE, O) GO TU 14	DE203470
	DF2D3205	LPOWER(LU) = LPOWER(LD) = INSTR(L)	DE203475
	DE2D3210	MX = MHX + INSTR(L)	DE2D3480
	DE203215	CALL OPRATR (OPX, IEOX, MX, MX, NX, 1)	DE203485
	DE203220	MNX = 11X	DE2D3490
LPGWER(101) = 0	DE2D3225	GU TO 13	DE203495
	DE203230 1	4 LPOWER(LO1) = LPOWER(LO1) + INSTR(L)	DE203500
ty a My	DE203235	MY = MHY - JHSTR(L)	DE203505
MTYMSM a NX ## 2	DE201240	CALL OPRATH (OPY) IEOY/MY/MY/NY/1)	DE2D3510
	DE203245	MNY = MY	DE203515
CREAT TWO UNIT MATRICES	DE203250	GU TO 13	DE203520
	0E2D3255 1	5 IF(INSTR(L) .GE, 1000 .OR. IABS(INSTR(L)) .LT. 10) GO TO 16	DE203525
$DD \rightarrow I = 1.000$	DE203260		DE203530
5 UPx(1) = 0.	DE203265	GO TO 15	DE293535
MNX = MX + 1	DE2D3270 C	-	DE203540
DD 6 I = 1, HTYKSNAMNX	DE203275 C	BUTH MATRICES ARE STURED IN IUPSI	DE203545
6 UPX(I) = 1, = SIGN	DE203280 C		DE203550
$MTYHSH = HY \neq 2$	DE203285 1	6 L = L + 1	DE203555
00 7 I = 1,MTYHŜN	DE203290	LGC(LO) = IDP + LOC(LO)	DE203360
7 CPY(1) = 0	DE203295	LRCW(LO) = 8X	DE203565
$M_{\rm A} = M_{\rm A} + 1$	DE203300	LROW(LG1) = MY	DEZD3570
$DD = R = T^{h} H L A N 2 H^{h} H A$	DE2D3355	MTYMSI: = NX + NX	DE203572
B @PY(I) = 1.	DE273310	IF(LO .EQ. 1) LOPS = O	DE203583
IEOX = 1	DE203315	DO 17 I = 1, MTYMSN	DESDASBO
153Y = 1	DE203320 1	7 OPS(I + LOPS) = OPX(I)	0E20359V
MTeX = 14X	DE203325	LOPS = LOPS + MIYISN	DE203292
WMA = HA	DE203330	MTYASU = HY * NY	05703600
IF(INSTR(L) .NE. 200) GD TO 13	DE203335	DO 18 I = $1.MTYMSN$	06203602
I = L2CS	DE203340 1	B UPS(I + LDPS) = PPY(I)	05202010
9 I = I = 1	02203343	LUPS = LUPS + MIYASN	05213012
IF(INSTR(I), EC, 300) IDP = = 1			05203620
1F(IPR0J .4E. 1) 60 10 11		5561 9156-7178	DE2"3630
IF(IABS(INSTR(I)),GE, 10) GO TO 11			DE203635
1F(INSTR(I) .GE, 0) 62 TU 10	DE203302 C		DE203640
PX = "X + INSTR(I)	05203375	LIND	5550000
. 60 13 I T	05203330		
	05203360		
AUD THE PROJECTION MATKIVES	DE203302		
1. HU - HU - 11.FTF/T)	DE2D3395		
10 PT = 0T = 10018(1) 11 TELINETULTA AT 10000 CU TO 9	DE203400		
11 12(10)(11) *2(* 100) 04 10 2	DE203405		\sim
3510A 8548 0KAT 00 10 35 00 0 3.57 4 55 - 07	nE203410		õ
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The coordinate-dependent factors of two-dimensional equations are added to the system by the subroutine OPG2D2 in Figure 4.3.2. OPG2D2 is patterned after the one-dimensional program OPGEN2 and, again, many statements in OPG2D2 are similar to those in OPGEN2. The basic difference between the two programs is that OPG2D2 must convert the pair of one-dimensional matrices obtained from OPG2D1 for each differential operator into a single coefficient matrix by taking the Kronecker product of the matrices. This operation is performed by the subroutine KRON given in Figure 4.3.3.

The major new problem encountered when changing from one to two-dimensional differential equations is the rapid growth in the size of the matrix equivalents of the operators. In the one-dimensional program, the largest possible matrix equivalent of a differential operator results when a ninth order polynomial approximation is made of a first order differential operator and produces a matrix with 90 elements. With two-dimensional problems, the Kronecker product of two of these one-dimensional matrices results in a matrix with 90^2 = 8100 elements. Although matrices of this size are within the limits of the storage capacities of modern computers, they require a sizable fraction of core memory and operations with them consume considerable computational time. Therefore, the generation of very high-order matrix equivalents for twodimensional problems requires much more consideration than is necessary in the one-dimensional case. Since the number of matrix elements in a two-dimensional problem varies approximately as n^4 where n is the order of the approximating polynomial, it is often wise to use a lower order element,

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	SUBKOUTINE OPG202(INSTR,CONST,LOC/LPOWER,LROW,L,ND,XD,YD,WTX,WTY * N,NPT,OPS,LOPX,OP,TEMP)	DE203645		00 9 I = 1, MX X = X + STEPX	DE203915
	THE SUBROUTINE OPG232 TAKES THE COMPANYATE AND COMPANY	DE2D3655		Y = YS	DE2D3920
	A TWO-DINEYSTONAL PARTIAL DIEEEPENTIAL CONSTRATE-INDEPENDENT FACTORS O	F DE2D3660		DD 9 K = 1 MY	DE203925
	THEM WITH THE APPROPRIATE COOPERATE CONTINUE AND MULTIPLES	DE2D3665		Y = Y = STEPY	DE2D3930
		DE2D3670	C		DEZU3935
	DOUBLE PRECISION DOUBLE	DE2D3675	C	EVALUATE THE FUNCTION AT THE NODAL POINT	DE203940
	DIMENSION INSTR(1), CONST(1), COC(1), CONST(1), CONST(1)	DE2D3680	C		02203943
	* XD(1)/YD(1)/WTX(1)/WTY(1), GPS(1), GP(1), TEUR(1)/	DE2D3685		FUNCT = FCN2D(X/Y/INSTR/CONST/LOCL/LOC(1) = WEIGHT	JE203950
	MNX = ND	DE203690		IK = IK + 1	05203972
	HNY = ND	DE203692		IJ = IK = MCP	DE203950
	<i>i</i> l] = L + 1	06203700		KJ = IJ	DE203930
	MX = LRGW(L)	05203703		IF(LOC(L1) .GE. 0) KJ = IJ + MOP	DE203075
	MY = LROW(L1)	06209710	C A		DE2D3930
	PDAX = PDAX + MX + MNX	DE203715	ç	THE COMPUTED MATRIX IS STORED IN IDPI	DE20235
	IF(MNX .EQ. ND .AND. MNY .EQ. ND) GD TD 4	DE203720	Ç		06203990
		06203725		DD A 1 = 1'VOb	DE203990
	EVALUATE THE NEXT MATRIX OPERATOR	02203730		J = IJ + nDP	DEZDADOD
		05202730		КЈ = КЈ + МДР	DE204000
	CALL KRON(CPS, MX, MNX, LOPX, OPS, MY, MNY, LOPY, TEMP, MOP, MNOP, O)	06203740		9 UP(KJ) = FUNCT + OP(IJ)	DE204010
	13 = 0	06203760		IF(LUC(L1) .LT. O) RETURN	DE204015
		06203765			DE204020
		06203750			DE204025
	NJ = KJ + MNDP	DE203765			DE204030
		DE203770			DE204035
	1J = 1J + 1	DE203775		END	DE204040
		DE203780			
		DF2D3785			
•		DE203790			
2	DEGLE = DOUGLE + DBLE(TEMP(IK) + DP(K + KJ))	DE203795			
_		DE203800			
6		DE2D3805			
10	LDPX = / PPY + Y + WHY	DE203810			
		DE203815			
		DE203820			
		DE203825			
	THE NODAL POINTS ARE STERRED OFF	DE2D3830			
		DE203835			
	HEIGHT = (VTX(h) ++ IPAVER(L) + (VTV(h) ++ IPAVER(L))	DE203840			
	IF(MX .EQ. 1) GD TO S	DE2D3845			
	STEPX = 1.0 / ((FX = 1) + wTY/WAA	DE203850			
	X = XD(IPT) = STEPY	DE2D3855			
	GG TO 6	DE2D3860			
- 5	STEPX = 0.5 / CTX(1)	DE2D3855			
	X = XD(NPT)	DE2D3870			
6	IF(XY +EQ. 1) GD TO 7	DE203875			
	STEPY = $1 + G \neq ((+Y - 1) + U + V + (h))$	DE203880			
	YS = YD(NPT) + STEPY	DE203385			
	60 TC 8	DE203890			
7	STEPY = 0.5 / TTY(N)	UE203895			
	YS = YD(1 PT)	DE203900			
ę	IK = 0	UEX03405			
		DF503410			œ

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Figure 4.3.3

	SUBROUTINE KRON(A,MA,NA,IAADD,B,MB,NB,IBADD,C,MC,NC,ICADD) KRON TAKES THE KRONECKER PRODUCT OF THE IMAI BY INAI MATRIX IAI AND THE IMAI BY INBI MATRIX IBI AND RETURNS THE RESULTING IMC! BY INCI MATRIX IN ICI.	DE2D4045 DE2D4050 DE2D4055 DE2D4060 DE2D4065
	DIMENSION $A(1)_{B}(1)_{C}(1)$ MC = MA + MB NC = NA + NB IJC = ICADD	DE2D4075 DE2D4030 DE2D4035 DE2D4035
	IBA = IBADD + MB * NB + MB + 1 $JACUL = IAADD = MA$ $DU 1 JA = 1 MA$ $JACUL = JACGL + MA$	DE2D4095 DE2D4100 DE2D4105 DE2D4105 DE2D4110
	JBCUL = IBA DD 1 JR = 1≠NB JBCUL = JRCUL = MB IJA = JACUL	DE2D4115 DE2D4120 DE2D4125 DE2D4125 DE2D4130
	$D_{\Box} \ I \ I \ A = I_{J} M A$ $IJA = IJA + 1$ $D_{\Box} \ I \ IB = I_{J} M B$ $IJC = IJC + 1$	DE2D4135 DE2D4140 DE2D4145 DE2D4150
1	C(IJC) = A(IJA) * B(JBCDL - IB) Return End	DE2D4155 DE2D4160 DE2D4165

Figure 4.3.4

	FUNGTION FCN2D(X,Y,INSTR/CONST/LUCÍ/COC2)	DE2D4170 DE2D4175
	THIS FUNCTION SUBPROGRAM RETURNS THE VALUE OF A TWO-DIMENSIONAL Function stored in Lucations 'Loc2' to 'Loc1' of the Array	DE2D4190 DE2D4185 DE2D4190
C	DECL 7890-7940	DE204195 DE204200
č	IF(INSTR4) ,EQ. 101) GO TO 24	DE204205 DE204210
CC	DECL 7945-7995	DE204219 DE204220 DE204225
C	24 I = I + 1 F(I) = Y	DE2D4230 DE2D4235
	IOP(I) = 0 $GO TO 22$	DE2D4240 DE2D4245
CC	DECL 6000-8360	DE2D4250 DE2D4255 DE2D4260
	END	DE2D4265

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particularly if several different matrix equivalents are needed simultaneously.

After converting the matrix of unknown coefficients into vector form by applying the Kronecker product, OPG2D2 multiplies the resulting product matrices by any necessary function matrices, as specified by equation (3.5.30). The function values are obtained from the function subprogram FCN2D given in Figure 4.3.4. FCN2D is similar to FCN except that it has two independent variables and calls auxiliary functions FN2D1 - FN2D5, which are dependent on both X and Y.

Last, Figure 4.3.5 contains DIFF2D, the two-dimensional version of the subroutine DIFFEQ. The principles of operation of DIFF2D are identical to those of DIFFEQ, although a large number of indexing changes have been made. By calling DIFF2D with a suitable set of instructions in the array INSTR, the matrix equivalent of any two-dimensional partial differential equation may be obtained.

Figure 4.3.5

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	<pre>SUBROUTINE DIFF2D(OH,HON,WTX,WTY,NELMTS,INSTR,CONST,XD,YD,</pre>	DE2D1385 DE2D1390 DE2D1490 DE2D1405 DE2D1405 DE2D1410 DE2D1415 DE2D1425 DE2D1425 DE2D1435 DE2D1435 DE2D1435 DE2D1440 DE2D1445 DE2D1440 DE2D1455 DE2D1450 DE2D1455 DE2D1460 DE2D1455 DE2D1460 DE2D1465 DE2D1465 DE2D1485 DE2D1485 DE2D1485 DE2D1495 DE2D1495 DE2D1495 DE2D1495 DE2D1515 DE2D1515 DE2D1535 DE2D1535 DE2D1545 DE2D1555 DE2	000 000	<pre>7 STEPY = 1.0 / WTY(N) YS = Y0(NPT) 8 DD 9 L = 1,LF,2 DECL 3180=3195 DD 9 I = 1,NRX X = X + STEPX Y = YS DD 9 J = 1,NRY Y = Y - STEPY i J = IJ + 1 9 D(IJ) = D(IJ) + FCN2D(X,Y,INSTR,COHST,LOCL,LOCLI) * WEIGHT 11 CONTINUE CALL HULL(0,NR2D,ND2D,NTDTAL,IPANK,IROW,JCOL,DN,HDN,NULLV,HADDDN) IF(IRANK .NE. NR2D) GD TU 991 DECL 3245=3320 END</pre>	DE2D1665 DE2D1660 DE2C1665 DE2D1670 DE2D1685 DE2D1680 DE2D1685 DE2D1700 DE2D1700 DE2D1705 DE2D1705 DE2D1725 DE2D1725 DE2D1735 DE2D1735 DE2D1745 DE2D1756
	CECL 3075-3110	DE2D1575 DE2D1580 DE2D1585			
·	<pre>IF(LF .EQ. 0) GD TO 11 IF(LAMDA) GD TC 990 IF(NRX .EQ. 1) GD TO 5 STEPX = 1.0 / ((NRX = 1) * WTX(N)) X = XD(NPT) = STEPX GD .D 6 5 STEPX = 0.5 / NTX(N) X = XD(NPT) 6 IF(NRY .EQ. 1) GD TD 7 STEPY = 1.0 / (('NRY - 1) * WTY(N)) YS = YD(NPT) + STEPY GD TD 8</pre>	CE201590 DE201595 DE201605 DE201605 DE201610 DE201610 DE201615 DE201625 DE201625 DE201635 DE201635 DE201645 DE201645 DE201650			
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CHAPTER V

SOLUTION TECHNIQUES

There are two possible approaches for solving the rectangular matrix equations which result from the discretization of differential equations by the projective method in Chapter 3. In one, boundary conditions are added to the system until the coefficient matrix becomes non-singular. This matrix equation may then be solved using standard methods to yield an approximate particular solution of the differential equation. In the other, the matrix equation is solved by using the matrix generalized inverse concept to determine the approximate general solution of the differential equation. Boundary conditions are added subsequently if a particular solution is desired. Since the latter method is applicable to a much wider range of problems than the first, it is the approach followed in this thesis.

The general solutions which result from the application of matrix generalized inversion to the singular matrix equations of Chapter 3 are completely analogous to the general solutions encountered classically in the analytic theory of differential equations. In both cases, the general solution of an operator equation is given by an inhomogeneous solution and a set of homogeneous solutions which form a basis in the nullspace of the operator. Furthermore, the procedures used to obtain both types of general solutions are independent from the addition of boundary conditions. The numerical method developed in this thesis is unique in this feature among all of the available approximate solution techniques for the solution of differential equations.

In addition to the ability to produce approximate general solutions, a further advantage of applying matrix generalized inversion to the projective solution of differential equations is the highly efficient matrix solution algorithm that results. This efficiency is derived from substituting manipulations on several small full matrices in multi-regional solutions for the more costly treatment of one large sparse matrix.

The only published work which uses matrix generalized inversion with a finite element type of formulation is a paper by Berkovic [50]. However, since this paper does not deal with differential equations but rather the algebraic Hooke's law equation, the extremely beneficial results which follow from its application to differential equations are not evident there.

5.1 GENERALIZED MATRIX INVERSION

Let A be an $m \times n$ matrix of rank r and suppose that all vectors y satisfying the equation

Ay = f (5.1.1)

are desired. The simplest method of determing these vectors is to decompose A by the Gaussian elimination process, interchanging rows and columns as needed, until only zero pivots remain [48,

51]. Then A may be factored as

$$A = LU = \begin{pmatrix} L_1 \\ L_2 \end{pmatrix} \begin{bmatrix} U_1 & U_2 \end{bmatrix}$$
(5.1.2)

where L_1 and U_1 are non-singular. With this factorization (5.1.1) may be written

$$y_1 = U_1^{-1} L_1^{-1} f_1 - U_1^{-1} U_2 y_2$$
 (5.1.3a)

$$f_2 = L_2 L_1^{-1} f_1$$
 (5.1.3b)

Equation (5.1.3a) contains all of the useful information from (5.1.1); (5.1.3b) merely indicates that solution is possible only if no conflicts arise. Consequently, (5.1.3a) together with the equation $y_2 = y_2$ gives

$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} U_1 & L_1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} + \begin{pmatrix} 0 & -U_1 & U_2 \\ 0 & I \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} (5.1.4)$$

Since the columns of the second matrix in (5.1.4) form (n-r) independent nullvectors of the matrix A, a general solution of (5.1.1) is given by [52]

$$y = A^{+} f + Nz$$
 (5.1.5)

where A^+ and N are the row and column reinterchanged forms of

$$A^{+} = \begin{pmatrix} U_{1}^{-1} & L_{1}^{-1} & 0 \\ 0 & 0 \end{pmatrix}$$
(5.1.6)

$$N = \begin{pmatrix} -U_1 & U_2 \\ & \\ I \end{pmatrix}$$
(5.1.7)

and the (n-r) components of z are arbitrary.

The matrix A^{\dagger} is called a weak generalized inverse of A [53] and satisfies the following two properties

$$AA^{+}A = A$$
 (5.1.8a)

$$A^{+}AA^{+} = A^{+}$$
 (5.1.8b)

However, the products

 $AA^{+} \stackrel{\cdot}{=} \begin{pmatrix} \mathbf{I} & 0 \\ -1 \\ L_{2} & L_{1} & 0 \end{pmatrix}$ (5.1.9a) $A^{+}A = \begin{pmatrix} \mathbf{I} & U_{1}^{-1} & U \\ 0 & 0 \end{pmatrix}$ (5.1.9b)

are not symmetric in general.

The significance of equation (5.1.3b) may be understood by considering the vector norm |.| of the residual vector

 $|Ay - f| = |AA^{\dagger}f + ANz - f|$

$$= | (AA^{\ddagger} - I) f |$$

$$= | \begin{pmatrix} 0 & 0 \\ -1 & -1 \\ L_2 & L_1 & -I \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} |$$

$$= | \begin{pmatrix} 0 & \\ -1 & \\ L_2 & L_1 & f_1 & -f_2 \end{pmatrix} |$$
(5.1.10)

Therefore, if (5.1.3b) holds, the solution y will be exact; if (5.1.3b) is not true, y will satisfy (5.1.1) only approximately, the amount of error being dependent on the values of f_1 and f_2 .

In most physical problems that are properly formulated, $f \in \text{Range}(a)$ so that (5.1.3b) will be satisfied. However, if this is not true, a more accurate approximate solution to (5.1.1) is provided by the Moore-Penrose inverse of a matrix [54] which minimizes the norm (5.1.10) in an L₂ sense. The Moore-Penrose inverse is related to the weak generalized inverse (5.1.6) by [55]

$$A^{MP} = A^{*} (A^{*} AA^{*})^{+} A^{*}$$
(5.1.11)

where A* is the transpose conjugate of A.

It is also of interest to determine the solution of a system of equations part of which has already been solved in the form (5.1.5). Suppose, for example, that the system of equations (5.1.1) is augmented by another matrix equation

(5.1.12)

By = g

where B is an $m_b \ge n$ matrix of rank r_b . Inserting (5.1.5) into (5.1.12) gives

$$BA + BNz = g$$
 (5.1.13)

The solution of this equation is

$$z = (BN)^{+} (g - BA^{+}f) + Mu$$
 (5.1.14)

where the $(n - r) \times (n - r - r_b)$ matrix M is in the nullspace of the matrix (BN) and the vector u is arbitrary. It follows that the solution of (5.1.1) together with (5.1.12) is

$$y = (A^{+} - N(BN)^{+} BA^{+})f + N(BN)^{+}g + NMu$$
 (5.1.15)

It may be observed from (5.1.15) that each time an independent equation is added to the system, the solution contains one less unknown. Initially, when the *n* dimensional space of the vector y is created, all of its components are arbitrary. Then as conditions are imposed in the form (5.1.1), unwanted vector behavior is filtered out, leaving a solution of the kind (5.1.5). As more equations are added, fewer degrees of freedom remain, until, when *n* linearly independent conditions are imposed, only one particular vector survives.

In actual computation of the general solution of the matrix equation (5.1.5), the most economical procedure is to augment the matrix A with the forcing function f to form the new matrix equation [56]

Figure 5.1.la

				<u> </u>
4	. A(IJ) = STERE			9
	$\Delta(1) = \Delta(1)$	DECL8630	12 = 11 - 11	0201P901
	STORE = A(I)	DECL8625	DD 14 I = 1 IRANK	02010877
	IJ = IJ + N	DECL8620 C		0201724
	I = I + M	DECL8615 C	NOW THE BACK SUBSTITUTION STEP IS PERFORMED	DELLARIA DECLARET
	$DO 4 JJ = 1_{g}NTOTAL$	DECL8610 C		UECLOBEL DECLOBEL
	IJ = IR - H	DECL8605	II = IRANK#M + JI	BECL5873
	I = IRANK - H	DECLB600	JI = IRANK + 1	DECESSIO
		DECL8595	10 IF(IRANK .EQ. O .OR. IRANK .GE. NTOTAL) RETURN	DECLIÈCE
	LARGEST ELEMENT	DECL8590	9 CONTINUE	DELL-504
	THE IRANKITH ROW IS INTERCHANDED WITH THE RUN CONTAINING THE	DEC1.8585	8 CONTINUE	DECLERD2 DECLERD2
	THE ROW CONTAINING THE	DECL8580	JC = J	0EUL2824 p=r: p==
	IF(IR .EQ. IRANK) SO TU P	DECL8575	IR = I	
	IRANK = IRANK + 1	DECL8570	PIV = A(IJ)	DECLES
		DECL8565	IF(ABS(A(IJ)) .LE. ABS(PIV) .OR. J .GT. N) GO (O 8	DECLESS
	MATRIX IS UPDATED	DEC18560	A(IJ) = A(IJ) - FACTUR*A(JJ)	020000000
	IF AN ELEMENT LARGER THAN THEF HAS BEEN TOUTED THE THE THE	DECL8555	JJ = JJ + H	DECLARAD
	TO AN CLOWENT LARGER THAN ITOLI HAS BEEN FOUND, THE RANK OF THE	DECL8550	IJ = IJ + H	DECL8875
	TL(MD2/MIA) *FE* (PF) OD (D IA	DECL8545	U = J = II M = II M = AL	DECLORZO
	30 - 300 = 10	DECL8540	iJ = iZ	DECLERIS
	18208 - V 199 - Vent - 1.8	DECL8535	11 - 12 11 - 12	DECLEBIO
	2014 C - HTM 10-197 - D	DECL8530	IT(IKANK -52 N'U'AL/ 00 10 7	DECLESOS
	roll v lavenderstart Horzan – Mass	DECL852>	A(12) = FAULUK TETTATIVE OF TOTAL) OF TO 9	DECLESCO
3	TT: - : . 5F_7+ARS(PIV)	DECL8520	FALIUK # A(12)/3/UKE	DECLE795
-		DECL8515	12 = 12 + 1 Exclose _ A/101/STOPE	DECLE790
		DECL8510	To -	DECL8785
	F19 = 1	DECL8505 C	DD 9 1 - 1.M	DECL57E0
	- Al CHUQCHIAR FROM FROM FROM FROM FROM FROM FROM FRO	DECL8500 L	INE GAUSSIAN ECTNINGLION SIEL IS LEN ANNEL HER.	DECL8775
	(FTABS(A(I)) LE. ABS(PIV)) GO TO 3	UECL8493 C	THE CAUSSIAN FLIMINATION STEP IS PERFORMED NEXT	DECLATIO
		DEC10490	111 - 11	DECLE755
	00 3 I = 1.K	DECL8487	PTV = 0	DECLE750
	20 3 J = 1,N	DEC18489	STORE = PIV	DECLE755
		DECL0472	12 = 1PIV	DECLETEO
	THE LARGEST ELEMENT IN THE MATRIX IS LOCATED	DEC18475	$IPIV = M \neq (IRANK - 1) + IRANK$	DECL8745
		DECL0402	II = IRANK + I	DECL8740
	10 = 0	DECI 8465	7 IF (IRANK .GE. H) GO TO 10	DECLE735
-	PIV = 0	DECL8460	JCDL(IRANK) = IJ	DECLA730
2	JC21(J) = J	DECL8455	JCDL(JC) = JCDL(IRANK)	DECL: 720
-	DO 2 J = 1,NTOTAL	DECL8450	ij = j(ol(jC))	02011720
1	IROW(I) = I	D=CL8445	6 CONTINUE	5:0:013
	0 = 1 I = 1 H	DECL8440	A(IJ) = STURE	DECLETIO
		DECL8435	A(I) = A(IJ)	DECL9703
	THE REA AND COLUMN VECTORS ARE FIRST INITIALIZED	DECL8430	STORE = A(I)	DECENTED
		DEC1.8425	IJ = IJ + 1	05015572 05019700
	01MEN51ON_4(12);18CW(2);JCOL(6);4N(12)	DECI 8420	$D = 6 I = J_J J_J$	02203570
	DOUBLE PRECISION DOUBLE	DECL8415	$IJ = (JC - 1) \pm M$	22613652
		DECL8410	u = 1 + 4 - 1 - 1	
	ADRK LRRAVS DE DIMENSION 'M' AND 'N'.	DECL8405	$J = (IRANK - 1) \neq M + 1$	020100-2
	FORGING FUNCTION MATRIX AUGMENTED TO "A". "INUME AND "JUDLE AND	DECL8400 C		020132 0
	THE LAST '(NTETAL - IRANK)' COLUMNS OF 'AN', WHERE PE' IS A	DECL8395 C	THE LARGEST ELEMENT	05017502
	WATRIX IN THE ARRAY 'ANI. IT ALSO RETURNS THE PRODUCT (AFTE IN	DECL8390 C	THE IRANKITH COLUMN IS INTERCHANGED WITH THE COLUMN CONTAINING	02010000 020102445
	MY SY INI MATRIX 'A' AND STORES THE RESULT AS AN THAN BITTAN	DECL8385 C	-	0501653
	OLL COMPUTES A LINEARLY INDEPENDENT SET OF NULLVELTORS FOR THE	DECL8380	5 IF(JC .EQ. IRANK) GO TO 7	
	THE REPORT OF AND INCOME THE	DEC18375	IRDW(IR) = IJ	02013047
	SUBROUTINE NULL(A,M,N,NTOTAL,IRANK,IROW,JCOLJAN,MAN,MAN,MANDAM	DECLE370	IRGW(IRANK) = IRGW(IR)	05012540
	THE REPORT OF AN MAN MAN MADDAN	DECI 8365	IJ = IROW(IRANK)	DECL3635

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1PIV = 11 - 1 DO 13 J = J1,NTOTAL 12 = 12 + 11 JJ = I P I VLL = 12 DUJBLE = 0.00 IF(I .EQ. 1) GU TO 20 VO 12 K = 2,1 JJ = JJ → M LL = LL - 1 12 DOUBLE = DOUBLE - DBLE(A(JJ)*A(LL)) 20 JJ = JJ - HLL = LL - 113 A(LL) =(SNGL(DOUBLE) - A(LL))/A(JJ) 14 CONTINUE JJ = (IRANK - 1)*MLL = MADDAH - MAN C C C THE NULL VECTORS ARE TRANSFERED TO THE ARRAY 'AN', REINTERCHARGING COLUMNS IN THE PROCESS С DO 17 J = JI.NTOTAL JJ = JJ + HLL = LL + MAN 00 17 I = 1, MAM IJ = JCOL(T) + LLIF(I .GT. IRANK) GD TO 15 (LL + I)A = (LI)MAGO TO 17 15 IF(I .EQ. J) GU TO 16 AN(IJ) = U. GO TO 17 16 AN(IJ) = 1. 17 CONTINUE 18 CONTINUE KETURN

المراجع والمراجع المحاري المراجع والمتحمي والمراجع والمراجع والمراجع والمراجع والمحموم والمحمولة المراجع والمراجع

END

DEC18910 DECI.8915 DECL8920 DECI.8925 DECI.8930 DECI.8935 DECI.8940 DECI.8945 DEC1,8950 DECL8955 DECL8960 DECL8965 DEC1,8970 DECI.8975 DECL8980 DECL8985 DECI,8990 DECI.8995 DECL9000 DECL9005 DEC1.9010 DECL9015 DEC1.9020 DECI.9025 0ECL9030 DECL9035 DEC1.9040 DEC1.9045 DECL9050 DECL9055 DECL9060 DEC1.9065 DECL9070 DECL9075 DECLOOSO DECL9085

DECI,8905

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Figure 5.1.1b

$$\begin{bmatrix} A & f \end{bmatrix} \begin{pmatrix} -y \\ 1 \end{pmatrix} = \begin{bmatrix} 0 \end{bmatrix}$$
 (5.1.16)

In this way, only the quantity $U_1^{-1} U_2$ needs to be evaluated for the augmented matrix and, provided that interchanges are not allowed on the last column, the product A^+f will appear in it. If the actual matrix A^+ is required, it can be obtained by augmenting to A the identity matrix instead of f. It must be stressed however, that the inverse A^+ should never be computed purely for the purpose of determining the solution of a matrix equation since this operation alone requires more computation than is needed to find the product A^+f .

Figure 5.1.1 contains a Fortran subroutine NULL which operates on the above principles. The input to the program is the M x NTOTAL array A of which N columns belong to the coefficient matrix and the rest are forcing functions. It returns the nullvectors of the system in the array AN, storing them as MAN component vectors, where MAN need not equal N.

5.2 GENERAL SOLUTIONS

The similarity of solutions produced by generalized inversion of the matrix equivalent of a differential equation and by the analytic theory of differential equations is established in the following two theorems. THEOREM 5.2.1 (DISCRETE FORM OF THE FIRST FUNDAMENTAL THEOREM OF CALCULUS). Let D^+ be the generalized inverse of a differentiation matrix as defined by (2.1.11) and suppose that y is given by $y = D^+z$. Then

$$Dy = z$$
 (5.2.1)

<u>**PROOF:**</u> In the definition of D, r = m so that

$$Dy = DD^{+}z = \begin{bmatrix} L_{1} \end{bmatrix} \begin{bmatrix} U_{1} & U_{2} \end{bmatrix} \begin{pmatrix} U_{1}^{-1} & L_{1} \\ U_{1} & L_{1} \\ 0 \end{bmatrix} \begin{bmatrix} z \end{bmatrix} = z$$

THEOREM 5.2.2 (DISCRETE FORM OF THE SECOND FUNDAMENTAL THEOREM OF CALCULUS). Let z be any vector such that z = Dy where D is a differentiation matrix as defined by (2.1.11). Then

 $D^+z = y + c$ (5.2.2)

where c is a vector with arbitrary but equal components.

<u>PROOF</u>: This result is a direct consequence of equation (5.1.5) and the fact that the nullspace of D contains equi-component vectors.

The relationship between the discrete and the continuous forms of these theorems will be illustrated by the following

simple example. Consider the problem of determining the general solution of the matrix equivalent of the equation

$$\frac{dy}{dx} = f \tag{5.2.3}$$

when y is approximated by a quadratic polynomial. According to (2.1.16), the rectangular matrix equation is

$$\frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$$
(5.2.4)

This coefficient matrix may be decomposed as

$$\frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} = \frac{1}{L} \begin{pmatrix} 1 & 0 \\ -\frac{1}{3} & 1 \end{pmatrix} \begin{pmatrix} -3 & 4 & -1 \\ 0 & -\frac{8}{3} & \frac{8}{3} \end{pmatrix}$$
(5.2.5)

so that

$$L_{1}^{-1} = \begin{pmatrix} 1 & 0 \\ \frac{1}{3} & 1 \end{pmatrix} \qquad U_{1}^{-1} = \begin{pmatrix} -\frac{1}{3} & -\frac{1}{2} \\ 0 & -\frac{3}{8} \end{pmatrix} \qquad U_{2} = \begin{pmatrix} -1 \\ \frac{8}{3} \end{pmatrix}$$

Consequently, the general solution of (5.2.4) is

$$\begin{pmatrix} y_{1} \\ y_{2} \\ y_{3} \end{pmatrix} = L \begin{pmatrix} -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{8} & -\frac{3}{8} \\ 0 & 0 \end{pmatrix} \begin{pmatrix} f_{1} \\ f_{2} \end{pmatrix} + \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix} c \quad (5.2.7)$$

where the constant c is arbitrary. By the second fundamental theorem of calculus, the general solution of (5.2.3) is

$$y = \int f \, dx + C' \qquad (5.2.8)$$

where c' is an arbitrary constant. It can be easily verified that multiplying (5.2.7) by the matrix (5.2.4) or multiplying (5.2.8) by $\frac{d}{dx}$ returns each of these equations to their original forms. Hence, Theorems 5.2.1 and 5.2.2 imply for discrete representations of differential equations the same properties that are obtained for continuous spaces from the two fundamental theorems of calculus.

Consider next the differential equation

$$\frac{dy}{dx} = y \tag{5.2.9}$$

According to (2.3.10), the discretized form of this equation in a quadratic space is

$$\frac{1}{L} \begin{pmatrix} -3 & 4 & -1 \\ 1 & -4 & 3 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = \frac{1}{3} \begin{pmatrix} 2 & 2 & -1 \\ -1 & 2 & 2 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix}$$

Using (5.1.5), the homogeneous solution of (5.2.10) is

$$\begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = \frac{1}{L^2 - 6L + 12} \qquad \begin{pmatrix} L^2 - 6L + 12 \\ -1/2 L^2 + 12 \\ L^2 + 6L + 12 \end{pmatrix} \qquad [c]$$

Equation (5.2.11) gives all of the possible quadratics which have their derivatives equal to their L_2 projections on the space of linear polynomials in an interval of length L.

From (5.1.5), the coefficients of the approximate general solution of the arbitrary differential equation (3.1.1) will be given by

$$y = A^{+}f + Nz$$
 (5.2.12)

Here, for a p'th order differential equation, z is in general a vector of p arbitrary constants. However, if an interpolation point happens to coincide with a singular point in the differential operation, the rank of the matrix A may be decreased by one and the number of arbitrary constants in (5.2.12) increased by one. In this case, it is best to add a boundary condition equation at the singular point to the differential equation in order to remove the undesired ambiguity.

A feature of the subroutines DIFFEQ and DIFF2D that has not yet been described is the fact that they call the subroutine NULL near the end of the programs. The effect of this call statement is to generate solutions of the form (5.2.12) for each of the matrix equivalents developed. Thus, DIFFEQ and DIFF2D not only evaluate matrix equivalents for arbitrary differential equations, but determine their general solutions as well.

5.3 PARTICULAR SOLUTIONS

One of the most important features of the analytic theory of differential equations is that particular solutions of differential equations may be obtained by adding boundary conditions to general solutions. This procedure will now be derived for the numerical solution of differential equations using the approximate general solution technique of the preceding section and equation (5.1.15).

In general, a boundary condition for a p'th order ordinary differential equation is an equation of the form

$$B_{\sum_{k=1}^{N}} B_{k} y(x_{k}) = g(x_{o})$$
(5.3.1)

where the $\{B_k\}$ are linear operators of order less than P, g is a given function and the $\{x_k\}$ are specified points in an interval I. If y is the projective solution of a differential equation, then it **is** given by

$$y = \sum_{i=0}^{n} y_i b_i$$
 (5.3.2)

Therefore, (5.3.1) becomes

$$\begin{array}{ccc} B & n \\ \Sigma & \Sigma & b_{k} \\ k=1 & i=0 \end{array} \begin{array}{c} b_{k} & y_{i} = g_{o} \\ i \end{array}$$
(5.3.3)

where

$$b_{k_{i}} = B_{k} b_{i} (x_{k})$$
 (5.3.4)

$$g_{o} = g(x_{o})$$
 (5.3.5)

The numbers b_{k_i} are easily evaluated by using the elementary differentiation and function matrices of Chapter 2.

If q independent boundary conditions are specified for a differential equation, then q independent equations of the kind (5.3.3) are produced. These may be written in matrix form as

$$By = g$$
 (5.3.6)

In conjunction with the general soution (5.2.12), this yields, equation (5.1.15)

$$y = (A^{+} - N(BN)^{+}BA^{+})f + N(BN)^{+}g + NMu$$
 (5.3.7)

where u has (p - q) arbitrary components. Thus, as long as it does not overspecify the problem, any number of boundary conditions may be added to the general solution (5.2.12) at one time. With each equation added, the vector of unknown coefficients becomes shorter until, when enough conditions are imposed, only a particular solution remains.

It is important to observe that any particular solution obtained from a general solution of a differential equation by

this procedure will satisfy both the differential equation and the boundary conditions, irrespective of what they might be. Hence, the locations of the points x_k where the boundary conditions (5.3.1) are applied is immaterial. If the point locations do not coincide with the endpoints of the interval, then the solution region will overlap the interval defined by the boundary conditions. However, this will not affect the validity or the accuracy of the particular solution within the confines of the specific boundary value problem.

This development is in sharp contrast to the solutions obtained by the finite element method or by ordinary projective methods in general where the solution region invariably coincides with the region defined by the boundary conditions. In fact, a large part of the difficulties and complications encountered in ordinary finite element analysis [22,23] is a direct result of defining unwieldy elements, such as the triangular, tetrahedral and iso-parametric elements, which are used to form complicated geometric shapes. With the necessity of matching boundary and solution regions removed, simple element shapes such as the rectangle and cube have an obvious advantage.

The procedure of specifying boundary conditions by restricting the behavior of the general solution of a differential equation is implemented for ordinary differential equations by the subroutine BOUND, shown in Figure 5.3.1. This subroutine accepts any number of arbitrary boundary conditions of the type (5.3.1) and generates particular solutions

					2501 5225
	SUCZO STATE REUNDLC, NORTS, NONHOM, NULLD, NULLV, Y, NWBC, NR, ND,	DECL4735		DD 20 ::EL = 1 ::ELMTS	54615010
	JUST CONST. XSC . MIMAXD. NELMTS. WT LAMDA, RITEP)	DECL4740		I = I + NO	05015015
		DECL4745		IF(XBC(NBCP)) LE. XD(I)/ GU 10 5	DECLEORO
	HE WE WARE THE BOWDARY CONDITIONS TO A GENERAL SOLUTION STORED	DECL4750	20	O CONTINCE	0=015025
	IN THE ARRAY ICL. ANY NUMBER OF BOUNDARY CONDITIONS MAY BE	DECL4755 C		NOW ADD THE ADDROLUTION FRANCENT FACTORS	DEC15030
	CREATED THE NUMBER OF SOLUTIONS FOUND IS DESIGNATED BY	DECL4760 C		NUW ADD THE COORDINATE-DEPENDENT TACTORS	05015035
	THE AD NOTAN CONDITION INSTRUCTIONS NEED TO BE STOPED	DECL4765 C		AND REPORT AND A REAL AND A REAL AND	070150-0
	THE APP VINCTALL STARTING IN LOCATION (NUN!)	DECL4770		CALL OPGEN2(INSINGLOISIGLOCS LEGPEN)	570190-5
	IN THE ARRENT INDIA OF DIMETRIC IN CONTINUES.	DECL4775		* UP3,LUP3,EXP7	620-50-50
	$\nabla T = \nabla T $	DECL4730		XO = XO(NPT)	520 5055
	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	DECL4785		MX = FEDK(F)	DECL51+2
	(1) = (1) + (1)	DEC14790		DD 21 I = 1 MX	DECLEDES
	Englisher Ten and a second sec	DECL4795	2	1 TEMP(1) = ALPHA(XBC(ABCP)) TAMAXOJ(ABCP)	DECLISOTO
		DECL4800 C		THE ADDRESS AND THE ADDRESS AN	5=015075
	LOJIC PRELIZION NOVER CONC	DECL4805 C		THE BOUNDARY CONDITION OPERATOR IS STORED IN THE REAR D	07015090
	$L_{0}L_{0}L_{0} = L_{0}$	DECL4810 C		•	05015055
		DECL4815		i J = - M X	55015070
	HOLE - HOLE. - FRIEL - NOIT	DECL4820		$U_1 Z_3 J = 1 N_0$	5=1.5-=5
	- FREE	DECL4825		DOUBLE = 0.CO	DECLENCO
		DECL4530		IJ = IJ + iX	DECLSICS
		DECL4835	_		D=C15110
	NECE - NOCHON	DECL4840	2	2 DOBLE = DOUBLE + DBLECTEAPCITADECT + 1377	DEC15115
		DECL4845	Z:	3 B(J) = SAGE(DOBEE)	DECLENZO
,		DECL4850		IJ = IEQ - NUL	DEC15125
-		DECL4855			2012012
	NETERMINE OF THERE ARE MORE BOUNDARY CONDITIONS	DECL4860		I + (LAF)A = 2 - KJ	061.5195
		DECL4865			05015140
-	ETILSTR(14:14) _NE. 4000) GD TB 4	DECL4870			DECL51-5
	IF(IFO_FO_C) RETURN	DECL4875		IJ = IJ + DL	DECL5150
	GC TR 13	DECL4880		$I_{0} I_{0} = I_{0} + NOLD$	DECL5153
4	$6001 = 6001 \pm 1$	DECL4885		NJ = NJ + 2075	CEC1516Q
•		DECL4890		$KJI = KJ = KI_{I}$	DECLEISE
	FIRST THE COORDINATE-INDEPENDENT FACTORS OF THE CURRENT	UELL4895 C		THE CONDITIONS ON THE ARBITRARY CONSTANTS IN THE GENERAL	DEC19170
	SECHDARY CONDITION ARE FOUND	DECL4900 C		SOLUTION ASE STOPED IN THE ARRAY 'VI	DEC15175
		020L4409 C		SECTION THE STORED IN THE RECENT	CEC15187
	CALL DPGEN1(INSTRANUNANDANDALDCALDALFCNALFADPSALROWALPOWERAUPAO)			DO(B) = -DB(E(V(T, L)))	DECL5185
	IF(LO .LE. 1 .ANC. LF .LE. 1) GD TO 992	DECL4415			DEC15190
	IF(LF .LE. 1) 60 TO 11	DEC14920			DEC15195
	$h_{\rm BM} = iABS(LFC)(LF - 1)) + 1$			DO(A + K - 1)	DECL5200
11	1F(LD .LE. 1) GD TO 12	DECL4930		16(1077.14804) GB TO 14	DECL5205
	hum = hAx0(hum, IABS(LOC(LO - 1)) + 3)			$D_{01}B_{2} = D_{01}B_{2} + D_{3} \in (B(K) \times C(K + KJ)))$	DECL5210
12	IEQ = IEQ + 1	02014999	1.	A CONTINUE	25015215
	IF(IEQ .GT. 1) GD TO 32	DEC14940	1.	Y(TJ) = SUCL(DOUPLE)	DEC15220
	$DC 31 I = 1_2 \text{MSIZE}$	DEC14955			DECL5235
31	f(1) = 0.			Y(1,0) = SSG(0,00,0)	DECLE230
32	IF(L3 - EQ. C) GD TO 7	DECL4900			DECL5235
	LEPS = + LREW(1)+RD	02024700		$7 16/16 = 60 = 0$ $-7R_{\odot}$ (AMDA) GO TO 9	DECL52-0
	DO 6 L = 1,LC/2	DECLATIO		a safes sheet na source non mark and source a	DECL52-5
	IF(LOC(L) .LT. 0) GD TO 990	DECLA912 C		THE MATRIX IS AUGMENTED BY THE POINT VALUES OF THE BOUNDARY	DĘCL5250
		DEC14985 C		CHUDITIAN FORCING FUNCTION	DEC15255
	THE ELEMENT NUMBER OF THE BOUNDARY CONDITION PUINT IS THELT	DEC14990 C		TANTI IA CENTRAL E	02012200
		DECL4995		b = 8 + 1 + 1 + 1 + 2	DEC152±5
	NBCPT = NBCPT + 1	DECL5000			DECL52TC
	I = 0				
					0
					\mathbf{N}

.

* •

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, , ...
Figure 5.3.1b

LDCL1 = IAPS(LFCN(L)) IF(INSTR(LDCL1) +LT, -1000) NPT = - INSTR(LDCL1) - 1000 IJ = IEQ + NUL#2 3 f(IJ) = Y(IJ) + (FLDAT(LOCL1)/FLDAT(LFCN(L))) * FCU(XD(NPT),INSTR,CONST,LOCL1,LFCN(L + 1)) 9 IF(IEQ +LT, NULLD) GD TO 3 13 IF(LAPDA) IEO = NPL NGEC = NDMENT + IEQ CALL NULL(Y,NUL,NFREE,NONHOM,L,IRDW,JCOL,B,NONHOM,NWBC,0) IF(L +DE, HUL +AND, LAMDA) GD TO 991 THE PRUDUCT (C * B' IS EVALUATED IJ = - ND LD = NDELTS*NVBC LF = 1 IF(LAPDA) LF = 2 %J = - NDMENT CD J 0 : = 1,NGAC %J = %J + HONHOM L = 0 LC 10 N = 1,NELMTS	DECL5275 DECL5280 DECL5285 DECL5295 DECL5300 DECL5305 DECL5315 DECL5320 DECL5320 DECL5320 DECL5330 DECL5335 DECL5340 DECL5345 DECL5350 DECL5355 DECL5350 DECL5375 DECL5375 DECL5375	ND = 0 RETURN 391 WRITE(6,202) L,IEQ ND = 0 RETURN 100 FORMAT(1H1/1H0,43X,'THE SOLUTION WITH BOUNDARY CONDITIONS') 201 FORMAT(1-ERROP IN BOUNDARY CONDITION INSTRUCTIONS') 202 FORMAT(1-THE CALCULATED RANK',I2,' DDES NOT EQUAL THE TRUE RANK', * I2) END	
IJ = IJ + ND DB 10 LOCL1 = 1×LF DD 10 1 = 1×ND c'= L + 1 IK = L + 'LOPTS UDUBLE = 0.00 SD 15 K = 1×NDOHEM IK = IK + NDOTS 15 DDUBLE + DBLE(C(IK)*B(K + KJ)) IF(LOCL1 .EC. 1) GJ TO 16 Y(I + IJ + LC) =-SNGL(DOUBLE) GC TD 10 16 Y(I + IJ) = SNGL(DOUBLE) IC CONTINUE IF(LAYDA .DRNOT. RITEP) RETURN IN NON-EIGENVALUE PROBLEMS, THE PARTICULAR SOLUTION IS PRINTED %RITE(5,100) CALL RITE(Y,NDPTS,NWBC,O/LAMDA) GD TO 1 930 GAITE(0,201)	DECL53390 DECL5395 DECL5395 DECL5405 DECL5405 DECL5415 DECL5420 DECL5425 DECL5425 DECL5425 DECL5435 DECL5440 DECL5445 DECL5455 DECL5455 DECL5455 DECL5475 DECL5475 DECL5475 DECL5475 DECL5475 DECL5475	FUNCTION ALPHA(X, I, ND, XO, WT) THE ALPHA FUNCTION RETURNS THE VALUE AT X DF THE I'TH NEWTON- COTES INTERPOLATION POLYNOMIAL (I = 1,ND) OF (ND - 1)'TH CODEP (ND .LE. 10) IN AN INTERVAL OF LENGTH 1./WT STARTING AT XO: DIMENSION FAC(10) DATA FAC / 112624120720504040320362880. Q = (ND - 1) \approx WT \approx ABS(X - XO) ALPHA = 1.0 / (FAC(1) \approx FAC(ND - I + 1); IF(MOD(ND - I,2) .NE. 0) ALPHA \approx - ALPHA DO 1 M \approx 1,ND IF(M .EQ. 1) GD TO 1 ALPHA = ALPHA \approx (Q - (M - 1)) 1 CONTINUE RETURN END	DDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDDD

r - - - -

^RI | E [0, 20]

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in the form (5.3.7) from the general solutions produced by the subroutine DIFFEQ. In defining the boundary operators B_{k} , BOUND uses the same instruction code as DIFFEQ and also relies on the subroutines OPGEN1 and OPGEN2 to generate their matrix equivalents.

The point locations x_k in equation (5.3.1) are specified in BOUND by adding an extra instruction at the end of the instructions for each simple operator. This instruction must have a value of -(200 + n) where n is the location in the array XD of the coordinate of the point of application of the simple operator. This coordinate need not be equal to that of an interpolation node.

The principles developed in this section for the specification of boundary conditions are also applicable to partial differential equations in N-dimensional space, except that the points x_k in equation (5.3.1) need to be replaced by (N - 1)dimensional surfaces. As discussed earlier, the logical solution region for these problems is the N-dimensional cube with the basis functions of equations (3.5.6) and (3.5.1). In this case, the number of linearly independent functions in the domain space of the equation is

$$N_{d} = \prod_{k=1}^{N} (n_{k} + 1)$$
(5.3.8)

Acting on this basis with a partial differential equation having only a p_i 'th order differential operator in the x_i direction, yields a range space of dimension

$$N_{r} = (n_{i} - p_{i} + 1) \prod_{\substack{k=1 \\ k \neq i}}^{N} (n_{k} + 1)$$
(5.3.9)

Consequently, the number of arbitrary constants in the solution is

$$N_{n} = p_{i} \prod_{\substack{k=1 \\ k \neq i}}^{N} (n_{k} + 1)$$
(5.3.10)

This is exactly p_i times the number of interpolation nodes in an (N - 1) - dimensional surface orthogonal to the x_i direction. Therefore, if p_i boundary conditions are specified on this set of interpolation nodes, the particular solution of the differential equation will be unique.

In general, the dimensionality of the null space of an approximate general solution of a partial differential equation will be

$$N_n = \prod_{k=1}^{N} (n_k + 1) - \prod_{k=1}^{N} (n_k - p_k + 1)$$
 (5.3.11)

where p_k is the order of the differential equation in the x_k direction. Therefore, a unique solution of a differential equation will result when N_n linearly independent equations are added to its general solution. This is most readily accomplished by enforcing the boundary conditions at a sufficient number of distinct points in each of the (N - 1)- dimensional subregions.

The subroutine BOUN22 shown in Figure 5.3.2 performs this boundary condition procedure for two-dimensional partial dif-

Figure 5.3.2

					05202441
	SUBBOUTINE BOUNDO(C.NDPTS,MONHOMSNULCD,NULLV,Y,NWBC,NRX,NRY,ND,	DE2D2190		KJ = KJ = NDPTS	05202465
	INSTR. CONST. NUM, XD, YD, NELMTS, WTX, WTY, LAMDA, RITEP, YI)	DE202195		KJ1 = KJ + NDPTS/2	DE202470
-		DE202200		DO 17 J = 14NUMBER	05202475
	BOUNDD ADDS THE BOUNDARY CONDITIONS TO A TWO-DIMENSIONAL	DE202202		IJ = IJ + HUL	DF2D2490
Ĭč	GENERAL SOLUTION OF A PARTIAL DIFFERENTIAL EQUATION STORED IN THE	DE202210			DE2D2485
	ARRAY ICI.	DE202219		NJ = NJ T (UPI) N N = VII - PROTE	DE2D249C
lè		DE202220		Double = Dote(v(1))	DE202495
1	DIMENSION C(1),Y(1),INSTR(1),CONST(1),XD(1),YD(1),WTX(1),WTY(1),	DE202230		000000 = 00000000000000000000000000000	DE20250C
	* YI(1) • • • • • • • • • •	05202235			DE2C2505
	DIMENSION IROW(100), JCDL(200), LDC(100), LPOWER(100), LROW(100),	DE202240	e		DE2D2510
	LFCN(100),B(4096),DP(4096),TEMP(4096),DPS(1)	DF202245	č	DECL 5195-5220	DE202515
Ł	NDSQ = ND ++ 2	DE202250	ċ		DEZC252C
iS	DECL 4810-4905	DE202255		IF("NOT, LAMOA) GD TO 17	DE202525
19	BELE 401044-05	DE202260	C		05202331
1	CALL TREATH AND THE NUME NO NO NO NO CALTER FOR SELEONS	DE202265	C	DECL 5230->290	05202934
ł.		DE202270	C		06202545
L	A Branchyon y - Constraint (Second Second	DE202275		* # FCNZD(XD(NPT))YD(NPT))INSINJCUNSIJCUCLIJCFCN(C = 1)	DE202550
Ľ	DECL 4915-4960	DE202280	ç	DECL #300-5338	DE202555
L		DE202285	5		02202565
1	LOPS = 0	DE202290	Ĺ,	11 - NDSO	DE20256:
1	00 17 L = 1/LD/2	05202292		to = NBC = NFEMTS = NDS0	DE50257C
i.	c	05202305	r		DE202575
	C DECL 4975->005	DE202310	č	DECL 5345-5375	DE2D2580
	C	DE202315	č		CE20258:
Ì.	I = I + NDSQ	DE202320	-	IJ = IJ + NDSQ	DE202590
	IF(XBC(NBCPT), GE, XD(I), AND, ABC(NBCPT), LE, YD(I + ND;	SQDE202325		DÖ 10 LOCL1 = 1/LF	DE202593
	AND, AND (NPCP) + CE, TETT AND, TOURDON TO CARD AND	DE202330		DD IO I = 1/7DSQ	DE20200C
		DE2D2335	C		05202000
	20 CONTINUE CALL DEG2D2(INSTP/CONST/LGC/LPOWER/LEOW/L/ND/XD/YD/WTX/WTY/	DE2D2340	ç	DECL 5400->535	05202010
		DE2D2345	C		DF20262C
	$x_0 = x_1(k_0)$	DE202350		END	
	ŶŎ = ŶĎ(KJ)	DE207352			
	WX = WTX(NEL)	05202365			
	WY = XTY(HEL)	05202370			
	MX = LRJW(L)	DF2D2375			
	MY = LRGW(L + 1)	DE202380			
1		DE202385			
		DE202390			
	VALUE & ALPHAIXU(HP1/JIJ)/AJAUANA/	DE202395			
		DE202400			
	A VILLE & ALCHALYD (NPT) JJHYJYOJWY)	DE2D2405			
		DE202410			
1	$D_{1}^{2} = 30 J = 1, NDSQ$	DE202412			
	DOUBLE = 0.00	05202429			
	ÍJ = ÍJ + K	05612427			
-1	00 6 1 = 1,K	DE202430			
	6 DOUBLE = DOUGLE + DBLE(YI(I) = OP(I + IJ))	DF2D2440			
	30 B(J) = SNGL(DJ:BLE)	DE202445			
	IJ = IEQ - KUL	DE202450		·	ليب
1		DE202455			0
1	IF(LAMUA) KJ # 4 # KJ				5

ferential equations. In most operations, BOUN2D is similar to BOUND. However, BOUN2D requires two coordinates for every boundary condition point location and these are stored in the arrays XD and YD.

Finally, note that if equation (5.3.6) is solved directly, it yields

$$y = B^{+} g + N_{B} w$$
 (5.3.12)

This equation defines a function space in which every function satisfies the boundary conditions. Hence, matrix generalized inversion can be used to advantage in conventional projective methods where such function spaces are often desired. In particular, many finite element formulations result in a matrix equation of the form [22]

Ay = f (5.3.13)where A is a symmetric matrix. Consequently, inserting (5.3.12)into (5.3.13) and pre-multiplying by N_B^T to retain symmetry of the coefficient matrix, gives

 $N_B^{T} A N_B w = N_B^{T} f - N_B^{T} A B^{+} g \qquad (5.3.14)$

When solved for w, (5.3.14) yields, in conjunction with (5.3.12), a solution of the boundary value problem which satisfies the boundary conditions (5.3.6).

The above process is generally performed implicitly in the finite element method [22] by numbering adjacent points with the same number and by eliminating rows and columns in the coefficient matrix in order to specify Dirichlet boundary conditions. Equation (5.3.14) explains why the procedures work and extends them to problems involving mixed boundary condition.

5.4 SOLUTION IN MULTIPLE REGIONS

Since the accuracy of numerical approximations such as (5.2.11) increases as the size of the interval used decreases, a basic method for obtaining highly accurate solutions of differential equations over a large region is to determine the general solution of the equation in several adjacent sub-regions or elements and then to tie these solutions together with continuity conditions. Provided that the continuity requirements are sufficiently stringent and of the right number, the composite solution obtained in this way will provide a good approximation of the general solution over the entire region.

In forming such composite solutions for a differential equation, one must be aware that they are, in general, not optimal in the sense of definition (3.1.3). This is because the projection matrices in a multi-regional solution are designed to minimize the residual function in each element separately and not in the overall region. If it is desired, the optimal approximate solution of a differential equation can be determined by evaluating the projection matrix (2.3.9) for every multi-regional combination of basis functions used. However. this procedure is not recommended as a practical solution technique for differential equations because the evaluation of complicated projection matrices requires considerable computation while the projection matrices for individual elements are pre-calculated. Thus, in most cases, given equal amounts of computer time, a much more accurate solution will result if in the place of evaluating fancier projection matrices, more ele7

ments are used.

In many respects, these ideas are similar to those encountered in ordinary finite element analysis. The main difference is that while in the finite element method continuity conditions are usually satisfied implicitly by using compatible elements, with the piecewise general solution procedure, the requirements are best imposed explicitly. The superiority of using overall conjugate approximation basis functions to local ones has, however, already been established in finite element analysis by Brauchli and Oden [29].

The most obvious method of obtaining a consistent set of continuity requirements for the overall solution of a p'th order differential equation is to set the solution and its (n - 1)'st derivatives equal across inter-element boundaries. This procedure has the effect of transforming the distinct sets of interpolatory basis functions in each of the elements into a single set of spline-like functions in the overall region. Therefore, for an ordinary p'th order differential equation that has had its solution region divided into m elements, the following p(m - 1) boundary conditions are imposed

$$\frac{d^{i}}{dx^{i}} \quad y(x_{n}^{s}) = \frac{d^{i}}{dx^{i}} \quad y(x_{0}^{s-1}) \qquad \begin{array}{c} 0 \le i \le p-1 \\ 1 \le s \le m-1 \end{array}$$
(5.4.1)

where the index s designates the element from which the approximate value of y should be taken. When assembled for the entire region, these conditions yield the matrix equation

Cy = 0 (5.4.2)

where C is a $p(m-1) \times (mn)$ matrix with components

$$C_{ij} = \begin{cases} D_{nj}^{(i)} & \text{if } j \text{ is in element } s \\ -D_{ij}^{(i)} & \text{if } j \text{ is in element } s+1 \\ 0 & \text{otherwise} \end{cases}$$
(5.4.3)

Together with (5.2.12), this gives, from (5.1.15)

$$y = A^{+}f - NC(CN)^{+}CA^{+}f + NMu$$
 (5.4.4)

where M is a $pm \ge p$ matrix of the nullvectors of C and u is an arbitrary p component vector. Equation (5.4.4) gives the general solution of the differential equation, including an approximate inhomogeneous solution and p approximate homogeneous solutions, each of which is C^{p-1} continuous in the entire region.

The continuity conditions (5.4.1) are closely associated with the theory of spline approximation [57]. In this theory, interpolating polynomials are defined in which some of the function specifying parameters are the values of the derivatives of the function at the endpoints of the interval. For example, in the case of cubic splines, the four parameters defining the polynomial are the function and derivative values at the endpoints of the interval. When applied to approximating functions, derivative continuity of cubic splines are established by setting adjacent derivative parameters equal, leaving two free parameters in each interval.

The functions corresponding to the solution (5.4.4) also

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satisfy the continuity conditions (5.4.1) at the edges of elements and are, therefore, similar to spline functions. In fact, when both sets of functions are *n*'th order polynomials, they span identical function spaces, although their representations are different. In the case of cubic equi-spaced Lagrangean interpolation polynomials, for example, the four function specifying parameters are the nodal point values $y_i = y(x_i)$, $x_i = x_o + (L/3)i$, *i*=0,1,2,3. When this polynomial is required to satisfy derivative continuity at both ends, the nodal values will be given by

y = Gu (5.4.5)

where G is a 4 x 2 matrix and u is an arbitrary two component vector. Here, as with the cubic splines, there are two independent parameters in each element, although in this case the components of u do not equal the values of the function at the endpoints of the interval.

The advantage of using the second type of polynomial instead of the first, is that the equi-spaced Lagrangean interpolation polynomials, as well as the procedures in equations (5.4.2) to (5.4.4), are both simple to define algebraically and easy to work with computationally - even in high order cases. By comparison, the calculation and use of general high order spline functions requires exceedingly complicated algebraic manipulation [57].

As an illustration of the procedure, suppose that the differential equation (5.2.9) is solved in n consecutive intervals of length L. In the *i*'th interval, the solution is given by (5.2.11) with $c = c_i$. According to (5.4.2), in this case CN=0 and this yields

where

$$a = L^2 - 6L + 12$$

 $b = L^2 + 6L + 12$
(5.4.7)

Consequently

$$c_{i} = \left(\frac{b}{a}\right)^{i-j} c_{1}$$
 (5.4.8)

and the entire solution contains only one arbitrary constant.

The solution (5.2.11) with (5.4.8) can be shown to converge to the exact solution of (5.2.9) by observing that

limit
$$c_n = 1 + nL + \frac{n^2L^2}{2!} + \frac{n^3L^3}{3!} + \dots = e^{nL}$$
 (5.4.9)
L $_{n \to \infty}$

Thus the value of the approximate solution at the third point in the n'th element is

Figure 5.4.la

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<pre>SUBROUTINE CONTIN(DN/NR,ND,MDN,NDN,C,NDPTS,NDNHOM,NULLD,NUCLV,</pre>	DECL3330 C DECL3335 C DECL3355 C DECL3355 C DECL3355 C DECL3355 D DECL3355 D DECL3355 D DECL3355 WEIGHT = $WT(N) \neq \#IPOW$ DECL3355 WEIGHT = $WT(N) \neq \#IPOW$ DECL3355 WEIGHT = $WT(N) \neq \#IPOW$ DECL3355 JON = JON + MON DECL3355 JON = JON + MON DECL3375 JC = JC + NC DECL3375 JC = JC + NC DECL3385 GO TO 25 DECL3395 25 IJ1 = I + JC DECL3395 25 IJ1 = I + 1 DECL3395 25 IJ1 = I + C DECL3395 25 IJ1 = I + 1 DECL3405 K2 = 0 DECL3405 K2 = 0 DECL3415 C DECL3425 C DECL3425 C DECL3425 C DEUBLE = DBLE(C(IJ1)) DECL3435 DBUBE = DBLE(C(IJ1))	00500500500500500000000000000000000000
<pre>%C = 10EN#NULLD*(NELMTS = 1) NC = NUNDUY + 2 = 1GEN NOWHOM = NC = MC N = MC#VC CD 3 K = 1/N 3 C(x) = 0. i = 0 if(NEL:TS .EQ. 1) GD TD 28 f4 = ND##2 THE DIFFERENTIATICY MATRIX IS STORED IN 'DP' CD 1 K = 1/N 1 CP(K) = 0. J = ND + 1 UD 2 K = 1/N,J 2 CP(K) = 1. IEC = 1 IPTA = 0 NCC = ID UCP = NUNDAN#MC</pre>	DECL3445 K1 = K1 + 10C DECL3445 K1 = K1 + 10C DECL3455 DNW = DN(K + JDN)+WEIGHT DECL3466 IF(N +EQ. 1) GD TD 20 DECL3465 DDUBLE = DDUBLE + DBLE(OP(K1)*DNW) DECL3475 DDUBLE = DDUBLE + DBLE(OP(K1)*DNW) DECL3475 DDUBLE = DDUBLE + DBLE(OP(K2)*DNW) DECL3486 12 CDNTNUE DECL3485 IF(N +EQ. 1) GD TD 5 DECL3485 IF(N +EQ. 1) GD TD 5 DECL3485 IF(N +EQ. 1) GD TD 5 DECL3495 5 IF(N +EQ. 4ELNTS) GD TD 6 DECL3495 C CIJJ = SNGL(DDUBLE) DECL3505 6 CDNTINUE DECL3505 6 CDNTINUE DECL3515 IJ = NCC DECL3525 IF(NCC +EQ. NR) GD TD 28 DECL3525 IF(NCC +EQ. NR) GD TD 28 DECL3535 C IN EACH PASS, THE ORDER OF THE DIFFERENTIATION MATRIX IS DECL3545 C DECL3545 C DECL3545 C DECL3545 C DECL3550 CALL OPRATR(OP, IED, NCC, IJ, ND, 1) DECL3555 IFON + 100 DECL3555 C DECL3555 C	
IN EIGENVALUE PROBLEMS, THE CONTINUITY CONDITIONS ARE IMPOSED TWICE 4 DD 17 L = 1/IGEN 1 = I - 1 20'. = - MO'. IF(L, EQ. 2) CON = - ND JC = - MC	DECL3555 IPUW = IPUW + 1 DECL3560 GD TO + DECL3565 19 CALL NULL(C;HC;NUNOWN;NC;IRANK;IROW;JCOL;DP;NUHOWN;NDNHCM;n) DECL3570 IF(IRAHK .ME.MC) GD TO 990 DECL3575 IPUW = NDPTS*NONHOM DECL3580 CD T I = 1;IPOM DECL3590 NCC = MDNM DN DECL3595 IF(LAMUA) GD TO 9	

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Figure 5.4.1b

C DECL3870 IJ = MON*NON DECL4140 USING THE EQUATION (C(LAST COLUMN) = DN(LAST COLUMN) THE DD 18 I = 1,IJ DECL3875 CEC14145 ÷C THE INHOMOGENEOUS SOLUTION IS EVALUATED DFCL3580 18 C(I) = DN(I)02014150 • 0 DECL3885 14 WRITE(6,100) 2501-153 IJ = NUPTS=APEREL DECL3890 IF(LANDA .OR. .NOT. RITEG) RETURN DECL-1:7 UD1 = HDN#HPEREL- NCC DECL3895 C 080141+5 CO 8 N = LINELMIS DECL3900 C THE GENERAL SOLUTION IS PRINTED DEC14170 JDN = JDN + NCC DECL3905 C DECL4175 00 8 1 = 1,MON DECL3910 CALL RITE(C;NDPTS;NONHOM;O;LAMDA) DECL4180 IJ = IJ + 1DECL3915 RETURN DEC14195 B C(IJ) = DN(I + JON)DECL3920 990 WRITE(6,200) IRANK, NM DEC:4170 IPOW = NOPTS+NR + MON DECL3925 0 = 011 50.4195 IJ = NUPTSANULLV - IPDA DECL3930 RETURN 54514210 JD1. = NONMULLY - NCC DECL3935 100 FORMAT(1H1/1H0,54X, THE GENERAL SOLUTION!) 0501-115 02 10 '1 = 1, NELMTS DECL3940 200 FORMATCI-THE CALCULATED RANKIJI4, DES NOT EQUAL THE TRUE RANKI, DECLARIC 13 = 13 + 12ga DECL3945 × 14) SEC14215 JON = JON + NCC DECI 3950 END DECL4220 IJ1 = IJ - NOPTS DECL3955 IJON = JDN DECL3960 DO 10 J = 1,HR DECL3965 IJ1 = IJ1 + HOPTS DECL3970 DC 10 I = 1,MDY DECL3975 IJON = IJDN + 1 DECL3980 10 C(1 + 101) = DN(100N)DECL3955 9 10 = 0 DECL3990 c DECL3995 С USING THE EQUATION (C = DN + DP, (DP = C+)), THE HOMOGENEOUS DECL4000 SOLUTION IS EVALUATED DECL4005 č DECL4010 JD'' = -NCCDECL4015 IPON = - NPEREL DECL4020 13 11 % = 1, RELNTS DECL4025 30% = 00% - MCC DECL4030 IPIN = IPDN + NPEREL DECL4035 02 11 1 = 1,40% DECL4040 100N ± 1 + 00N DEC1.4045 1C = 1C + 1DECL4050 IED = IPOW - NUNDWN DECL4055 IJ = IC - MOPTS DECL4060 00 11 J = 19101.00 DECL4065 IJ = IJ + .DPT5 DECL4070 1EC = 1EC + NONBWH DECL4075 IJI = IJDU + MDN DECL4080 BOUBLE = DALE(C(IJ)) DECL4085 00 13 K = 1/MPEREL DECL4090 IJ1 = IJ1 + HD1 DECL4095 . 13 1208LE = D"UBLE + DBLE(DN(IJ1)=0P(K + IED)) DECL4100 11 C(IJ) = SHOL(DOURLE) DECL4105 GE TO 14 DECL4110 28 IF(I .GT. 0) 60 TO 19 DECL4115 DECL4120 IF THERE IS ONLY ONE ELEMENT, THE SOLUTION IS SIMPLY TRANSFERED DECL4125 FROM DU TO C DECL4130 DECL4135 ___

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limit
$$y_{3n} = c_1 e^x$$
 (5.4.10)
 $h \to \infty$
 $L \to 0$

where x = nL.

The subroutine DIFFEQ in Figure 4.2.4 has been designed with multi-element solutions in mind. The program calculates the matrix equivalent of a differential equation in NELMTS consecutive intervals and evaluates the general solution of the equation in each of these separately. The continuous general solution is produced from these sectional solutions by the subroutine CONTIN shown in Figure 5.4.1. By calling DIFFEQ and CONTIN in succession, highly accurate, continuous multielement general solutions of ordinary differential equations can be obtained.

The procedure of imposing continuity requirements on multielement solutions of partial differential equations is very similar to that used with ordinary differential equations except that it must be done repeatedly and in different directions. For two-dimensional problems, the region is divided into matching rectangles in each of which the interpolation polynomials are of the same degree. The derivative continuity conditions (5.4.1) are imposed along each line of interpolation nodes in both the x and the y directions. These must be sufficient to satisfy the continuity requirements of the partial differential equation in a direction normal to the edge of a rectangle. Derivatives tangent to an edge are, of course, equal whenever the function is made equal on both sides of the edge. The degree to which cross derivatives are continuous is given in the following theorem.

<u>THEOREM 5.4.1.</u> Let two functions z(x,y) and w(x,y) be given by

$$z(x,y) = \sum_{i=0}^{m} \sum_{j=0}^{n} Z_{ij} l_i^{(m)}(x) l_j^{(n)}(y) \qquad \substack{a \le x \le b\\ d \le y \le e}$$

$$w(x,y) = \sum_{i=0}^{m} \sum_{j=0}^{n} W_{ij} l_i^{(m)}(x) l_j^{(n)}(y) \qquad b \le x \le c \\ i \le y \le e$$

where $l_i^{(m)}$ is defined in (3.5.2). Then, whenever

$$\frac{\partial^{i} z}{\partial x^{i}} \begin{vmatrix} & = & \frac{\partial^{i} w}{\partial x^{i}} \end{vmatrix} \qquad \begin{array}{c} i = 0, \dots, p \qquad (5.4.12) \\ x = b \qquad \end{array}$$

the cross derivatives $\frac{\partial^{p+q}}{\partial x^p \partial y^q}$ and $\frac{\partial^{p+q}}{\partial y^q \partial x^p}$ are continuous across the line x=b

$$\frac{\partial^{p+q} z}{\partial x^{p} \partial y^{q}} \begin{vmatrix} = \frac{\partial^{p+q} w}{\partial x^{p} \partial y^{q}} \end{vmatrix} = \frac{\partial^{p+q} w}{\partial x^{p} \partial y^{q}} \begin{vmatrix} q=0,\ldots,n & (5.4.13a) \\ x=b \end{vmatrix}$$

$$\frac{\partial^{p+q} g}{\partial y^{q} \partial x^{p}} \begin{vmatrix} z = \frac{\partial^{p+q} g}{\partial y^{q} \partial x^{p}} \end{vmatrix} = \frac{\partial^{p+q} g}{\partial y^{q} \partial x^{p}} \begin{vmatrix} z = 0 \\ z = b \end{vmatrix}$$
(5.4.13b)

PROOF: The elements of the bottom row of the matrix form of $u = \frac{\partial p}{\partial x^{p}} \left(\frac{\partial q_{z}}{\partial y^{q}} \right) \text{ are}$ $U_{m-p,j} = \sum_{\substack{k=0 \\ k=0}}^{m} \sum_{\substack{l=0 \\ l=0}}^{n} D_{m-p,k}^{(p)} Z_{kl} D_{lj}^{(q)} (5.4.14)$ where $D^{(p)}$ is the p'th order differentiation matrix. Similarly, the elements of the top row of the matrix form of $v = \frac{\partial^p}{\partial x^p} \left(\frac{\partial^q w}{\partial y^q} \right)$ are

$$V_{1j} = \sum_{k=0}^{m} \sum_{l=0}^{n} D_{1k}^{(p)} W_{kl} D_{lj}^{T(q)}$$
(5.4.15)

Hence

$$r_{j} = U_{m-p,j} - V_{1j}$$

$$= \sum_{k=0}^{m} \sum_{l=0}^{n} \left\{ D_{m-p,k}^{(p)} Z_{kl} - D_{1k}^{(p)} W_{kl} \right\} D_{lj}^{T(q)}$$
(5.4.16)

Now (5.4.12) implies that

$$\sum_{k=0}^{m} D_{m-p,k}^{(p)} Z_{kl} = \sum_{k=0}^{m} D_{1k}^{(p)} W_{kl} \qquad (5.4.17)$$

so that $r_{j} = 0$ and (5.4.13a) is true. The proof of (5.4.13b) is similar.

As was the case with DIFFEQ, the two-dimensional version of the program DIFF2D also computes the general solution of a differential equation in NELMTS separate elements. These must be placed in a matching pattern as described earlier. These elemental general solutions are then made continuous by the subroutine CONT2D shown in Figure 5.4.2. CONT2D imposes C^{p-1} continuity on the solution of a partial differential equation in the x direction and C^{q-1} continuity in the y direction, where p and q are the highest orders of the x and y derivatives

Figure 5.4.2

1							
	SUBROUTINE CONT2D(DN;NRX;NRY;ND;MD);NDN;C;NDPTS;NONHOH;NULLV;	DE2D1755			$JCADD = MCNDN \pm N1$		DE202025
1	* NCCONDATELCXATELCYANELMTS, WTX, WTY, LAHDA, RITEG)	DE2D1760			NUNA = ISTART + MNDN + N1		DE202030
1 e		DE2D1765			JLINE = NDOA = LINADD		DE202035
17	CONTED TAKES THE GENERAL SOLUTION OF A THOUSTHENSIONAL PARTIAL	06201770			DO 6 NITNE # 1.NO		DE202040
17	DIEEEDENTAL CONTRACTOR IN CONTRACTOR AND ESTADISTIC	ESDE201775			JETNE - JETNE - ITNADD		DE202045
12	CALINITY OF THE CURTION AND THE DEPARTIC REGIONS AND ESTADLISH	06201779			10H = 11THe = MSH		06202050
12	CONTROLIT DE THE FUNCTION AND ITS DERIVATIVES ACROSS THE	05201700			400 - 10400 - NON		0520203055
15	CLEMENT BUUNDAPIES.	02201762			JC = JCADD = HUN		05202030
16		DE201790			* # 1 + 1		05202000
1_	DIMENSION DN(I))C(I))WIX(I))WIY(I))IELCX(I))IELCY(I)	DE201795	<u>د</u>				06202003
10		DEZDIAOU	ç		0=0	GL 3030=3705	06202070
10	DECL 3385-3410	DE2D1809	Ç				DE202075
10		DE201810			KJ = JDN = KADD		DESDSOBO
	LUGICAL LAMDAJFIPSTJLASTJRITEG	DE201815	С				DE202085
ì.	NPEREL = ND ++ 2 - NRX + NRY	DE2D1820	C		DEC	CL 3710-3720	DE2D2090
10		DE201825	C				DE202095
l ē	DECL 3420-3445	DE201830			KJ ≖ KJ → KADD		DE202100
ĪČ		DE201835			DNW = DN(KJ) + WEIGHT		DE202105
1	MC = IGEN # NCCOND	DE201840			IF(FIRST) GO TO 20		DF202110
1		DE2D1845			DOUBLE = DOUBLE + DBLE(DP)	(K13*0NW)	DE202115
		DE201850		20	IF(LAST) GO TO 12		05202120
		Dc2D1855			DOUB2 - DOUB2 - DBIELOPIKS	2 1 + DNW 1	DE202125
	P(HDN = NC + HDN)	DE201940		1.	CONTINUE	c)+pinny	05202120
15		05201865		42	TEVETOETA CO TO E		06202135
12	UECL 3400#3485	02201000			17(71K317 GL 10 3		05702120
16		02291870			C(14) = 200C(DUDBCE)		00202140
	NK = NKX	UE201875		2	IF(GASI) GU IU O		UE232142
	LINADD = 1	DE201880			C(1J1) = SNGL(DUUB2)		05202150
	KADD = ND	DE2D1885		- 5	CUNTINUE		DEZDZ155
	GO TO 31	DE201890			FIRST = FALSE.		DE202160
	30 NR = NRY	DE2D1895			GU TO 33		DE202165
1	LINADO = ND	DE2D1900	Ç				DE202170
	KADD = 1	DESD1902	C		DEC	CL 3780-4215	DE2D2175
	31 ICONT = O	DE2D1910	C .				DE2D2180
0		DE201915			END		DE202185
10	DECL 3490-3580	DE2D1920					
10		DE201925					
1	ISTART = 1	DE201930					
	IF(L = EQ = 2) ISTART = 1 + NDN / 2	DE201935					
	32 FIRST = TRUE.	DE201940					
	LAST . FALSE.	DE2D1945					
	33 ICONT # ICONT + 1	DE2D1950					
	LE(KADD - E0- 1) 60 TO 34	DE201955					
	$\mathbf{x}_{i} = \mathbf{x}_{i} (\mathbf{x}_{i} (\mathbf{x}_{i}))$	DE201940					
		05201965					
1		DE201970					
1	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	DE201975					
1	91 - IECONIICUM - 17 97791 - 07910 - 07910	06201000					
	chara se BETANI = MIY(N) = MILIM	05201400					
1		05201000					
	34 N H IELLY(ILINT)	05501990					
1	IF(N +LT + C) 60 TU 17	DE202422					
1	IF(N , EQ, 0) 50 TO 32	0E202000					
1	AL = IELCY(ICDHT + 1)	DEZDZOOS					
1	WEIGHT = WTY(N) =# IPOW	DE2D2010					
1	35 IF(N1 .EQ. 0) LAST = .TRUE.	DE202015					
1	51 = 5 = 1	DE505050					
i i							œ
1							
1							

in the equation, respectively.

5.5 EIGENVALUE PROBLEMS

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The procedures developed thus far are not capable of solving eigenvalue problems. These are problems of the form

$$D_1 \ y = k \ D_2 \ y \tag{5.5.1}$$

that are accompanied by homogeneous boundary conditions. The difficulty stems from the fact that the constant k in this equation is not determined until all of the boundary conditions are specified. Therefore, it is not possible to treat differential equation eigenvalue problems independently from the boundary conditions with a numerical method and only particular solutions can be obtained directly. Of course, if a complete set of independent particular solutions of a differential equation are determined, together they constitute its general solution.

Despite the impossibility of separating the differential equation (5.5.1) from its boundary conditions, it is still extremely desirable to evaluate each contribution in the discretization process independently. A simple but effective device for developing such a procedure is to write the matrix equivalent of (5.5.1) in each element as

$$Az = 0$$
 (5.5.2)

where A is the m x2n matrix

$$A = [D_1 | D_2]$$
 (5.5.3)

$$\mathbf{z} = \begin{pmatrix} \mathbf{y} \\ -\mathbf{z} \\ -\mathbf{k}\mathbf{y} \end{pmatrix}$$
(5.5.4)

Equation (5.5.2) can be easily solved for z using (5.1.5)

$$z = Nu$$
 (5.5.5)

where u has (2n - m) arbitrary elements. The (n - m) continuity and boundary conditions can then be imposed on both the top and the bottom halves of z. Since these conditions are all homogeneous, they will reduce the number of arbitrary constants in (5.5.5) to m but will not alter its form:

 $z = M_W$ (5.5.6)

Consequently, setting the bottom half of z equal to -k times the top half results in the rectangular matrix eigenvalue problem

$$M_{b}W = -k M_{t} W \qquad (5.5.7)$$

where M_b and M_t are $n \ge m$ matrices of rank m. Thompson and Weil have proved [58] that the rank reducing numbers of this equation are identical to the eigenvalues of the standard, square matrix eigenvalue equation

$$Lw = -kw$$
 (5.5.8)

where L is the $m \times m$ matrix

$$L = M_t^+ M_b \tag{5.5.9}$$

Equation (5.5.8) may be solved by using standard techniques to determine the eigenvalues -k and the eigenvectors w. Then the vector y is found by using (5.5.6) and (5.5.4).

A simple problem will illustrate these properties. Consider the boundary value problem

$$\frac{d^{2}y}{dx^{2}} = -\lambda^{2}y$$
(5.5.10)
$$y(0) = 0 \qquad \frac{dy}{dx}(1) = 0$$

For the sake of simplicity, let y be approximated by a single quadratic element. Equation (5.5.10) becomes

$$\begin{bmatrix} 4 & -8 & 4 \end{bmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} = -\lambda^2 \begin{pmatrix} \frac{1}{6} & \frac{4}{6} & \frac{1}{6} \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} (5.5.11)$$

The scheme is to solve the equation

[1 -2 1 1 4 1] z = 0 (5.5.12)

with

$$\begin{pmatrix} z_1 \\ z_2 \\ z_3 \end{pmatrix} = \frac{24}{\lambda^2} \begin{pmatrix} z_4 \\ z_5 \\ z_6 \end{pmatrix}$$
 (5.5.13)

The solution of (5.5.12) is

where u is a vector of five unknowns. The boundary conditions imply that

Thus

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and (5.5.14) becomes

(5.5.17)

Imposing (5.5.13) results in

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$$\begin{pmatrix} 0 \\ 6 \\ 8 \end{pmatrix} = \frac{24}{\lambda^2} \begin{pmatrix} 0 \\ 3/4 \\ 1 \end{pmatrix} w$$
 (5.5.18)

This is equivalent to the standard matrix eigenvalue equation

$$\begin{bmatrix} 0 & 0 & 1 \end{bmatrix} \begin{pmatrix} 0 \\ 6 \\ 8 \end{bmatrix} W = \begin{bmatrix} 8 \end{bmatrix} W = \frac{24}{\lambda^{4/2}} W$$
 (5.5.19)

Therefore the approximate eigenvalue is $\lambda = \sqrt{3} = 1.73$ and the approximate eigenvector is

$$y = \begin{pmatrix} 0 \\ .75 \\ 1. \end{pmatrix} c$$
 (5.5.20)

This is to be compared to the exact solution $\lambda = \frac{\pi}{2} = 1.57$ and

$$y = c \sin \frac{\pi}{2} x$$
 $(y(\frac{1}{2}) = .707).$

It may be added that the above procedure is not stationary in the eigenvalues so that their accuracy will be of about the same precision as that of the eigenvectors. If more accurate eigenvalues of (5.5.1) are desired, they can be obtained by inserting the eigenvectors calculated by the method of this section into the Rayleigh quotient [43]

$$k = \frac{\langle y \mid D_1 y \rangle}{\langle y \mid D_2 y \rangle}$$
(5.5.21)

In the above example, for instance, using (5.5.20) in (5.5.21) results in λ = 1.65 - a 50% decrease in the error of the eigenvalue. It is interesting to note, however, that the new eigenvalue-eigenvector combination results in a much larger residual function than the original pair.

All of the computer programs for the discretization of differential equations given in the preceding sections of this thesis have been designed with the capability of handling differential equation eigenvalue problems in addition to the standard ones. Therefore, computer programs of all of the operations described for the solution of differential equation eigenvalue problems up to equation (5.5.7) are already available. The remaining steps are performed for both ordinary and partial differential equations by EIGEN shown in Figure 5.5.1. In

			~		02010214
	THE NEW WE NELVES NURSENWRG . C. YR Y Y A ROOTRARD TI	DECL5542	5	AND DE THE ETCONVANCES ARE DOINTED	025215820
	SUBROUTINE EIGENTY, ND MELMISTADE LATAR DU CONTRA LE	DECL5550	C .	ALL OF THE EIGENVALUES AND FRANCES	DECLEEZE
x	E NEIGEN)	DECL5555	C		02015830
	THE THE AND EXCENNECTORS	DECL5560		WRITE(6,101)	P101 5435
	THIS SUBRE TIME EVALUATES THE EIGENVALUES AND FIGENVE THE NUMBER	DEC1 5565		CALL RITE(RODTR/LD/1/0/1)	557:5610
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		DECL5530		IF(LU .G). KEISEN/ LO - GEISEN	DEC15435
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	AT (EXS (C), TROW(100), JCOL(200), LOC(100), B(100), DP(1)	DECI 5590		IJ = 0	
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	ADELIS = MELMES # MP	DECL560>	L.	- THET ARE ALSO INVITE DOMESTIC DE TRA	DECLEREC
9	1J = N.30400ELTS	DECL5610	C		DFC15335
, .	THE REPORT OF A STREED FOR LATER	DFCL5615		DD 29 J = 1 J L U	00115320
	THE LEFT HAND SIDE OF THE MATRIX NEEDS TO BE STORED FOR CATER	0501 5620		B(j) = 0.	
		000015405		$KJ = N_{\rm e}BC * (LDC(J) - 1)$	05010572
	-22-0E	DECLOOZO		$p_{1} = 25$ $I = 1 + 2p_{1} = 15$	DECISATO
		DECLOGIO			05015905
	DD 19 I = 1 J I J	DECL5635		13 - 13 + 1	02615910
. 9	$Y_{\mathcal{E}}(\mathbf{I}) = Y(\mathbf{I})$	DECL5640		DOORFE = 0°CG	DEC1 5915
•••		DECI 5645		DDUB2 = 0.DC	0000000000
		05015650		$\mathbf{I}\mathbf{K} = \mathbf{I} + \mathbf{M} \mathbf{H} \mathbf{F} \mathbf{L}^{T} \mathbf{S}$	
	THE WATTER AS FIRST REDUCED TO STANDARD FORM				25019425
	INT MARKIN STOLE IS THE REFERENCE TO THE	DECUSOSS			DECT2030
		DECL5660		IK = IK + AGECIS	D=CL5935
	CALL NULL(YANDELTS, NABLANELALATANA JUULACANADOS ANDOS OF	DECL5665		IJI = K + KJ	-27: 4010
	IFO = 1.WBC	DEC1.5670		DDUBLE = DDUBLE + DBLE(YK(IK)*C(IJI))	
	15/1 / 5- UXBC) 60 TO 991	DEC1 5675		24 DDUB2 = DDUB2 + DBLE(YR(IK) * Y(IJ1))	
			•	OP(1.1) = SrGL(ODU3LE)	
		DECESSOO		$B(t) = A(t) + ABS(5P(T_i))$	CEC12425
	THE SINGLE STANDARD SQUARE	DECL5665			02015950
	THE EISENVALUES AND EIGENVECTORS OF THE CONCLUSE	DECL5690		25 (1(1)) = 5/0L(5002)	04015955
	PATRIX ARE EVALUATED	DECL5695		B(J) = NDELTS/B(J)	0511 5012
		DFCL5700	1	29 CONTINUE	555:5575
	CALL RIMAT(C, VERDETISYISNWBC/NWBC/O/LD)	DECI 5705		IJ = 0	
			c		02013-59
:		DECESTIO	ř	THE FLEENVECTORS ARE ALSO SCALED	350L5962
20	$L_{L_{1}}^{(1)} = 1$	DECL5/15	č	THE EIGENEEDISKS AND PORTE	DECLEARD
	IF(LO .EQ. 1) GU 'C 2/	DECL5720	C		07615995
	K = LC + 1	DECL5725		$DD 28 J = 1 \mu LU$	35014013
		DECI 5730		DO 28 I = 1,NDELTS	2200000
	THE ETGENVALUES ARE DROERED ACCORDING TO THEIR MAGNITODES	00015735		r + LI = Li	
				$Y(T,I) = CP(T,I) \neq B(I)$	DECLECIC
	287766201 - 2162 -	DECL5740			DECL6015
		DECL5745		25 11(14) # 11(14) # 3(9)	DECL6020
	BC 21 1 = 1 j K	DECL5750		IJ = - NUELIS	0=01 4025
	VALUE = ABS(FICTI(LOC(I)))	DECL5755		UD 26 I = 1,LD	DEC
		DECI 5760		IJ = IJ + HDELTS	
	AV ' A ' # 55 21 - 1 # [14]2	02012745	c	· · ·	CECT4035
	$C_{1} = C_{1} = C_{1$	DECLOTOD	ž	THE TRANSEGRAED ELCENVECTORS ARE PRINTED	020160-0
	IF(AdS(ROB)I(LEC(0)))	DEC1.5770	L C	THE TRANSPORTED EIGENVEUTORS THE TRANSPORTE	- DECLED-5
	AVTOR = ABREATED (COCOTA)	DEC1.5775	6		22014520
	£ = 130(I)	DECL5730		WRITE(6,102) I	DECLARES
	incli = LPC(J)	DEC1 5755		CALL RITE(V,NDELTS,1,1,1)	25
		00000000		26 CALL RITE(VIJNHELTSJ1J1)	DECT9366
		DECLOTAD			
23	C CETHINDE	DECL5795			
27	7 50 22 I.= 1/L ⁰	DECL5800			_
2:	> REGTR(I) = REGTI(LOC(I))	DECL5805			
~ (0723 I = 1/L0	DFCL5810			10
	$\mathbf{F}_{\mathbf{T}}$			•	01
- 23	5 NUJ1111				

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. .

RWBC = LODECL6065 RETURN 991 WRITE(6,202) L, NWBC DEC16070 110 = Ò DECL6075 DECI.6080 RETURN DEC1.6085 101 FORMAT(1H0,58X, THE EIGENVALUEST) 102 FURMAT(1H-, 56X, 'EIGENVECTOR NUMBER', 12) DEC1.6090 202 FORMATCI-THE CALCULATED RANK 1, 12, 1 DOES NOT EQUAL THE TRUE RANK 1, DECL6100 12) DEC1.6105 END DECL6110

Figure 5.5.1b

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order to solve the standard matrix eigenvalue problem, these subroutines call a single precision version of the program RILMAT, published in the Union Carbide programming handbook [59]. Since the eigenvalues and eigenvectors of a general, real matrix may be complex, all of these programs have two arrays for each of these quantities, one for the real and one for the imaginary part.

5.6 NONLINEAR EQUATIONS

The techniques developed thus far provide a powerful method for solving linear differential equations. In this section, the application of this method to problems involving nonlinear differential operators will be considered. The procedure used is the well-known Newton's method for solving nonlinear operator equations and the discussion will draw heavily from reference [60]. In this procedure, the solution of a nonlinear equation is obtained by solving a sequence of linear equations, the solutions of which converge to the solution of the original equation.

The first concept which needs to be introduced when dealing with nonlinear differential equations is the Fréchet derivative of an operator. This may be given as follows [60]: <u>Definition 5.6.1.</u> Let A(y) be a nonlinear operator. The Fréchet derivative of A(y) at the point y_0 will be a linear operator $A'(y_0)$ such that

limit $\frac{|| A'(y_0) \Delta y - A(y_0 + \Delta y) + A(y_0) ||}{|| \Delta y || + 0} = 0$ (5.6.1)

provided that such an operator exists.

As a consequence of this definition, the quantity

$$r(y,z) = A(y) - A(z) - A'(z) (y - z)$$
(5.6.2)
will have the property that

the second se

· ;

limit
$$\frac{|| r(y,z) ||}{|| y-z ||} = 0$$
 (5.6.3)
 $||y-z|| \to 0$ || $y-z ||$

Thus, if the two functions y and z are sufficiently close, the norm of the expression (5.6.2) approaches zero at least as quickly as the square of the norm of the difference of the two functions y and z.

Consider now the problem of solving the nonlinear differential equation

A(y) = 0 (5.6.4)

and suppose that some function $z \in \Sigma$ is given. Then, from (5.6.2),

$$A(y) = A(z) + A'(y - z) + r(y, z) = 0$$
 (5.6.5)

In Newton's method, it is assumed that a function z can be found sufficiently close to the solution y of (5.6.4) so that, by virtue of (5.6.3), the quantity r(y,z) in (5.6.5) can be neglected [60]. If this is done, then (5.6.5) yields

$$A'(z) \ y \simeq A'(z) \ z - A(z)$$
 (5.6.6)

This equation is linear in the unknown y and can be solved by the methods in the previous section. However, since (5.6.6) is an approximation, the solution obtained in this way will not satisfy (5.6.4) exactly. Presumably, though, it will be closer to the correct solution y than was the function z in (5.6.6) and can be used as z in (5.6.6) to obtain an equation having an even better approximate solution. By repeating the above steps enough times, the solution of the linear equation (5.6.6) can be made arbitrarily close to the solution of the nonlinear equation (5.6.4). The conditions on the function z and on the operator A(y) for this procedure to converge are discussed in reference [60].

The effectiveness of Newton's method in treating nonlinear differential equations is limited by two principal defects. The first of these occurs whenever the Fréchet derivative of the operator A(y) does not exist in a neighborhood of the solution y. In these cases, Newton's method is invalid and cannot be applied to determine the solution. The other defect in Newton's method is the difficulty of determining an initial function z within the radius of convergence of the solution for some problems. Unless an a priori estimate of such a function is known, the problem of finding a useful initial function may require considerable effort and ingenuity.

In evaluating the Fréchet derivative A'(y) in (5.6.6) for arbitrary nonlinear differential equations, it is profitable to use the following two properties [60]

.

$$(A(y) + B(y))'|_{y=z} = A'(z) + B'(z)$$
 (5.6.7a)

$$(A(B(y)))'|_{y=z} = A'(B(z)) B'(z)$$
 (5.6.7b)

These properties permit one to form the Fréchet derivative of a differential operator by using the ordinary rules of differentiation, treating the unknown y as the independent variable

$$A'(z) = \frac{dA(y)}{dy}\Big|_{y=z}$$
 (5.6.8)

In this way, the Fréchet derivative of an arbitrary differential operator can be easily determined.

As a simple illustration of the procedure, consider the initial value problem

$$A(y) = \frac{d^{3}y}{dx^{3}} + y \frac{d^{2}y}{dx^{2}} = 0$$

$$y(0) = \frac{dy}{dx}(0) = 0$$
 (5.6.9)

$$\frac{d^{2}y}{dx^{2}} = 1$$

which governs the flow of a laminar boundary layer on a flat plate parallel to the stream [49]. The Fréchet derative of this operator is

$$A'(z) = \frac{d^3}{dx^3} + \frac{d^2z}{dx^2} + z \frac{d^2}{dx^2}$$
(5.6.10)

7

Consequently, equation (5.6.6) becomes

$$\frac{d^{3}y}{dx^{3}} + \frac{d^{2}z}{dx^{2}}y + z \frac{d^{2}y}{dx^{2}} = z \frac{d^{2}z}{dx^{2}}$$
(5.6.11)

Given an initial value for z, say z = 0, this equation may be solved by using the subroutines DIFFEQ, CONTIN and BOUND to obtain an approximate solution of (5.6.9). This approximate solution may then be used for z in (5.6.11) and the process repeated. It will be shown in the next chapter that after several interations, a highly accurate solution of (5.6.9) is obtained.

There are two minor program modifications of the linear subroutine package which increase their efficiency with nonlinear problems. One of these is to store the coordinate independent factors obtained from OPGEN1 in the first iteration. Since these factors contain no functions, they are unchanged by the iteration procedure and may be used repeatedly. Consequently, the step, CALL OPGEN1, may be skipped after the first iteration. The second modification is to place the z function values in a labëled common block accessible to subprogram FCN and the auxiliary functions FN1 -FN5. Then, whatever the value of z in the current iterate, the values of the functions dependent on z can be easily generated.

CHAPTER VI

AUTOMATIC SOLUTION OF DIFFERENTIAL EQUATIONS

The computer programs given in preceding chapters provide a computationally efficient means of solving any ordinary or two-dimensional partial differential equation. However, in their present form, these programs are impractical for everyday use by a wide spectrum of engineers because of their special input and output characteristics. This chapter presents the results of a major effort to write a differential equation solving program which is compatible with ordinary engineering training and experience. In order to do so, it has been necessary to develop a special-purpose computer language which is capable of deciphering mathematical statements and generating computer understandable instructions according to the code given in Chapter 4. By incorporating this specialpurpose language and the programs given previously, the solution of any linear ordinary or two-dimensional partial differential equation requires no analysis beyond the ability to write the differential equation on a standard keypunch or teletype terminal. In addition, with the simple modifications described in section 5.6, many nonlinear differential equations can be solved simply by evaluating the Fréchet derivative of the differential operator analytically and forming equation (5.6.6) beforehand.

The use of the automatic differential equation solving programs will also be illustrated in this chapter by solving numerous examples from linear and nonlinear ordinary differential equations and from linear two-dimensional partial differential equations. Whenever possible, the results from the programs will be checked with analytical solutions or with previously computed answers. These comparisons will indicate the extremely accurate solutions obtained by using high-order polynomial matrix equivalents for differential operators.

The development of convenient computer languages for differential equations has attracted considerable interest in recent years. The main impetus behind this movement has been a desire to free the numerical solution of differential equations from cumbersome Fortran programming. However, since all of the languages developed to date have been tied to inefficient finite difference formulae, none of these languages has received widespread acceptance.

For initial value problems, a formal computer language and associated program compiler has been designed by Barton, Willers and Zaher [61]. With this system, commands such as INTEGRATE, WITH INITIAL CONDITIONS, END and the mathematical symbols + - * / + ' define the operations to be performed by the computer. As described in [61], the language is capable of handling arbitrary systems of ordinary differential equations for which initial conditions are known.

A language has also been written for partial differential equations, this one by Cardenas and Karplus [62]. Called PDEL, the language is a superset of PL/1 and may be mixed

with it. There are ten functionally different types of statements in PDEL and the syntactic definition of each has been designed to simplify coding procedures for differential equations. At the time reference [62] was written, the PDEL translator was able to handle five different types of differential equations.

A second language for partial differential equations has been written by Roberts and Boris using a symbolic style of Algol [63]. In this language, real procedures such as CURL, CROSS and DELSQ are defined which produce the numerical operations corresponding to their analytic or algebraic counterparts. By providing a library of such procedures, Symbolic Algol allows three dimensional partial differential equations to be written in a form very closely resembling the formalism of mathematical physics.

6.1 SIMPLIFIED INPUT AND OUTPUT PROCEDURES

In this section, computer programs are described which provide a convenient engineer-to-computer interfacing for the programs given in Chapters 4 and 5. With these additions, the operation of the differential equation solving programs in this thesis is made extremely easy: the user simply writes the differential equation and associated boundary conditions on a standard keypunch or teletype terminal in easily recognizable mathematical form and the computer automatically returns its general and particular solutions in both analytical and graphical forms.

The input portion of the new programs is based upon a new

symbolic computer language for the representation of differential equations. Called DECL for Differential Equation Computer Language, this language has been designed to maintain the free-form writing of differential equations as far as is possible on a computer. Accordingly, the elementary units of the language are taken to be the standard ASCII 60 character set, with numbers obeying the same syntax rules as numbers in Fortran. The characters X, Y and Z are reserved to represent the dependent and independent variables of differential equations in the following way

one dimension
<independent variable>::=X
<dependent variable>::=Y

two dimensions

<independent variables>::=X|Y

<dependent variable>::=Z

The metalinguistic notation used here is that of reference [64]. When any of the remaining letters is used by itself, it denotes a numerical constant

<constant>::= < any individual letter except X,Y or Z>
In DECL, multi-literal names are not allowed.

Functions are represented in DECL by the following character strings

```
<sin x>::= SIN(X)
<cos x>::= COS(X)
<log x>::= LOG(X)
<exp x>::= EXP(X)
<auxiliary functions>::= FN1(X) \sim FN5(X)
<x<sup>number</sup>>::= POW<number>(X)
```

The use of the auxiliary functions FN1 to FN5 has already been explained in Chapter 4. The reason for adopting the above notation for exponentiation is that DECL, in contrast to such number manipulation languages as Fortran and Algol, is a symbol manipulation language. Thus, while the double star in the Fortran expression X**2 or the up arrow in the Algol expression X+2 means that the numerical value of X is to be squared, the DECL expression POW2(X) refers to the functional operation of multiplying the function X with itself. Accordingly, the form of exponentiation in DECL is chosen to be the same as for other functional operations with the function preceding its argument.

In DECL, the arguments of functions may be composed of any combination of the independent variable(s), constants, numbers, other functions, and the symbols +-*(). In evaluating a function with a compound argument, DECL follows the same structural conventions as Fortran. Note that the symbol / may not appear in a functional expression in DECL and, as a result, division can only be obtained by applying the function POW-1.

Differential operators are specified in DECL by the following expressions

$$< \frac{d}{dx} \text{ or } \frac{\partial}{\partial x} > ::= D/DX$$

$$< \frac{\partial}{\partial y} > ::= D/DY$$

$$< \frac{d^{n}}{dx^{n}} \text{ or } \frac{\partial^{n}}{\partial x^{n}} ; n = 2, \dots, 9 > ::= D < n > / DX < n >$$

$$< \frac{\partial^{n}}{\partial y^{n}} ; n = 2, \dots, 9 > ::= D < n > / DY < n >$$
valid DECL expression may follow a differential operator

symbol.

Any

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A DECL statement is a string of valid DECL expressions in a predefined and proper order. DECL statements may be written anywhere on an 80 column computer card with the first character appearing in column one. Blanks may be inserted between any two DECL expressions in a DECL statement as long as they do not alter the appearance of the expressions themselves.

There are six functionally different statements in DECL:

(1) Control Statement

<control statement>::=*<number><title>

The number in a control statement may be any integer from 1 to 9 and specifies the order of the approximating polynomial to be used in the solution. The characters given in the title will be used to provide a heading for the output. An example of a control statement is

* 9 TITLE

(2) Differential Equation Statement

<differential equation statement>::=

DE:<any valid DECL differential equation> A DECL differential equation is any DECL expression containing at least one differential operator and one and only one equal sign. Differential equation statements may be continued on consecutive computer cards provided that the letters DE:begin each one of them. A typical differential equation statement is

DE: D2/DX2(Y) = K* SIN(X) * Y

(3) Boundary Condition Statement <boundary condition statement>::=

BC:<any valid DECL boundary condition>

A DECL boundary condition is any DECL expression containing one equal sign in which each of the dependent and independent variables is followed by a number or pair of numbers in parenthesis. These numbers provide the coordinate values of the boundary condition locations. An example of a boundary condition statement is

BC: D/DX(Y(0.0)) + EXP(1.5) + Y(1.5) = P

(4) Constant Specification Statement

<constant specification statement>::=

CST:<constant>=<number>

Constant specification statements are used to provide numerical values for known constant parameters in differential equation or boundary condition statements. A typical constant specification statement is

CST:P=3.141592

Any constant in a differential equation statement which does not have its value defined by a constant specification statement is assumed to be a symbol for an eigenvalue. Although they may appear repeatedly in a differential equation statement, only one eigenvalue symbol may be used per differential equation. Eigenvalue symbols are not allowed in boundary condition statements.

(5) Geometry Statement

one dimension

<geometry statement>::=

X=<number>,<number>(<integer>ELEMENTS)

. 138
two dimensions

<geometry statement>::=

REC: X=<number>,<number>(<integer>SUBDIVISIONS),

Y=<number>,<number>(<integer>SUBDIVISIONS) In one dimension, the geometry statement specifies the endpoints of a line segment to be used in the solution and provides the number of equal elements to be taken in that segment. In two dimensions, it specifies the edges of the rectangular region to be used and the number of equal elemental subdivisions in each direction. If the number of elements or the number of subdivisions to be used is not specified, one element or one subdivision is assumed. Unequal one-dimensional elements or irregular two-dimensional regions...may be defined by using more than one geometry statement. Examples of geometry statements are

X=0.0, 1.0 (2 ELEMENTS)

REC:X=0.0, 1.0 (2 SUBDIVISIONS), Y= 0.0,2.0

(6) Output Statements

regular problems

<output statement #1>::=WRITE<parameters>SOLUTIONS
<output statement #2>::=PLOT<parameters>SOLUTIONS

where

<parameter #1>::=GENERAL
<parameter #2>::=AND
<parameter #3>::=PARTICULAR
eigenvalue problems
<output statement #1>::=WRITE<number>EIGENVECTORS
<output statement #2>::=PLOT<number>EIGENVECTORS

Output statements control the manner in which the results are displayed by the computer. Typical output statements are

WRITE GENERAL AND PARTICULAR SOLUTIONS

- PLOT PARTICULAR SOLUTION
- WRITE 10 EIGENVECTORS
- PLOT 2 EIGENVECTORS

A DECL program consists of DECL statements in the order given above. Each program must contain one control statement, one differential equation statement (which may be continued on any number of cards), any number of boundary condition statements, the number of constant specification statements required by the differential equation and boundary condition statements, one or more geometry statements and two output statements. For programming convenience, the two output statements may be replaced by a single blank card. In this case, the default option of the program is to write and plot both the general and particular solutions of ordinary differential equations, write and plot only the particular solution of partial differential equations, and to write and plot the first five eigenvectors of eigenvalue problems. Examples of valid DECL programs will be given in the following sections.

Computer programs which read boundary value problems written in DECL and translate them into instructions for the differential equation solving programs given earlier are presented in Figures 6.1.1 and 6.1.2. The first of these is the Fortran subroutine INSTR for use with ordinary differential equations and the second is INST2D for use with two-dimensional partial differential equations. These compilers are written in Fortran instead of a machine language so that they are machine

Figure 6.1.1a

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DECL PTD

TEC1 =17

1911 913 2401 913

DEC1 805

DECL 101 DECL 335

2801 813

DECL F-5

DECL REC

0601 8÷3

1501 F 1 1501 F 1 1501 F 5

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SUBROUVINE INSTRIINSTRICONSTINUMDE, XBC, XD, WT, NELMTS, DECL 340 NPROD = 0* ANGE, NOOM, LAMDA, RITEP, RITEG, PLOTP, PLOTG, NEIGWR, NEIGPL) DECL 345 LASIGN = 0 DECL 350 NSKIP = 0THIS SUBREUTINE RELES DIFFERENTIAL EQUATIONS, BOUNDARY COND-DECL 355 ISIGN = 1ITICHS AND GEOMETRICAL DATA AND CONVERTS THEM INTO CODED DECL 360 LAMDA = .FALSE. DECL 365 EVALUE = BLANK INSTRUCTIONS. THE DIFFERENTIAL EQUATION AND COUNDARY CONDITIONS DECL 370 RSIDE = .FALSE. FCHCND = .FALSE. DECL 375 ARE REFURNED IN THE ARRAY 'INSTR' AND THE POINT LOCATIONS ARE PARCND . . FALSE. DECL 380 STORED IN IXDI. ARGUE = .TRUE. DECL 385 WRITE(6,100) **DECL 390** INTEGER#2 INSTR DECL 395 GO TO 1 DIMENSION INSTR(1); CONST(1); XBC(1); XD(1); WT(1) DECL 400 C DI'ENSIGN ARRAY(80), ISKIP(100), THERE(100), TPAR(100), TYPE(8) DECL 405 C A BOUNDARY CONDITION CARD HAS BEEN READ LOGICAL RSIDE, YCND, FONOND, PARCND, ARGUE, LAMDA, RITEP, RITEG, PLOTP, DECL 410 C ₽LOTG,CHARTE DECL 415 13 IF(PRCARD .EQ. STAR .DR. PRCARD .EQ. EX) GD TD 990 COMON ARRAY, ISKIP, IHERE, IPAR DECL 420 WRITE(6,106) HELL LARACAPPINUS DECL 425 INDEX = INDEX + 1 EQUIVALENCE (TYPE(1),STAR), (TYPE(2),DEE), (TYPE(3),BEE), DECL 430 INSTR(INDEX) = 4000* (YPE(4), GEE), (TYPE(5), EX), (TYPE(6), DBLU), (TYPE(7), DECL 435 NUMDE = INOEX PEA), (TYPE(8), BLANK) DECL 440 14 INDEX =INDEX + 1 DATA BLANK, LBRACK, RBRACK, PLUS, MINUS, STAR, SLASH, EQUAL DECL 445 INSTR(INDEX) = 3000 * DECL 450 IF(+NOT, RSIDE) GD TO 990 DATA BEE/CYE, DEE, EEE/EF / GEE/EYE, EL , EN , PEAJOH JARE/ES, DBLU/EX , DECL 455 RSIDE = .FALSE. * 2 m V DECL 460 C /103/1+C,1+D,1+E,1HF,1HG,1HI,1HL,1HN,1HP,1HD,1HR,1HS,1HW,1HX, DECL 465 C × A DIFFERENTIAL EQUATION CARD HAS BEEN READ IriY Z * DECL 470 C DECL 475 15 YCHD = .FALSE. READ THE DATA CARDS IF(IPARA .NE. C .DR. FCNCND .DR. PARCND) GD TD 990 DECL 480 PRCARD = AKRAY(1) DECL 485 1 RELD(9,110,END=999) ARRAY DECL 490 INDEX = INDEX + 1INSTR(INDEX) = 1000 DECL 495 GO TO THE LECATION INDICATED BY THE TYPE OF CARD DECL 500 IATS = - 1 DECL 505 WRITE(6,102) (ARPAY(I), =4,80) 00 2 I = 1,8 DECL 510 C IF(ARRAY(1) .EQ. TYPE(I)) GD TD (3,15,13,4,7,82,84,86),I DECL 515 C III INDICATES THE LOCATION ON THE DATA CARD 2 CENTINUE DECL 520 C GO TO 990 DECL 525 1 - 3 3 FRCARD # STAR DECL 530 16 I = I + 1DECL 535 C DETERMINE THE ORDER OF ELEMENT REQUESTED AND PRINT THE TITLE DECL 540 C IF SOME LOCATIONS WERE SKIPPED ON THE DATA CARD, RETURN TO THEM DECL 545 C HCOM = IFIX(FNUM(ARRAY/2/11) + +1) IF(ND2.1 .GE. 10) GD TD 990 IF(IL .GT. 80) IL = 80 DECL 550 IF(NSKIP .EO. 0) GD TD 58 DECL 555 DO 57 M = NSTART, NPROD DECL 560 IF(1 .NE. ISKIP(M)) G0 T0 57 ARITE(0,111) DECL 565 I = IHERE(M)ARITE(6,102) (ARRAV(I),I = 11,80) DECL 570 IF(ARRAY(I = 1) .E2. MINUS) ISIGN = -1 # ISIGN DECL 575 57 CONTINUE INITIALIZE THE PARAMETERS-DECL 580 C DECL 585 C BLANKS AND ASTERISCS ARE IGNORED INCEX = 0 DECL 590 C CSTS = 0DECL 595 58 IF(ARRAY(I) .EQ. BLANK) GD TD 18 ASCPT = 0 DECL 600 IF(ARRAY(I) .EQ. STAR) GO TO 54 IPARA = 0DECL 605 INDEX = INDEX + 1

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Figure 6.1.1b

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      CHECK FOR: !(!
      IF(ARFLY(I) .NE. LBRACK) GD TD 56
      IP_RA = IP_PA + 1
      INSTR(INDEX) = 300
      IF (FCNCND .CR. PARCND) GD TO 54
      M = 1
      N = IPARA
   59 M = M + 1
      IF (ARR.Y(M) .ED. WHY .OR. ARRAY(M) .EQ. SLASH) GO TO 60
      IF(ARRAY(M) .EC. LERACK) N = N + 1
      IF(ARRAY(M) .EQ. RERACK) N = N - 1
IF(N .3E. IPARA) GO TO 59
      IF(FONEND) GD TD 54
      FCNCND = .TRUE.
ARSUE = .FLLSE.
      IFPARA = IPARA
      60 TO 54
   60 PARCHO . .TRUE.
      IATP = 1
      63 TO 34
С
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                  171
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   55 IF(ARRAY(I) .NE. RBRACK) GD TO 17
      IPARA = IPARA = 1
      INSTR(INDEX) = 400
      GG TO 55
000
                  TYT
   17 IF(ARRAY(I) .NE. WHY) GD TD 22
      IF(FENEND) GD TD 990
      YCHD = TRUE.
      ARGUE - .THUE.
      INSTR(INDEX) = 200
      IF(ARRAY(1) .Eg. DEE) GD TO 55
NBCPT = NBCPT + 1
      ABC(NBCPT) = FHUM(LRRAY, 12, 13)
      I = I3
      INDEX = INDEX + 1
      INSTR(INDEX) = - (300 + NBCPT)
      GO TO 55
C
                  1 X 1
С
٤
   22 IF(ARRAY(I) .NE. EX) GD TO 23
      INSTR(INDEX) = 100
      ARGUE = .TRUE.
      60 TO 54
С
C
                  1±1
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DEC1. 880 C			DECLIISC
DECL 885	23	IF(ARRAY(I) .NE. EQUAL) GO TO 24	DECL1155
DECL 890		IF(PARCND) GD TD 50	DECL1150
DECL 895		IF(RSIDE) GD TO 990	DECT1165
DECL 900		RSIDE = .TRUE.	DECL1170
DECL 905		ARRAY(I) = PLUS	DECLII/P
DECL 910 C			DECLI:00
DECL 915 C		1+1 UR '-'	02021100
DECL 920 C	-	THE WAR AND ADDRESS OF TO 33	DECL1190
DECL 925	24	IF (ARRAY(I) .NE. PLUS .AND. ARRAY(I) .NE. MINUS) 60 10 32	120111-0
DECL 930	50	IF(,NET, ARGOE) GD ID 990	02021205
DECL 935			DECI 1210
DECL 940		1037R(1)DEX = 500	DECL1215
DECL 940		GO TO SS	DECL1220
DEC1 955	25	IE (NOT YOND CR. TPARA .NE. 0) GD TO 26	DECL1225
DECL 960		YCND = FALSE.	DECL1230
DECL 965 C			CECL1235
DECL 970 C		IF THE + DR - IS WITHIN A PARENTHESIS, THE REMAINING TERMS NEED	DECL1240
DECL 975 C		TO BE SKIPPED	DECL1245
DECL 980 C			DECL1250
DECL 985	26	IF(.ND1, PARCND) GO TO 30	DECLIZES
DECL 990		INDEX = INDEX - 1	DECLIZES
DECL 995		IF(IPARA .EQ. C) GD TD 63	DECLIZED
DECL1000		IF(NSKIP .;;E. 0) GD 10 28	05011270
DECLIDOS	27	NPRUD = NPRUD + 1	DECL12/0
DECLIGIO			02011200
DECLIDID		IDSP(A)OOD = I + I	DEC11290
		IAR(NYRDY) = IFANA	0=011295
DECL1020		GU TO 62	DECLIBOD
DECI 1035	28		DECE1305
DECL1040		1F(I .LT. 15K1P(N)) GD TD 61	DECL1310
DECL1045	29	CONTINUE	DECL1315
DECL1050		N = N + 1	DECT1350
DECL1055	61	LASIGN = N = 1	DECL1325
DECL1060		ILASIN a I	DECL1330
DECL1065	62	I = I + 1	02CL1333
DECL1070 C			02001340
DECL1075 C		LOCATE THE END OF THE PARENTHESIS CONDITION	DECL 1342
DECLIOSO C			02011330
DECLIO85		$1 + \{1, 0\}, \{7\}$ but $0, 790$ 1 + (1, 0), (7) but $0, 7901 + (1, 0), (7)$ but $0, 7001 + (1, 0), (7)$ but $0, 7001 + (1, 0), (7)$ but $0, 7001 + (1, 0), (7)1$	DEC11340
DECLIONO		$IF(ARRAY[1]) = U_{2} = LORAC(A) IPARA = IPARA = 1$	n=[1]345
DECT1040		$\frac{1}{1} \left[\left(A R A A Y \right) \right] = \frac{1}{1} \left[A R A A Y \right] = \frac{1}{1} \left[\left(A R A A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left(A R A Y \right) \right] = \frac{1}{1} \left[\left($	DECI 1370
DECI 1105		$\frac{1}{10} \frac{1}{10} \frac$	DECL137
DECL1110	63	IF(LASIGN .FO. 0) CD TD 65	DECLIBE
DECI 1115 C	•••		DECL1385
DECL1120 C		DETERMINE THE SIGN OF THE CURRENT SIMPLE OPERATOR	DECL1390
DECL1125 C			DECLI39
DECL1130		IF(NSKIP .EQ. 0) NSKIP = 1	DECL1400
DECL1135		NSTART = NPROD - NSKIP + 1	DECL1405
DECL1140		$I = I \Delta T S$	DECL141
DECL1145		IHERE(LASIGN) = ILASIN + 1	DECL141
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Figure 6.1.1c

CONST(NCSTS) = FNUM(ARRAY, 13, 14) DECL1420 IF(LASIGN .EQ. NPROD) GO TO 55 IF(14 .EQ. 13) GD TO 990 DECL1425 NSKIP = NSKIP + 1I = I4 = 1DECL1430 M = LASIGN + 1GO TO 53 DECL1435 00 64 K = 1,11PR00 DECL1440 C 64 IHERE(P) = ISKIP(M) ISTR! DECL1445 C 65 \5.1P = 0 DECL1450 C 42 IF(ARRAY(I) .NE. ES .OR. ARRAY(I1) .NE. EYE DECL1455 * .OR. ARRAY(12) .NE. EN) GO TO 43 NPR05 = 0 DECL1460 INSTR(INDEX) = 11 ISIGI = 1DECL1465 PARCHO . FALSE. GD TD 52 DECL1470 INDEX = INDEX + 1 DECL1475 C IF (ARRAY(I) .EC. EQUAL) GD TD 23 10051 DECL1480 C 30 IF(INS(R(1:CEX - 1) .EQ. 1000) INDEX = INDEX - 1 DECL1485 C 43 IF(ARRAY(I) .NE. CEE .OR. ARRAY(I1) .NE. OH IATS = IDECL1490 * . UR, APRAY(12) .NE, ES) GD TD 44 INSTR(INDEX) = 1000 DECL1495 INSTR(INDEX) = 12· C SIGNS ON THE RIGHT SIDE OF AN EQUATION ARE REVERSED DECL1500 GD TD 52 DECL1505 :0 DECL1510 C IF(.NOT, RSIDE .AND, ARRAY(I) .EO. PLUS) GO TO 55 'EXP' DECL1515 C IF(RSIDE ARRLY(I) .EQ. MINUS) GD TD 55 DECL1520 C INSTR(INDER) = 2000 44 IF(ARRAY(I) .NE. EEE .OR. ARRAY(I1) .NE. EX DECL1525 * .DR. APRAY(12) .NE. PEA) GD TO 45 ISIGN = - 1 * 1515% DECL1530 INSTR(INDEX) = 13 60 TO 55 DECL1535 32 ll = I + 1 GD TO 52 DECL1540 12 = 11 + 1DECL1545 C 13 = 12 + 1'LOG' DECL1550 C I4 = I3 + 1 · DECL1555 C 45 IF(ARRAY(I) .NE. EL .OR. ARRAY(I1) .NE. OH IF(14 .GE. 80) GG TO 49 DECL1560 # .OR. ARPAY(12) .NE. GEE) GO TO 46 C DECL1565 1NSTR(INDEX) = 141D/DX1 . с DECL1570 GD TD 52 0 DECL1575 IF(ARRAY(I) .NE. DEE .OR. ARRAY(I1) .NE. SLASH * .DR. LPRAY(12) .NE. DEE .OR. ARRAY(13) .NE. EX) GO TO 33 DECL1580 C IEN! DECL1585 C 1F(FCNCND) 60 TB 990 DECL1590 C 46 IF(ARRAY(I) .NE. EF .DR. ARRAY(I1) .NE. EN) GD TD 49 INSTR(INDEx) = - 1 DECL1595 I1 = IFIX(FNUM(ARRAY,I2,I3) + .1) IF(I3 .NE. I2 + 1 .0R. I1 .GT. 5) GD TO 990 I = 13 DECL1600 GO TO 54 DECL1605 INSTR(INDEX) = 14 + 11 j.c DECL1610 1DN/DXN1 GD TO 52 DECL1615 ; C DECL1620 C 33 IF(ARRAY(I) .ME. DEE .DR. ARRAY(I2) .NE. SLASH IA NUMBER! .02. AFRAY 13) .NE. DEE .OR. ARRAY (14) .NE. EX) GD TO 36 DECL1625 C * DECL1630 C IF(FONCND) GD TO 990 49 NCSTS = NCSTS + 1DECL1635 INSTR(INDEX) = - IFIX(FNUM(ARRAY, 11, 12) + .1) ARGUE = .TRUE. DECL1640 IF(IZ .NE. 11 + 1) GO TO 990 INSTR(INDEX) = - (100 + NCSTS)DECL1645 1 = 14 + 1CONST(NCSTS) = FNUM(ARRAY, I, I1) DECL1650 Gg Tg 54 IF(I1 .EQ. 1) GD TD 47 DECL1655 C 1 = 11 = 1DECL1660 ١č 100/1 GO TO 54 DECL1665 DECL1670 C C 35 IF(ARR_Y(I) .NE. PEA .OR. ARRAY(I1) .NE. DH WHATEVER REMAINS MUST BE A CONSTANT DECL1675 C # .UR. LRRAY(I2) .NE. DBLU) GD TO 42 DECL1680 C NCSTS = NCSTS + 147 IF(FCNCND) GD T0 990 DECL1685 INSTR(INDEX) = - (200 + NCSTS)

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DECL1925 DECU1930 DECL1935 2E2119-0 DECU1945 DECL1950 DECL1955 4

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DECL1690

DEC11695

DECL1700

DECL1705

DECLIPIC

DECL1715

04CL1720

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CECL1895

DECL1492

DECLIBRE

DECL1900

DECL1905

DECL1910

DEC11915

DECLIPZO

Figure 6.1.1d

INSTR(INDEX) = 800 XD(INDEX) = ARRAY(I) ARGUE = .TRUE. GD TO 54 52 i = 12 53 ARGUE = .FALSE. IF(FONUND) GO TO 55 FCNCND = .TRUE. IFPARA = IPARA 54 IF(YOND) GD TO 990 55 CONTINUE 18 IF(I .LT. 80) 60 TO 16 6<u>9</u> TO 1 EVALUATE THE CONSTANT SPECIFICATION STATEMENT 4 IF(PRCARD .EQ. CEE) GD TD 5 IF(PRCIRD .EQ. STAR .OR. PRCARD .EO. EX) GD TD 991 PR04RD = CEE ARITE(0,107) 5 30 75 1 = 5,80 IF(ARRAY(I) .NE. BLANK) GD TO 76 75 CONTINUE GE TE 991 75 II = I + 1VALUE = FNUM (ARRAY, 11, 12) IF(11 .E2. 12) GD TO 991 ARITE(0,105) (ARRAY(13),13=1,12) 02 77 13 = 1, 1+ DEX 1F(XD(13) .EQ. ARRAY(13)) GD TD 78 77 CONTINUE SO TO 1 78 NCSTS = NCSTS + 1 INSTR(13) = -(100 + NCSTS)CENST(NCSTS) = VALUEGD TO 77 EVALUATE THE GEOMETRY STATEMENT 7 IF(PRCARD .EQ. EX) GD TD 71 IF(PRCARD .EQ. STAR) GD TD 991 DETERMINE IF THE PROBLEM IS AN EIGENVALUE PROBLEM 00 79 I = 1/INDEX IF(INSTR(I) .Eg. 800) GO TO 80 79 CONTINUE GO TO 91 BO IF(LAMDA .AND. XD(I) .NE. EVALU) GD TO 991 IF(LAMDA) GD TO 79 LANDA = .TRUE. EVALU = XD(I)GO TO 79 51 IF(INDEX .LE. 2) GB TO 991

NUMDE = NUMDE + 1 DECL1960 DECL2230 DFCL1965 INDEX = INDEX + 1DEC12735 INSTR(INDEX) = 4000**DECL1970** DECL22-C WRITE(6,104) (I, INSTR(1), I = 1, INDEX) DECL1975 DECLOSAS DECL1980 8 I = 1 DECL2250 NRANGE = NOOM DECL1985 DECL2255 DECL1990 9 N = 0DEC12250 DECL1995 C DECL2265 DECL2000 C DETERMINE THE ORDER OF THE RANGE SPACE DECL2270 DECL2005 C DECL2275 DECL2010 10 I = I + 1DECLZZED IF(INSTR(I) .GE. 10) GD TO 11 DECL2015 DEC12265 IF(INSTR(I) .LE. -10) GD TO 11 DECL2020 DECLEZED N = N + INSTR(I)DEC12025 DEC12275 DECL2030 11 IF(INSTR(I) .LT. 1000) GU TO 10 DECLIZED DECL2035 NRANGE = MINO(NRANGE>NDDM + N) DEC12305 IF(INS R(I) .LT. 4000) GD TO 9 DECL2040 DECL2310 IF(NRALGE .GE. 0) GO TO 12 DECL2045 DECL2315 NDOM = NDOM - NRANGE DECL2050 DECL2320 DECL2055 NRANGE = 0 DEC12775 DECL2060 12 WRITE(6,105) NODP, HRANGE **DECT2330** DECL2065 NELMTS = 0 DEC12335 NPT = 0DECL2070 DECL23-0 71 UD 72 i = 2,80 DECL2075 DECL27-5 DECL2080 IF(ARRAY(I) .ED. EQUAL) GD TD 73 DECL2350 72 CONTINUE DECL2085 DECL2355 DECL2090 GD TO 991 DECL2350 DECL2095 73 I = I + 1DEC12355 DECL2100 C DEC12370 DECL2105 C FIND THE LOCATIONS OF THE ENDPOINTS OF THE ELEMENTS DECL2375 DECL2110 C DECL2320 DECL2115 X1 = FRUM(ARRAY)[]]DECL2385 DECL2120 IF(I1 .EQ. I) GD TD 991 DECL23=0 DECL2125 IF (NELITS .NE. O .AND. X1 .NE. X2) GO TO 991 DECL2395 DECL2130 I = I1 + 1DECL2400 DECL2135 XZ = FHUM(ARRAY)IIII)DECL2405 DECL2140 IF(I1 .EQ.]) GO TO 991 DECL2410 DECL2145 I = I1 + 1DECU2415 DECL2150 C DECL2420 DECL2155 C READ THE NUMBER OF ELEMENTS IN THE INTERVAL DECLZ425 DECL2160 C DECL2430 DECL2165 NEL = IFIX(FNUH($\Delta RRAY$, I, I) + .1) DECL2435 DECL2170 IF(NEL LE. 0) NEL = 1 DECLESSO DECL2175 XLI:GTH = (x2 - X1)/NEIDECL24-5 DECL2180 STEP = XLNGTH / (NDOM - 1) DECL2450 XPT = X1DECL2185 DECL2455 DECL2190 C DECLASED DECL2195 C DETERMINE THE INTERPOLATION POINTS DEC12-55 DECL2200 C DECL2470 DD 74 I = 1/NEL DECL2205 DECL2475 DECL2210 NELMTS = NELMTS + 1 02012-80 DECL2215 wT(NEL:ITS) = 1.0 / XENGTH DECL2485 DECL2220 XPT = APT - STEPDECL2490 DECL2225 90.74 I1 = 1, NDDECL2495 ____

Figure 6.1.1e

74	XPT = XPT + STEP NPT = NPT + 1 , XD(NPT) = XPT FRCARD = EX GC TO 1	DECL2500 DECL2505 DECL2510 DECL2515 DECL2520 DECL2525	<pre>104 FORMAT(1H-,54X,'THE INSTRUCTION CODE IS'//(1H ,25X,5(I10,J4))) 105 FORMAT(1H1,27X,'THE POLYNOMIAL ORDER OF THE DOMAIN SPACE IS'; * 12,' AND THE ORDER OF THE RANGE IS',I2) 106 FORMAT(1H-,51X,'THE BOUNDARY CONDITIONS ARE'//) 107 FORMAT(1H0,52X,'CONST SPECIFICATION LIST'//) 108 FORMAT(56X,75A1)</pre>	DECL2763 DECL2755 DECL2775 DECL2775 DECL2775 DECL2755 DECL2755
83 83 54	DEFINE THE CUTPUT PROCEDURE : IF(PRCARD .NE. EX) GD TO 991 PRCARD = DSLU IF(LWOA) GC TO 83 RITEP = CHARTF(ARRAY,PEA) RITEG = CHARTF(ARRAY,PEA) RITEG = CHARTF(ARRAY,GEE) GO TO 1 NEIGWR = IFIX(FNUM(ARRAY,6,I) + 0.1) IF(I .EQ. 6) GO TO 991 GC TO 1 : IF(PRCARD .NE. DBLU) GU TO 991 : IF(PRCARD .NE. DBLU) GU TO 991	DECL2530 DECL2535 DECL2540 DECL2545 DECL2555 DECL2555 DECL2560 DECL2565 DECL2565 DECL2575 DECL2575 DECL2580 DECL2585 DECL2585	109 FORMAT(I-ERROR:') 110 FORMAT(BOA1) 111 FORMAT(1H1) END	DECL2790 DECL2795 DECL2800 DECL2805
85	PLOTP = (PLRTF(ARRAY,PEA) PLOTG = CHIRTF(ARRAY,GEE) GD TD 85 (NEIGPL = HINO(IFIX(FNUM(ARRAY,5,I) + 0.1),NEIGWR) IF(I .EQ. 5) GD TD 991 IF(I .EQ. 5) GD TD 991	DECL2595 _ DECL2600 DECL2605 DECL2610 DECL2615 DECL2615	FUNCTION CHARTE(A+C)	DEC16555
85	GITU GA IF(PCARD .NE. EX) GD TO 991 IF(LAMDA) GD TO 87 RITEP = .TRUE. RITEG = .TRUE. PLOTP = .TRUE. PLOTP = .TRUE. GO TO do	DECL2625 C DECL2635 C DECL2635 C DECL2640 C DECL2640 C DECL2650 DECL2655	CHARTF DETERMINES WHETHER OR NOT CHARACTER 'C' IS CONTAINER IN THE ARRAY 'A'. LOGICAL CHARTF DIMENSION A(1) CHARTF = TRUE.	DECL6590 DECL6595 DECL6603 DECL6603 DECL6610 DECL6613 DECL6623
87	NEIGWR = 5 NEIGPL = 5 WRITE(====================================	DECL2660 DECL2665 DECL2670 DECL2675 DECL2685 DECL2685 DECL2690 DECL2695	DD 1 I = 1,80 IF(A(I) .EQ. C) RETURN 1 CONTINUE CHARTF = .FALSE. RETURN END	DECL6630 DECL6630 DECL6635 DECL6645 DECL6645 DECL6650
90 91	WRITE(6,103) ARRAY(I), I, IPARA, RSIDE, YCND, FCNCND, PARCND, ARGUE WRITE(6,104) (I, INSTR(I), I=1, INDEX) GD TO 1 WRITE(0,109) WRITE(0,102) ARRAY GD TO 1	DECL2700 DECL2705 DECL2715 DECL2715 DECL2720 DECL2725 DECL2725 DECL2725		
99 00 01 02 03	STOP STOP FGRMAT(1H-,51X'THE DIFFERENTIAL EQUATION IS'//) FORMAT(1H-,55X,'THE POINT COORDINATES'//) FORMAT(1H0,25X,80A1) FORMAT(1-ERROR IN DIFFERENTIAL EQUATION: ',A1,7I3)	DECL2735 DECL2740 DECL2745 DECL2750 DECL2755	·	

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Figure 6.1.2a

	SUBROUTINE INST2D(INSTR,CONST/NUMDE/X8C/Y8C/X0/YD/WTX/WTY/	DE2D 20	b c	DECL 1140-1570	DE2D 470
	* IELCX, IELCY, NCCOND, NELMTS, NRANGX, NRANGY, NDDM, LAMDA, RITEP, I	DE2D 20	5Č		DE2D 47
	* RITEG, PLOTP, PLOTG, NEIGWR, NEIGPL)	DE2D 21	0	* ',OR, ARRAY(12) ,NE, DEE) GD TO 36	DE2D 480
C	r	DE2D 21	5	IF(FENEND) GO TO 990	DE2D 485
ē	THIS SUBROUTINE READS LINEAR TWO-DIMENSIONAL PARTIAL DIFFERENTIAL [DE2D 22	0	I = 13	DE20 490
ē	EQUATIONS. BOUNDARY CONDITIONS AND GEOMETRICAL DATA AND CONVERTS [DE2D 22	5	IF(ARRAY(13) .NE. EX) GO TO 90	DE2D 49
ē	THEM INTO CODED INSTRUCTIONS.	DE2D 23	0	INSTR(INDEX) = 1	DE20 500
č		DE2D 23	ŝ	GU TO 54	DE20 505
•	01/2ENSTON_VBC/11.VD/11.VETV/11.1ELCV/11.1ELCV/11	DE2D 24	ο c		DE2D 510
	DIMENSION INSTRUITCOUSTINA VBC(1) AD(1) MTX(1)	DE2D 24	5.Č	10/071	DE20 51
	(1) = (1)	DE20 25	ōč		DE20 520
		nE2n 25	5 5	AN TEXAPRAVITAN INE WHYS OF TO 990	DE2D 525
•	¢ (YPE(10)	DE20 24	n i	$\frac{1}{1} \frac{1}{1} \frac{1}$	DE20 630
		DE2D 26	Šr	INDIRITIVERI - 1	DE2D 53
		DE2D 27		0501 1600-1620	DE2D 540
6		0620 27			DE2D 544
	$= \frac{1}{2} \left(\frac{1}{2} \left(\frac{1}{2} \right) \right) \left(\frac{1}{2} \left(\frac{1}{2} \right) \right) \left(\frac{1}{2} \left(\frac{1}{2} \right) \right) \left(\frac{1}{2} \right) \left(\frac$	DE20 27		- DD ADDAV(13) AND DEEN CO TO 34	DE20 55
	T PEALS (TTPE(BIJEX/) (TTPE(J)JWHT/) (TTPE(LU)JDLANK/	0E20 20	6	# .DR: ARKAT(13) (NE: DEE) 00 10 30	0520 550
5		DEAD 20		1 - 14 + 1	DE20 33
6	UELL 4078425	DE20 29	5	4 - 34 - 4 TEXADDAV/17/1 NE - FW1 60 40 01	DE20 50
¢		0EZU 29.		1P(AKRAT(14) .NE. EX) GU 10 91	DE20 304
	* WHYJZEE L	0520 30		NEA: 1678 1774	
Ç		DESD 30		DECC 103541040	
C	DECL 405-405	DE20 31	<u> </u>		
C		0E20 31	2		0620 38
	* 1HY/1HZ /	UE20 32	C C	91 IF (ARRAY(14) .NE. WHT) GU TU 990	DE20 390
C		DEZD 32	2	INSTR(INDEX) = IFIX(FNUM(APRAY)I1)I2) + 0.1)	DES0 29
C	DECL 475-505	DESD 33	0	IF(12 .NE. I1 + 1) GD TD 990	DESD 600
¢		DE20 33	2 C		DE20 60
	DD 2 I = 1,10	DEZD 34	ο C	DECL 1650-2150	DESD 610
	IF(ARKAY(1) ,EQ, TYPE(I)) GD TU (3,15,4,7,82,84,150,150,82);I	DEZO 34	2 C		DE2D 51
C	r i i i i i i i i i i i i i i i i i i i	DESD 35	0	7 IF(PRCARD ,EO, ARE) GO TO 71	DE2D 620
C	DECL 520-925	DE2D 35	5 Ç		DE2D 62
C		DESD 36	O C	DECL 2160-2250	DESD 630
	IF(ARRAY(M) .EQ. ZEE .DR. ARRAY(M) .EQ. SLASH) GD TD 60	DE2D 36	5 C		DE20 63
C	, i i i i i i i i i i i i i i i i i i i	DEZD 37	0	NRANGX = NDBM	DE2D 640
C		DE2D 37	2	NRANGY = NDDM	DE2D 64
C	DECL 935-1025	DEZD 38	0	9 N = 0	DE2D 650
C	121	DE2D 38	2	M = O	DE2D 65
C	1	DE2D 39	o c		DE20 660
	17 IF(ARRAY(I) .NE, ZEE) GD TO 22	DESD 39	2 C	DECL 2265-2290	DE2D 66
C	1	DE2D 40	Č Č		DE2D 670
C	DECL 1050-1115	DE2D 40	5	IF(INSTR(I) .GE, 0) GO TO 92	DE2D 67
C	. 1	DE2D 41	0	N = N + INSTR(I)	DE5D 68
	IF(ARRAY(I) .NE. EX) GO TO 89	DE2D 41	5	GU TO 11	DE20 68
C		DE2D 42	0	92 M = M + INSTR(I)	DE2D 690
C	DECL 1125-1135	DE2D 42	5	11 IF(INSTR(I) .LT. 1000) GD TD 10	DE20 69
C		DE20 43	0	NRANGX = MINO(HRANGX,NDOM + N)	DE2D 70
C C	iγi	DE2D 43	5	NRANGY = MINO(PRANGY=NDDM = M)	DE20 70
Ċ		DE2D 44	0	IF(INSTR(I) .LT. 4000) GU TE 9	DE20 710
	89 IF(ARRAY(I) .NE, WHY) GO TO 23	DE20 44	5	IF(NRANGX ,GE, 0) GO TO 93	DEZC 71
i	INSTR(INDEX) = 101	DE20 45	0	NDON = NDUH - PRAHGX	DE20 72
	ARGUE = .TRUE.	DE2D 45	5	NRAHGY = HRANGY - MRANGX	DE20 72
	62 TO 54	DE20 46	0	HRANGX = 0	DE20 73
C		DE20 46	5	93 IF(NRANGY .GE. 0) 63 TO 12	DE2D 73
-		-		· · · · · · · · · · · · · ·	
					<u>ب</u>
		•			4

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Figure 6.1.2b

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00000	12	NDOM = NDOH = NRANGY NRANGX = NRANGX = NRANGY NRANGY = 0 WRITE(6,105) NDDM,NRANGX,NRANGY ND = NDOM + 1 NDSQ = ND +* 2 WRITE(6,101) DECL 2335-2430
		FIND THE NUMBER OF ELEMENTS IN THE X DIRECTION
	94	NELX = IFIX(FNUH(ARRAY/I/II) + 0.1) IF(NELX .LE. 0) NELX = 1 XLNGTH = (X2 - X1) / NELX DC 94 I = I1/80 IF(ARRAY(I) .EG. EQUAL) GD TD 95 CC"ITINE GD TD 991
C	95	<pre>1 = I + 1 Y1 = FRUM(ARRAY; I; I1) IF(I1 .EQ. I) GD TO 991 I = I1 + 1 Y2 = FRUM(ARRAY; I; I1) IF(I1 .EQ. I) GD TO 991 I = I1 + 1 FRUM ARRAY; I; IN THE M DIRECTION</pre>
C		FIND THE NUMBER OF ELEMENTS IN THE Y DIRECTION
c		NELY = IFIX(FNUM(ARRAY)I/I) + 0.1) IF(NELY *LE* 0) HELY = 1 YLNGTH = (Y2 - Y1) / NELY X2 = X1 = XLNGTH DD 97 J = 1>NELX X2 = X2 + XLNGTH Y1 = Y2 + YLNGTH DD 97 I = 1>NELY Y1 = Y1 - YLNGTH NELMTS = NELMTS + 1 WTX(NELMTS) = 1*C / XLNGTH XTY(NELMTS) = 1*C / YLNGTH XSTEP = XLNGTH / NDDM YSTEP = YLNGHT / NDDM X22 = X2 + XSTEP
Ċ		ADD THE POINTS IN EACH ELEMENT SEPARATELY
č		DD 96 J1 = 1>ND X22 = X22 + XSTEP Y11 = Y1 + VSTEP DD 96 I1 = 1>ND Y11 = Y11 - VSTEP NPT = LPT + 1 XD(NPT) = X22

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DF2D	740		96	YD(NPT) = Y11	DEZDIGÍO
DEZD	745			NPT1 = NPT = NDSQ + 1	DEZDIOIS
DE2D	750	·		WRITE(6,112) NELHTS, (XD(J), J = NPT1,NPT,ND)	DESDICSO
DE2D	755			NPT1 = NPT1 = 1	DE201025
DE2D	760			DO 97 II = 1,ND	DE2D1030
DE2D	765			NPT1 = NPT1 + 1	DE201035
DE2D	770		97	WRITE(6,113) YD(NPT), (J, J = NPT1,NPT,ND)	DE201040
DE2D	775			PRCARD = ARE	DE201045
DE2D	780			GO TO 1	DESDIDSO
DE2D	785	C			DEZUIOSA
DESD	790	C		A CONTINUITY STATEMENT HAS BEEN READ	DE201060
DE2D	795	C			DE201063
DE2D	800		150	IF(PRCARD .NE. ARE) GO TO 151	02201070
DESD	805			PRCARD = EL	02201073
DESD	810			I = 0	DESCION
DE 2D	815			J = 0	02201062
0E2D	820			NCCUNC = 0	06201090
DE2D	825			WRITE(6,114)	05201102
DESD	830		151	IF(PRCARD ,NE, EL) GO TO 991	00201100
DE2D	835			13 = 0	05201103
DE2D	840			DD 152 Il = 6;80	05201110
DESD	845			13 = 13 + 1	05201112
DESD	850			NELC(13) = IFIX(FNUM(ARRAY)I1)I2) + 0.1	05201120
DE5D	855			11 = 12	06201120
DESD	860		152	CONTINUE	05201135
DESD	865			NELC(I3+1) = 0	DE201100
DESD	870			WRITE(6,116) ARKAY(1), NELC	nE201145
DEZD	875			NCCUND = NCCOND = ND	DE201150
DESD	880			INDEX # 1	DE201155
DESD	885			IF(AKKAY(1) SEQS WHY) OU ID 104	DE201160
DEZD	890	ç		NUM THE M ADUTANUSTRY CONDITIONS IN THE ADDAY ITSICY	DF201165
DESD	892	5		POI THE X CONTINUITY CONDITIONS IN THE WARMY TEEN.	DE201179
DESO	900	Ç	1	NCOUND - NCOND : ND	0E201175
DESD	905		123	ACCUMU = ACCUMU + NU	DE201180
0620	910			0 = 0 = 4 16(69/1) - NE(6/INDEX)	DE201185
0520	917			THDEY - THDEY	DE201190
0520	920			$\frac{1}{1} \frac{1}{1} \frac{1}$	DE201195
0620	920			1 - 1 + 1	DE201200
0520	930				DE2D1205
0520	040			100 TO 1	DE20121
06/0	0/5				DE201215
DE 20	050	ž		PUT THE V CENTINUITY CONDITIONS IN THE ARRAY FIELCY!	DE201220
0220	055	ř		tol nucli Continuiti Contritune in the House cont	DE20122
0520	040		154	NCCUND # NCCOND + ND	DE201230
0520	965				DE2D123
0620	970			IELCY(I) . NELC(INDEX)	DE2D1243
0520	975			INDEX = IUDEX + 1	DE20124
DE2D	980			1E(NE) ((INDEX) -NE. 0) GD TO 154	DE20125
DF2D	985			$1 \pm 1 \neq 1$	DE2D125
DF20	290			IELCY(I) = 0	DE271260
CE2D	995			GUTUI	DE20126
DF2C	1000		82	IF(PRCARD .NE. ARE) GD TO 155	DE201273
DE20	1005			1 = 0	DE20127
				-	-

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DE2D1280 J = 0 1 DE2D1285 NCCUND = 0 DE201290 GO TO 156 DE2D1295 155 IF(PRCARD .NE. EL) GD TO 991 DE2D1300 C C UPDN COMPLETION, ADD ENDING VALUES TO 'IELCX' AND 'IELCY' DE201305 DE201310 C DE201315 156 J = J + 1DE2D1320 IELCX(J) = -1DE2D1325 1 = 1 + 1DE2D1330 IELCY(I) = = 1DE201335 IF (ARRAY(1) , EO, BLANK) GO TO 86 DE2D1340 C • DE2D1345 DECL 2545-2800 C DE2D1350 C DE201355 112 FORMAT(1H0,60X, IELEMENT NUMBER 1,13 / 1H0,20X,10G11.3) DE2D1360 113 FURMAT(1H0,9X,G11.3,I6,9I11) 114 FORMAT(1H1 / 1H0,42X, THE FOLLOWING ELEMENTS ARE TAKEN TO BE ADJACDE2D1365 DE2D1370 *ENT! / 1H0,7X, TYPE1, 12X, ELMENTS!//) DE2D1375 116 FORMAT,(1H ,6X, A2, ONT', 3X, 3913) DE2D1380 END

Figure 6.1.2c

independent and can be taken to any computer installation which accepts Fortran programs.

In defining the instructions for a differential equation statement or a boundary condition statement, INSTR and INST2D operate by evaluating the branching structure and syntax tree of the statement by noting the number and locations of the symbols +, -, (, and). The syntax tree is then reduced to the canonical form

 $\sum_{s=1}^{\Sigma} D_s y = f$ (6.1.1)

by applying the distributive law to compound operators enclosed in parentheses. Instructions generated by this procedure are compatible with the subroutines DIFFEQ and DIFF2D.

A separate subprogram FNUM, given in Figure 6.1.3, is called by the subroutines INSTR and INST2D to determine the value of a number stored in Hollerith form in the Fortran language. FNUM establishes the machine-independent, free-format feature of numbers in DECL since it accepts numbers in any Fortran permissible form and in any location.

The diagnostic capabilities of INSTR and INST2D have been designed to detect most errors in DECL that lead to logical inconsistency. For this purpose, a number of parameters and conditions are checked in the programs after each symbol in a statement has been translated. If an error is spotted, the programs write an appropriate message and proceed to the next DECL program.

Conversational versions of the programs INSTR and INST2D for use in time-sharing situations have also been

Figure 6.1.3

DECL6115 FUNCTION FRUM(ARRAY, I, J) DECL6120 DECL5125 C THIS FUNCTION SUPPROGRAM RETURNS THE NUMERICAL VALUE OF A NUMBER STORED IN HOLLERITH FORM IN TARRAY! STARTING IN LOCATION DECLETED C 111 DECL6140 DECL6145 LEGICAL NURBER, MINUS/EXPEN, DECMAL/ONLYE DECL6150 UIMENSION ARRAY(FO), ANUM(15) DECL6155 DATA ALUM / 1H ,1HC,1H1,1H2,1H3,1H4,1H5,1H6,1H7,1H8,1H9, # 1H., 1HE, 1H-, 1H+ / INITIALIZE THE TYPES OF NUMBERS POSSIBLE DECL6175 DECL6180 HUMBER & FALSE. DECL6185 BINUS = .FALSE. DEC16190 EXPON = .False. DECMAL = .; ALSE. LNLYE = .FALSE. . NU¹ = 0 DECL6210 J = 1 DECL6215 DFCL6220 С DETERMINE THE TYPE OF CHARACTER IN LOCATION J DECL6225 DECL6230 1 10 2 8 # 1915 DFCL6235 IF (ARRAY(3) .EG. ANUM(N)) GO TO 8 DECL6240 2 CENTINUE DECL6245 3 IF (EXPUN) GD TO 6 DECL6250 4 IF(HOT, NUPBER) GO TO 990 5 FNLM = NUM С ADD A HINUS SIGN DECL6270 ٢. IF(MINUS) FRUM = - FRUM IPL.NOT. DECMAL) RETURN 60 70 7 5 IF (.NDE, NUMBER) GD TO 5 2 ADD THE EXPONENT С C LPOWER = LPOWER + NUM IF MINUS) LPONER = LPOWER - 2 * NUM - FNUM = FNUM = 10.0 ** LPOWER RETURN g J = J + 1 IF(J .07. 20) GD TO 3 1F(N .EQ. 1) 68 TO 9 IF(N .GT. 11) GD TD 10 IF(DNLYE) GC TO 5 THE CURRENT CHARACTER IS A NUMBER C ċ NUMBER . .TRUE. NUM = 10 = NUM + N - 2 DECL6380 IF(,NOT, DECHAL) GO TO 1

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LPOWER = LPOWER - 1 DEC16390 GO TO 1 CEC:6393 DECLASS BLANKS MAY NOT APPEAR IN A NUMBERS EXCEPT BEFORE AN E 12516-1**1** DECLEA13 9 IF(NUMBER) CHLYE = .TRUE. CECL6413 GE TE 1 55016433 10 N = N - 11DEC16425 GD TO (11,12,13,13),N DECL6-30 DECL6160 C DECLO435 THE DECIMAL POINT IS LOCATED DECL6165 C 350164-3 DECL6170 C 180164-3 11 DECMAL = .TRUE. 18014451 LPOWER = 0 05000-55 69 TO 1 12 IF(EXPUN .OR. .NOT. DECMAL .OR. .NOT. NUMBER) GO TO 990 22016.00 DEC16455 DEC16195 C 2201.5472 THE CURRENT CHARACTER IS AN E DECL6200 C DECL6475 DFCL6205 C CECL6490 EXPON = .TRUE. DECLASS FNUM = NUM DECLOSED IF(MINUS) FNUM = - FNUM DEC16475 MINUS = .FALSE. DECLASCO NUBBER = .FALSE. DECLESIS UNLYE = .FALSE. 0=016510 NU(1 = 0)DECLASIS GG TO 1 0=014520 13 IF (NUMBER) GD TD 3 DEC16535 DECL6255 C DECLESOD THE CURRENT CHARACTER IS A PLUS OR MINUS SIGN DECL6260 C DEC1(595 DECL6265 C DECL65-C IF(N .EQ. 4) GO TO 1 DECLESUS MINUS = .TRUE. DFCL6275 CECLASS1 DECL6280 Go To 1 DECLOSE DEC16285 C DECLEFE? THE INPUT ARRAY DOES NOT CONTAIN A NUMBER DECL6290 C 020165-5 DEC1.6295 C 32014373 DECL6300 990 J = 1 DECLASTS. DEC1.6305 RETURN DECLASED DECL6310 END DECL6315 DECL6320 DECL6325 DECL6330 DECL6335 DECL6340 DEC1.6345 DECL6350 DECL6355 DECL6360 ___ DECL6365 50 DECL6370 DECL6375

DECL6383

written. However, since the programming changes required for this type of operation are minor, the listings of these versions of the programs are not presented here. The modifications are designed to make the input procedures more convenient in on-line operation. For example, the programs will prompt the user with the following statements

> WHAT IS THE DIFFERENTIAL EQUATION? STATE A BOUNDARY CONDITION WHAT IS THE VALUE OF THE CONSTANT P? IN REGION 1, WHAT ARE THE LIMITS OF THE X COORDINATE?

HOW MANY ELEMENTS ARE TO BE USED?

A company and a second

In this way, the operation of the program is made virtually foolproof and requires very little specialized knowledge.

Output programs have also been developed which write and graph the solutions of a differential equation in terms of their numerical values at the interpolation nodes. The program which writes the solutions is called RITE and is given in Figure 6.1.4. This subroutine has already been called by some of the programs in Chapter 5. A program to graph the solutions of ordinary differential equations is presented in Figure 6.1.5. Called GRAPH, the subroutine can plot up to ten functions simultaneously. Finally, Figure 6.1.6 contains a plotting program called PICT for use with two-dimensional partial differential equations. Typical output from these programs is shown in the following sections.

SUBROUTINE RITE(A,M,N,MADDA, IPART) DECL9090 C NECL9095 С THIS SUBROUTINE WRITES THE INF BY INF MATRIX TAF FOR OUTPUT BY DECL9100 C THE PROGRAM. IF 'IPART' IS FALSE, THAN THE LAST COLUMN OF (A) DECL9105 C C IS ASSUMED TO BE AN INHOMOGENEOUS SOLUTION AND IS WRITTEN FIRST DECL9110 DECL9115 LUCICAL IPART DECL9120 DIMENSION A(1) DECL9125 NM1 = N DECL9130 IF(IPART .NE. 0) GO TO 2 DECL9135 WRITE(6,100) DECL9140 NM1 = 11 - 1DECL9145 1ADD = MADDA + M*NM1DECL9150 WRITE(6,101) ($I \neq A(I + IADD)$, $I = 1 \neq M$) DECL9155 IF(NM1 .EQ. 0) RETURN DECL9160 WRITE(6,102) DECL9165 2 IADD = MADDA - MDECL9170 $UD \mathbf{1} \mathbf{J} = \mathbf{1} \mathbf{M} \mathbf{M} \mathbf{1}$ DECI.9175 IADD = IADD + MDECL9180 1 WRITE(ϕ_1 101) (I $\rho\Lambda$ (I + IADD) ρ I = 1 ρ M) DECL9185 RETURN DEC1.9190 100 FORMAT(1H-,53X, 'THE INHOMOGENEOUS SOLUTION'/) DECL9195 101 FORMAT(1H0,100(5x,5(16,E18.6)/1H)) DECL9200 102 FORMAT(1H-,54X, 'THE HOMOGENEOUS SOLUTION'/) DECL9205 END DECL9210

Figure 6.1.4

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Figure 6.1.5

	SUBROUTINE GRAPH(X,Y)N,NPER)	OFCL4225		
		DEC14020		DECL443
	THIS SUBRENTINE GRADUS THE INTERINTS (M. V. DN. & STANDARD CINE	02014230	$12 \operatorname{CHAR}(1) = PLUS$	DECL450
	PRINTED THE NUMBER AND AND A POINTS (AVI) ON A STANDARD LINE	UECL4233	61 = 0	DECLASS
	DE LA CORDENT AND AT BE ANT BULLIPLE OF NPER, THE NUMBER	DECL4240	J2 = C	n=** _**
	UP TAT CODENTIALE POINTS SO THAT SEVERAL FUNCTIONS CAN BE	DECL4245	7 I = I + 1	
	GRAPHED AT DIE TIME.	DECL4250 C		
		DECI 4255 C	TT IS THE NUMBER OF THE EUNCTION CONDUCT	
1	DIHENSION X(1),Y(1),IY(1),IX(500),CHAR(111),ANUH(6),SYMBO(710)	DEC14250 C	I IS THE BOILDER OF THE FONCTION GRAPHED	
	COMMON CHAR, IX; IY	DECL4200 0		DEC: 433
	DATA BLANK VERTINGE PLUS SYNDEL / 18 JULIU JULIUS VID	D2CL4253	1P(1 + GT + 10) I = 1	CEC1453
j		DECL4270	$J_{1} = J_{2} + I_{-}$	DEC1454
	IE(2001) 1002 1002 1002 1002 1002 100	DECL4275	JZ = J2 + NPER	08154F.
	Terror Down and the state of Refurn	DECL4230	JX = 0	
	I T (A(APER) +EQ. X(I)) REFURN	DECL4235	DD 14 J = J1,J2	
		DECL4290	1 + X L = X L	
	THE POINTS ARE SCALED PRUPERLY AND PUT IN THE INTEGER ARRAYS	DECI 4225	IECTVCIN THE LINES OF TO AC	
	1X1 7.0 (IA)	DEC14300 C	citation and friet of in It	DE01-34
		D=C14305 C		060.447
	$SC_{\perp}EX = 110.17(v(MDED) - v(1))$		IF THE F COURDINATE OF THE POINT CORRESPONDS TO THE LINE TO BE	0701451
		DECL4310 C	PRINTED, THE FUNCTEON NUMBER IS PLACED IN THE APPROPRIATE	DECU451
• •		DECL4315 C	LDCATION	2711473
	4 (1476 - 784) 	DECL4320 C		526 223
		DECL4325	CHAR(IX(JX)) = SYMBOL(I)	
	$IF(Y(I) \circ GT \circ YMAX) YMAX = Y(I)$	DECL4330	14 CONTINUE	
	I IF(Y(I) .LT. YHIN) YMIN = Y(I)	D=CL4335		
. '	IF(YMAX -EQ. YMIN) RETURN	0=014340		
•	$5C4LEY = 5 \cdot 1/(YYAX - YMIN)$	DECLA345		
	CO 15 ! = TANPER	02014343	TF(MUS(LIAE - 1910) .EQ. 0) GU TE 9	2801481
	15 IX(I) = 15 X(X(I) = X(I) ASCALEVALA	520L4350 C		DE01467
	$\frac{1}{2} = \frac{1}{2} + \frac{1}$	DECL4355 C	THE ARRAY 'CHAR' IS WRITTEN USING THE PROPER FORMAT	5=51-4-5
·		DEC1.4350 C		
•- *	2 = 17(1) = 171X((7(1) - 7HIN) *SCALEY) + 1	DECL4365	WRITE(6,101) CHAR	
		DECL4370	GD TD 3	
	TX01 AND FIYOF ARE THE LOCATIONS OF THE COORDINATE AXES	DECL4375	8 WRITE(6,102) CHAR	
1		DEC14380		0=
	$IXO = (AXO(IFIX(-X(1) \neq SCALEX + 1 + 1 + 1 + 1)))$	DEC14395		DEC1465
· • • •	IYO = IAXO(IFIX(-YMIN*SCALEY + 1.))	DECLARDO		DECL445
ا النش	IF(IXO , GT, 111) IXO = 1	02014375	IF(CHAR(IXO) .NE. VERT) GD TD 13	DECLAss
	IE(IYC - GT - 51) IYO = 3	DEC14393	CHAR(IXO) = PLUS	DE01465
		DECL4400	13 WRILE(G,103) YVALUE, CHAR	DECLAST
1		0ECL4405	GD TO L	DECLART
1	SCALEX = (X(NFER) - X(1))/5.	DECL4410	10 ANUM(1) = X(1)	DECLIAS
- E 4	SCALEY = (YMAX - YMIN) /S.	DECL4415	$D_0 11 I = 2.6$	0201425
1	YVALUE = YMAX + SCALEY	DECL4420	11 ANUM(I) = $ANUM(I - 1) + SCALEY$	
	LIXE = 52	DECI.4425	WRITE(6.104) ANIM	5221464
'	3 LINE = LINE - 1	06014430		DE01469
1.1	IF(LINE .EG. IVO) ON TO 16	05014425		DECLATE
	IF(LINE - E0 - IVO - 1) GO TO 10		100 TURNATCIALTIAC, DSX, TGRAPH OF THE SOLUTION ///)	DECLARC
1	IF(LINE LE. O) RETURN		101 FURMAT(1H ,15X,11141)	0501471
		DECL4445	102 FORMAT(1H , 10X, 1HY, 4X, 111A1)	0201471
1	THE CHARTER ADDAY LOUDD AN AVETAL FROM	DECL4450	103 FORMAT(1H ,5X,69,2,1X,111A1)	
	The CHARACTER ARRAY CHART IS INITIALIZED	DECL4455	104 FORMAT(1H ,10X,69,2,2(13X,69,2),6X,1HX,6X,2(69,2,13X),69,2)	8261 475
i		DECL4460	END	01114 2
1	99 4 I = 1,111	DECL4465		0501473
	$4 CH_{A}R(I) = BLANK$	DECL4470		
i	CHAR(IXO) = VERT	DECLAA75		
1	50 TC 6			<u></u>
i	16 01 5 1 = 1,111	DECL740V		(11
	5 CHAR(1) - UCD	UE4489		<u>.</u>
;	J CHARLES - PDR	DECL4490		ω.

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Figure 6.1.6

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	<pre>SUBROUTINE PICT(X,Y,Z,N,NUHBER) THIS SUBROUTINE PRODUCES A PERSPECTIVE GRAPH OF THE 'N' PDINTS '(XYY,Z)' GN A STANDARD LINE PPINTER' THE INTEGER 'NUMBER' INDICATES WHICH OF SEVERAL SOLUTIONS STORED IN THE ARRAY 'Z' IS TO BE GRAP HED. DIMENSION X(1),Y(1),Z(1) OIMENSION X(1),Y(1),Z(1) OIMENSION CHAR(111),SYMBDL(10) COMMON CHAR,XNER,YNEWJIX,IY,IZ DATA BLANK,VERT,HOR,PLUS,SLASH,WHY / IH ,IH ,IH-,IH+,IH/,IHY / DATA SYMBOL / 1HC,IH,JHZ,IH3,IH4,1H5,1H6,1H7,1H8,1H9 / NZ1 = N * (NUMBER - 1) + 1 NZ1 = N * (</pre>	DE2D2625 DE2D2635 DE2D2645 DE2D2645 DE2D2645 DE2D2645 DE2D2650 DE2D2650 DE2D2650 DE2D2670 DE2D2670 DE2D2675 DE2D2675 DE2D2710 DE2D2710 DE2D2710 DE2D2720 DE2D2720 DE2D2720 DE2D2720 DE2D2725 DE2D2755 DE2D2750 DE2D2755 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750 DE2D2750	CC	7 14 8 9 13 10 11 990	APPROPRIATE LOCATION DO 14 J = 1/N IF(12(J) .ME. LINE) GO TO 14 CHAR(IX(J)) = SYMBOL(IY(J)) CONTINUE IF(LINE .EO. 25) GO TO 8 IF(MOD(LINE - 1/10) .EQ. 0) GO TO 9 THE ARRAY 'CHAR' IS WRITTEN USING THE PROPER FORMAT WRITE(6/101) CHAR GO TO 3 WRITE(6/102) CHAR GO TO 3 WRITE(6/102) CHAR GO TO 3 VALUE = ZVALUE = ZSCALE IF(CHAR(1) .ME. VERT) GO TO 13 CHAR(1) = PLUS WRITE(6/103) ZVALUE,CHAR IF(LINE .EQ. 1) GG TO 10 GO TO 3 CHAR(1) = CHAR(1 = 1) + XSTEP WRITE(6/104) (CHAR(I),I = 1/6) KETURN N = 0 RETURN FORMAT(1H1/1H0,55%,IGRAPH OF THE SOLUTION'///)
3	LINE = 52 LINE = LINE = 1 IF(LINE .E3, 1) GO TO 16 THE CHARACTER ARPAY 'CHAR' IS INITIALIZED DO 4 I = 1,111 CHAR(I) = RLANK CHAR(I) = VEFT GO TO 6 OD 5 I = 1,111 CHAR(I) = HCP DD 12 I = 1,111,22 CHAR(I) = FLUS IF(LINE .E0, 27) CHAR(27) = WHY IF(LINE .E0, 27) CHAR(27) CHAR(27) CHAR(27) CHAR(27) CHAR(27) CHAR(27) CHAR	DE2D2805 DE2D2810 DE2D2810 DE2D2810 DE2D2810 DE2D2810 DE2D2825 DE2D2825 DE2D2835 DE2D2835 DE2D2840 DE2D2855 DE2	0000	1	SUBROUTINE SFALE(X,N1,N2,XMAX,XM1N,SIZE,XSCALE) THIS SUBROUTINE DETEPMINES THE MINIMUM AND MAXIMUM VALUES AND A SCALE FACTOR FOR THE ARRAY 'X'. DIMENSION X(1) XMAX = X(N1) XMAX = X(1) XMAX = X(1) XSCALE = 0. DU 1 I = N1,N2 IF(X(1) .LT. XMIN) XMAX = X(1) IF(X(1) .LT. XMIN) XMIN = X(I) IF(XMAX .E0. XMIN) RETURN XSCALE = SIZF / (XMAX = XMIN) RETURN END

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6.2 ORDINARY DIFFERENTIAL EQUATIONS

A computer program for the automatic solution of arbitrary linear ordinary differential equations is now virtually complete. The only remaining step is to provide a main program to link the various subprograms together in the proper order. Such a main program is given in Figure 6.2.1. The complete program package, including this main program and twenty-one independent subprograms, requires 190 K bytes of core storage. A large part of these memory requirements is for the three arrays D, C and COM in which all of the major matrix operations are performed.

The program has been extensively tested with both standard differential equations and eigenvalue boundary value problems and some of these results will be presented here. For the first example, consider the differential equation (5.2.3), the solution of which is the exponential function. A complete DECL program listing for this problem is as follows

> * 8 DIFFERENTIAL EQUATION WITH EXPONENTIAL SOLUTION DE: D/DX(Y) = Y (6.2.1) BC: Y(0.0) = 1.0 (6.2.2) X = 0.0, 10.0 (4 ELEMENTS) (6.2.3) WRITE PARTICULAR SOLUTION GRAPH PARTICULAR SOLUTION

This program specifies that equation (5.2.3) is to be solved using four equal eighth order elements in the interval [0.0, 10.0].

The above program was run on an IBM 360/75 computer (as were all examples in this thesis) and required 1.8 seconds

ուս անուստեղում անենան մասեն՝ եներ մոտուտ։ Սոտում ելուու է անեն անձում տեսու են ու են ու են ու ու ու ու ու ու ո Հայաստանությունները մասեն՝ եներ մոտուտ։ Սոտում ելուու է անեն անձում տեսուները են ու են անեն են հետ է ելու ու ու

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Figure 6.2.1

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<pre>THIS PROGRAM WAS WRITTEN BY JULTAN CSENDES DF MCGILL UNIVERSITY MONTREAL , QUEBEC, TO SOLVE ARBITRARY LINEAR ORDINARY DIFFERENTIAL EQUATIONS. THE DIFFERENTIAL EQUATIONS ARE READ IN THE DECL COMPUTER LANGUAGE AND THEIR GENERAL SOLUTIONS ARE METURNED IN TERMS OF BOTH POINT VALUES AND A GRAPHICAL PLOT. PARTICULAR SOLUTIONS ARE ALSO PRODUCED, IF BOUNDARY CONDITIONS ARE SPECIFIED. UNENSION INST(5001;CONSI(500),KD(500),WT(100);DN(6000),C16000); eQUIVALENCE (YR(1),C(4001)); (YI(1),DN(1001)) COMMON COM LCGICAL LANGARITEP;RITEG,PLOTP,PLOPG RELD THE DIFFERENTIAL EQUATION; BOUNDARY CONDITIONS AND GEIXETRIC INFORMATION 1 CALL INSTR(INST;CONST,NUMDE;XBC;X0,WT,NELMTS,NRANGE;NDOM; * LANGARITEP;RITEG;PLOTP;PLOTG,NEIGWR,NEIGPL) NR = NARMEE + 1 ND = NCOM + 1 NULLD;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 NULLD;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 NULLD;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 NULLD;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 NULLD;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 NULCS;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 NULCS;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 NULCS;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 NULCS;NULV;LAMDA) IF(ND - CO, OD TO 1 ND = NCOM + 1 ND = NCOM + 1 ND = NCOM + 1 NULCS;NULV;LAMDA) CALL CONTINUENTIAL SOLUTIONS CONTINUOUS AND PRINT THE GENERAL SOLUTION CALL CONTINUENT AND NON;NULLV;C,NDPTS;NONHOM;NULLD;NULLV; * NCOM + NCOM + NON;NULLV;C,NDPTS;NONHOM;NULLD;NULV; * NCOM + NCOM + NCOM + NCOM + NULLD;NULV; * NCOM + NCOM + NCOM + NULLD;NULV;C,ND + NCOM + NULLD;NULV; * NCOM + NCOM + NCOM + NULLD;NULV; * NCOM + NCOM + NCOM + NULLD;NULV;C,ND + NCOM + NULLD;NULV; * NCOM + NCOM + NCULV;C,ND + NCOM + NULLD;NULV; * NCOM + NCOM + NCULV;C,ND + NCOM + NULLD;NULV; * NCOM + NCOM + NCOM + NULLD;NULV;V,DN,NKBC;NR,ND; * NCOM + NCOM + NULD;NULV;V,DN,NKBC;NR,ND; * NCOM + NCOM + NULD;NULV;V,DN,NKBC;NR,ND; * NCOM + NCOM + NULD;NULV;NMADA;NTIEP) IF(ND - EQ. (D ED TO 1 IF(ND - EQ. (D E</pre>	DECL 5 C CALCULATE THE EIGENVALUES AND EIGENVECTORS DECL 10 C CALL EIGEN(Nn,Nn,NELMTS,NDPTS,NWBC,C,YR,YI,RODTR,RODTI,NEIGWR) DECL 20 IF(ND.EQ.0) GD TD 1 DECL 20 N = NDYT * NEIGPL DECL 30 TD 4, PCCT DECL 40 C GRAPH THE PARTICULAR SOLUTION DECL 50 C 4 CALL GKAPH(XD,DN,N,NDPT) DECL 60 TO 1 DECL DECL 60 TO 1 DECL DECL 50 C FND DECL 50 C GO TO 1 DECL 50 C GO TO 1 DECL 50 C FND DECL 50 C FND DECL 60 TO 1 DECL DECL 75 DECL DECL 75 DECL DECL 75 DECL DECL 10 DECL DECL 10 DECL DECL 10 DECL DECL 10 DECL <t< td=""><td>DECL 225505050505050505050505050505050505050</td></t<>	DECL 225505050505050505050505050505050505050
1F(ND .EQ. C) 60 TO 1	DECL 250	
NORT + ND + NELMIS	DECL 255	
N = NOPT # NVBC	DECL 260	
IF(+NGT, LAMDA) 60 TO 3	DECL 265	ഗ
	DECL 270	6
	DECL ATO	<u>.</u>

of C.P.U. time for solution. The value of the solution at x=10.0 was 22025.7, which is correct to five significant figures. This value is considered to be round-off error limited since the IBM 360/75 has a 24-bit mantissa in single precision.

Equation (6.2.2) has also been solved with the differential equation solving program using first, second, fourth and eighth order polynomial approximations and one, two and four element subdivisions in the interval [0.0, 10.0]. In each case, the difference between the answer and the exact solution was calculated and the L_2 norm of this difference function was evaluated. A similar analysis was performed for the initial value problem.

DE:	D2/DX2(Y) = -Y	(6.2.4)
BC:	Y(0.) = 0.0	(6.2.5)
BC:	D/DX(Y(0.)) = 1.0	

using second, fourth and eighth order polynomials in the interval [0.0, 10.0]. Figure 6.2.2 contains a plot of the L_2 norm of the computed error functions against the number of interpolation nodes for these problems. The lower limiting value of approximately 10^{-4} is interpreted to be the result of round-off error accumulation. The convergent behavior of these error norms can be approximated fairly well with the following empirical formula

 $|| y - y appox || = \frac{40.}{(n-p+1)^3} \left(\frac{8}{N}\right)^n$ (6.2.6) where n is the polynomial order, p is the order of the of the differential operator and N is the total number of points. This behavior indicates that higher order polynomials result not only in more accurate solutions for a given number



Figure 6.2.2 L_2 norm of the error in the solution. (6.2.1) ----; (6.2.4) ----.

of points but also in a faster convergence rate.

As a practical example of the utility of the differential equation solving program in finding general solutions for engineering problems, consider the equation

DE: D3/DX3(Y) + (A+B*SIN(X))*D2/DX2(Y)

+(C+D*SIN(X)+B*COS(X))*D/DX(Y)+E*Y = 0 (6.2.7) which governs the behavior of hydraulic copying mechanisms used in metal cutting [65]. If the homogeneous solutions of this equation decay, the mechanism is stable, if they grow, it is unstable.

Solutions of this problem have been obtained by choosing three linearly independent initial conditions, integrating (6.2.7) numerically for each case, and analyzing the results to determine their growth or decay [65]. When solved with the differential equation solving program of this thesis, none of these complicated procedures is necessary; the equation is simply read in as it appears in (6.2.7), and no boundary conditions are specified. The program automatically returns a full set of homogeneous solutions and, provided the region used encompasses several periods, it is only necessary to observe the shape of the solution in order to determine its stability.

To be specific, let the values of the constants in (6.2.7) be

CST:A = 0.55 CST:B = 0.08 CST:C = 0.825 CST:D = 0.004 CST:E = 0.34

The general solution of the equation in the interval X=0. to 30. is shown in Figure 6.2.3. It is obviously stable. If, however, the value of E is

CST:E = 0.76

the unstable behavior given in Figure 6.2.4 results. These conclusions are confirmed by reference [65]. The graphs in Figures 6.2.3 and 6.2.4 were taken directly from the program output and each analysis, using five ninth order elements, required only 3.0 seconds on an IBM 360/75.

The next examples will indicate the application of the automatic differential equation solving program to eigenvalue boundary value problems. Consider, first, the Mathieu equation

DE: D2/DX2(Y) -2. *Q*COS(2.*X)*Y = A*Y (6.2.8) where Q is a parameter. The solution interval of interest is [0,2 π] and, in order to obtain a complete set of particular solutions, both the Dirichlet and the Neumann boundary conditions need to be specified at the endpoints. Table 6.2.1 contains the first three eigenvalues calculated for these problems by the program in single precision on an IBM 360/75 along with the exact eigenvalues for several values of Q. In these calculations, two ninth-order elements were used and each computation required 3.3 seconds of execution time Note that all eigenvalues produced by the program are accurate to at least four significant figures. The corresponding eigenvectors displayed similar accuracy. Two of the graphs of the approximate eigenvectors which were produced automatically by the subrouting GRAPH for the Mathieu equation





Figure 6.2.4

TABLE 6.2.1

EIGENVALUES OF THE MATHIEU EQUATION N = NEUMANN, D = DIRICHLET

q	Boundary Conditions		Number 1		Number	r 2	Number 3		
<u> </u>	x = 0	x = L	Calculated	Exact (66)	Calculated	Exact (66)	Calculated	Exact (6 <i>6</i>)	
	 N	 N	4545	4551	4.37137	4.37130	16.0332	16.0338	
1	 П	N	1093	1102	9.0503	9.0477	25.0178	25.0208	
	N	D	1.8577	1.8591	9.0770	9.0784	25.0205	25.0208	
	D	D	3.9181	3.9170	16.0306	16.0330	36.0120	36.0143	
10	 N	N	7,7163	7.7174	-13.9359	-13.9369	21.1046	21.1046	
10	ח	N	7,9855	7.8861	-13.9351	-13.9365	26.7693	26.7664	
	N	D	-2.3994	-2.3991	15.5006	15.5028	27.6989	27.7038	
	D	D	-2.3824	-2.3822	17.3761	17.3814	37.4167	37.4199	
20	N	 N	1.1577	1.1543	27.5919	27.5946	-31.3087	-31.3134	
20		N	1,1638	1.1607	28.4692	28.4682	-31.3084	-31.3134	
	N		-14,487	-14.491	15.3736	15.3958	36.6396	36.6450	
	D	D	-14.488	-14.491	15.4662	15.4940	40.5943	40.5897	

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are presented in Figures 6.2.5 and 6.2.6.

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The next example presented is the analysis of the longitudinal vibration of a non-prismatic rod. This example is used because it is a semi-classical problem familiar to most engineers and yet presents difficult differential equations to solve in all but a few special cases. In addition, it is also a problem of considerable practical importance in ultrasonic processing and has a number of known solutions which may be used to assess the accuracy of the results [67].

The four shapes solved were the conical rod, the exponentially tapered rod, the catenoidal rod and a nonuniformly tapered rod having no simple analytical description. For the conical vibrating rod, the differential equation to be solved is

DE: D2/DX2(Y) + 2.0*POW-1(X-0.4)*D/DX(Y)=K*Y (6.2.9) and the boundary conditions are

BC: Y(0.0) = 0.

BC: D/DX(Y(1.0)) = 0.

In this case, the solution interval [0.0, 1.0] was divided into two equal ninth order elements.

The problem required 3.8 seconds of execution time to solve for the eigenvalues and eigenvectors of the system, including a graphical plot of the eigenvectors. The eight lowest eigenvalues are presented in Table 6.2.2 along with their analytically computed values. Notice that the first seven eigenvalues are accurate to four significant figures and have an apparently random error distribution. These errors can be attributed directly to round-off error accumulation in the

(6.2.10)

Figure 6.2.5 The first four eigenvectors of the Mathieu equation. Y(0) = 0, D/DX(Y(L))=0.





Figure 6.2.6 The first four eigenvectors of the Mathieu equation. D/DX(Y(0)) = 0, Y(L) = 0.

TABLE 6.2.2

Longitudinal Vibration Eigenvalues of Non-uniform Rods

	Cone Shaped		Expor	pential	Cateno	Catenoidal		
Eigenvalue Number	Calc.	Exact.	Calc.	Exact	Calc.	Exact		
1	0.80996	0.80976	.77317	.77310	6.1751	6.1730		
2	20.7603	20.7654	21.2535	21.2583	25.9185	25.9122		
3	60.2590	60.2520	60.750	60.745	65.387	65.390		
4	119,473	119.472	119.967	119.964	124.611	124.608		
5	198.39	198.43	198.921	198.922	203.557	203.565		
6	297.04	297.13	297.76	297.62	302.11	302.26		
7	415.62	415.56	415.99	416.05	420.67	420.70		
8	554.7	553.7	555.0	554.2	560.7	558.9		

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with the share of the second states of the second states of the second second

computation. The eighth eigenvalue is, however, slightly worse, as were the successively higher ones, and in these the discretization error is dominant.

The eigenvectors exhibited a similar behavior. The accuracy of the first to seventh eigenvectors was limited by round-off error to four significant figures while eigenvectors higher than the seventh were less accurate. It would appear that above the seventh eigenvector two ninth-order polynomials cannot approximate the **trug** solution with this accuracy. Consequently, if for some reason more accuracy is desired with these high-order eigenvalues and eigenvectors, the problem should be run with the solution region divided into more than two sections.

The exponentially tapered rod and the catenoidal rod were solved using similar data sets, except that their differential equation statements were [68]

DE: D2/DX2(Y) -2.5055 *D/DX(Y) = K*Y (6.2.11) for the exponentially tapered rod and

DE: D2/DX2(Y) -2.*B*(EXP(B*(1-X))-EXP(-B*(1-X)))* POW-1(EXP(B*(1-X)) + EXP(-B*(1-X)))*D/DX(Y)= K * Y (6.2.12)

CST: B = 1.925

for the catenoidal rod.

The eight lowest eigenvalues of these problems are also given in Table 6.2.2, along with the exact values. The agreement is similar to the conical rod.

The non-uniformly tapered rod shown in Figure 6.2.7 was also solved. In this case the differential equation was specified by

DE: D2/ DX2(Y)+FN1 (X) *D/DX (Y) = -K*Y (6.2.13) The function FN1 (X) was supplied by providing the numerical values of the curve in Figure 6.2.7 at the nodal points of two ninth-order interpolation polynomials.

Part of the output from the computer program is shown in Figures 6.2.8 and 6.2.9. The eigenvalues and eigenvectors are given by two sets of numbers, the first of which gives the real part of the solution and the second the imaginary part. In this case, the problem is self-adjoint and the imaginary





Profile of the non-uniformly tapered rod analyzed in Figures 6.2.8 and 6.2.9.

Figure 6.2.8

Part of the computer program output for the vibrating rod of Figure 6.2.7.

		COKE BOTT	PED VIBRATING RO	****						
			THE	DIFFERE	NTIAL EQUATION I	s				
			D2/DX2(Y) + FN1	{X} + D/DX(Y) = -	- K * Y				
			THE	BOUNDAR	Y CONDITIONS ARE					
				Y(1) =	0					
				D/DX(Y	(3)) = 0					
			т	HE GENE	PAL SOLUTION					
				THE	EIGENVALUES					
1 6	0.332795E 01 0.300963E 03	2 7	0,266087E 02 0,415225E 03	3 8	0,612339E 02 0,552151E 03	4 9	0,122639E 03 0,738636E 03	5 10	0,201319E 03 0,934256E 03	
1 6	0.0	2 7	0.0	3 8	0 • 0 0 • 0	4 9	0.0	5 10	0.0 0.0	
				EIGEN	ECTOR NUMBER 1					
1 6 11 16	0.122405E=06 0.571843E 00 0.113229E 01 G.152740E 01	2 7 12 17	0.111557E 00 0.721721E 00 0.124178E 01 0.157367E 01	3 8 13 18	0.224610E 00 0.865722E 00 0.133161E 01 0.160931E 01	4 9 14 19	0.336788E 00 0.100463E 01 0.140701E 01 0.163220E 01	5 10 15 20	0.453565E 00 0.113229E 01 0.147181E 01 0.164022E 01	
1 6 11 16	0.0 0.0 0.0 0.0	2 7 12 17	0.0 0.0 0.0 0.0	3 8 13 18	0.0 0.0 0.0 0.0	4 9 14 19	0.0 0.0 0.0 0.0	5 10 15 20	0.0 0.0 0.0 0.0	•
				EIGEN	ECTOR NUMBER 2					
1 6 11 16	0.656569E=06 -0.143854E 01 -0.956856E 00 C.931725E 00	2 7 12 17	-0,381497E 00 -0,149311E 01 -0,612743E 00 0,129451E 01	3 8 13 18	-0.741593E 00 -0.142665E 01 -0.243629E 00 0.159634E 01	4 9 14 19	-0.104507E 01 -0.124162E 01 0.143087E 00 0.180062E 01	5 10 15 20	+0.128131E 01 -0.956854E 00 0.539966E 00 0.187432E 01	
1 6 11 16	0.0 0.0. C.0 C.0	2 7 12 17	0.0 0.0 0.0 0.0	3 8 13 18	0.0 0.0 0.0 0.0	4 9 14 19	0,0 0,0 0,0 0,0	5 10 15 20	0.0 0.0 0.0 0.0	

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parts of the eigenvalues and eigenvectors are zero. (Nonself-adjoint problems may yield complex eigensystem solutions). Figure 6.2.9 contains the graph of the eigenvectors which was automatically produced by the program.

The differential equation which appears in the analysis of the transverse vibration of a beam is generally regarded to be a more difficult equation to solve than the longitudinal equation because of the higher order derivatives involved. Coded in DECL, the equation which governs the transverse vibrations of a non-prismatic beam is [68]

DE: D2/ DX2(FN1 (X)*(D2/ DX2(Y))) = K *FN2 (X)*Y (6.2.14) where FN1 (X) is the product of Young's modulus and the moment of inertia of the beam and FN2(X) is the product of the longitudinal density of the beam and its cross-sectional area. The boundary conditions of interest are those of simply supported beams

BC: Y(0) = 0. BC: Y(1) = 0. BC: D2/DX2(Y(0.)) = 0. BC: D2/DX2(Y(1.)) = 0. and beams with clamped ends BC: Y(0.) = 0. BC: Y(1.) = 0. BC: D/DX(Y(0.)) = 0. BC: D/DX(Y(1.)) = 0.

In order to demonstrate the accuracy of the program with this problem, the first case solved was that of a uniform







Figure 6.2.10

Profile of the haunched beam corresponding to Table 6.2.3. h = L/10.0.

beam. Again using two ninth-order elements, the problems required 3.5 seconds of execution time on the IBM 360/75 and the first four eigenvalues are presented in Table 6.2.3. This time it was found that only the first six eigenvalues were accurate to four significant figures.

The program was also used to analyze the vibration of the haunched beam [69] shown in Figure 6.2.10. The resulting eigenvalues with both sets of boundary conditions are given in Table 6.2.3.

Longitudinally accelerated transversely vibrating uniform beams with clamped ends were also analyzed with the program. In this case, the differential equation card read by the program was

DE: D4/DX4(Y)+A*D/DX((B-X)*D/DX(Y)) = K*Y (6.2.15) The eigenvalues generated by the program are presented in Table 6.2.4 along with variationally derived approximate

TABLE 6.2.3

Natural Frequencies of Vibrating Beams

	Prismatic Beam				Haunched Beam	
Mode Number	Simply Supported		Clamped Ends		Simply Supported	Clamped Ends
	Calculated	Exact (nπ)	Calculated	Exact		
1	3.1434	3.1415	4.7295	4.7300	.50089	.75222
2	* 6.2838	6.2832	7.8539	7.8532	.91343	1.17049
3	9.4241	9.4247	10.9949	10.9956	1.38097	1.63003
4	12.5673	12.5664	14.1387	14.1372	1.82442	2.07074
5	15.719	15.708	17.295	17.279	2.29374	2.51808
6	19.01	18.85	20.73	20.42	2.7978	3.0434
TABLE 6.2.4

Constant Values	Eigenvalue Number	Program Results	Rayleigh-Ritz Solution [6]	Solution by Kato's Method [67
A = 25.0	1	261.14	264.34	241.50
B = 0.25	2	2934.3	2932.2	2881.4
	3	12,757.8	12,756.1	12,669.1
	4	36,734.1	36,720.5	36,635.3
A = 25.0	1	102.64	104.34	51.29
B = 0.75	2	2354.61	2349.54	2237.66
	3	11,534.2	11.519.0	11.337.2
	4	34,586.8	34,574.8	34,344.0
A = 100.0	1	464.24	464.42	349.23
B = -0.50	2	3739.95	3738.80	3492.14
	3	14,553.6	14,536.3	14,210.1
	4	39,854.8	39,855.6	39,429.1

Eigenvalues of Accelerated Clamped Beams

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lower and upper bounds [70] for some value of the constants A and B.

6.3 NONLINEAR ORDINARY DIFFERENTIAL EQUATIONS

It was shown in section 5.6 that the solution of many nonlinear differential equations can be accomplished via Newton's method by solving a sequence of linear differential equations. Since a computer program has been developed in this thesis which is able to solve any linear ordinary differential equation, the extension of this program to solving nonlinear ordinary differential equations is straightforward. By converting the program to iterative usage, it is possible to determine the Newton sequence of solutions which converge, if they do, to a solution of the original nonlinear equation.

In order to make the linear program suitable for solving a sequence of linear differential equations, the existing programs need to be modified slightly, as has already been described in section 5.6. Figure 6.3.1 contains a main program which incorporates these modifications. This program has been designed to call the subroutines DIFFEQ, CONTIN and BOUND repeatedly, until the successive particular solutions generated by BOUND have converged satisfactorily.

In defining the instructions which determine the differential equation to be solved, the new program calls the subroutine INSTR, which was used in the original non-iterative program. The only additional data that need to be specified

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Figure 6.3.la

						DENL 275
		DENL	5		N = ND##2	DENL 280
	THIS PROGRAM WAS WRITTEN BY ZOLTAN CSENDES OF MCGILL UNIVERSITY,	DENL	10		CUALTY # 0 D[] 0] # 71V	DENL 285
	HONTREAL QUEBEC, TO SOLVE ARBITRARY NONLINEAR ORDINARY	DENL	20		$A C(1) \Rightarrow 0.$	DENL 290
	DIFFERENTIAL EQUATIONS BY NEWTON ITERATION. THE QUASI-LINEAR	DENL	25		J = ND + 1	DENL 293
	FORM OF THE DIFFERENTIAL EQUATION MUST BE READ IN THE DECK	DENL	30		$DD 9 I = 1_{P}N_{P}J$	DENL 305
	COMPUTER LANGUAGE, FULLWRED BY A DATE THE SOLUTION IS TO CONVERGE	DENL	35		COM(I) = 1.	DE'1 310
	TO COMPANY IN STGENVALUE PROBLEMSIS AND THE INITIAL VALUE OF THE	DENL	40		9 C(1) = 1	DENL 315
-	SOLUTION.	DENL	45		IE · NEADU(2)	DENL 320
	2010/10/1	DENL	50			DE4L 325
	DECL 50-75		60			DENL 330
Ē		DENL	65	С		DENL 333
	DIMENSION ENDRY(100)	DENL	70	Č	FIND THE DERIVATIVE OF THE SOLUTION	DENL 345
	COMMON /ANS/ AA,NEADD(10), E(3000),EDER14(3000)	DENL	75	C		DENL 350
-	LUGICAL LASI	DENL	80		CALL DPRATR(C/1/K/L/ND/D/	DEPL 355
	DECL 80-120	DENL	85		1F(L .E4) KKJ GU 10 40	DE"L 360
č	• • • •	DENL	90		$f_{\text{A}} = 1$	DE'IL 365
-	REAU(5,100) LAS, IGENUM, AA	DENL	100		IF(L .NE. ND) GU TO 15	DEAL 370
	WRITE(6,101) AA	DENL	105		1 = NEADD(1) + 1	DENL 3/2
	ITER = 0	DENL	110		J = NEADD(1) + NDPTS	DEVIL 385
	LAST = "FALSE"	DENL	115		$DD = 16 H = I_F J$	DENL 390
	NUMBES & NUMBE NOBTE - NOWNEINTS	DENL	120		16 EDERIV(N) = F(N)	DEML 395
	NEADALLY = 4#NDPTS	DENL	122		SU IU I/ NE CALL ODPATR/CONSTERSKALANDA13	DE"L 400
	N = ND + 1	DENI	125		17 CONTINUE	DE112 452
	00 5 I = 1,NULLD	DENL	140		JED = -ND + NFADD(1)	DE11 410
	N = N + 1	DENL	145		00 10 H = 1 HELHTS	DE1L 419
	5 NEADD($I + 1$) = NEADD(I) + NE + NE(MTS	DENL	150		JED = JED + ND	DENL 425
	J = NEADDINGLED + 17 + NO + NELTON	DEHL	155			DETIL 430
ř	INITIALIZE THE SCLUTION	DENL	160		1C = 1C + 1	DE"L 435
č		DENL	170		EQERIV(IE) = 0.	DE'IL 440
•	DD + S = 1, J	DENL	175		1J = I + K	DENL 447
	5 E(I) = AA	DENL	180		00 10 J ≖ 1≠ND	DE11 455
	9 ITEK = ITER + 1	DENL	185		IJ = IJ + K	DENL 460
	$M_{0} = M_{0} M_{1} = K_{0} + K_{0} T S$	DENI	190		EDERIV(IE) = EUERIV(IE) + CUR(II) + CUR(II) + CUR(IV)	DENL 465
	NUMDE = NUMBES	DENI.	192		1 = NEADD(1) = K + 1)	DE'L 470
c		DENL	200		J # NELMTS#K	DENL 475
č	DECL 125-290	DENL	210		IF(K .GT. 1:R) GD TO 12	0276 455
C		DENL	215	c		DENL 490
		DENL	220	C	DETERMINE IF THE ITERATIONS ARE COMPLETED	DENL 495
	IF(RU .EW. U) GU (U I I	DENL	7.25	С	TT (UD CO / TV (AST) 60 TO 14	DENL 500
		DENL	230		1P(DU, CA, W, CAST, CAST, CO, D, TA, TRUE, TELEDING (TTED) (T. 1. F=7) LAST = TRUE.	DE'4L 505
	1F(J .LT. 0) J = 4 # NDPTS	DENL	235		IF(ITER .IT. LAS) GO TU 9	DE"1 510
c		DENI	240		LAST = .TRUE.	DENL 512
č	EVALUATE THE DIFFERENCE NORM	DENL	250		GO TO 9	DEML 520
C		DENI	. 255		14 N = NUPTS + (14PC + (14D + 1))	DE11 530
i	$DO[7] = 1 \mu (0075)$ $L = 0 P (1 + J)$	DENL	. 260		CALL RITE(E,NDPTS,MD + 1,0,1)	
	= (1 + 1) = (1) + (1 + (DENI	. 265			7 1
	ENDRM(ITER) = ENERM(ITER)/NDPTS	OEUT	. 270			7 '7

£__;



Figure 6.3.1b

to the program are the maximum permissible number of iterations in the solution process and the initial value of the function z of equation (5.6.6) on the interpolation point set. These function values are stored in the array Z, which shares common storage locations with the function subprograms FCN and FNI-FN5. Since the function z and its derivatives often need to be approximated by several different orders of polynomials in a single equation, provision has been made in the program to store the numerical values of up to ten different functions.

Using the initial values in the array Z the program calls the subroutines DIFFEQ, CONTIN and BOUND in succession, in exactly the same manner as they were called by the noniterative program. The particular solution obtained in this way is then compared to the function z. If the L_1 norm of the difference of the point values of these two functions

$$L_{1} = \sum_{I=1}^{N} |Y(I) - Z(I)|$$
 (6.3.1)

is less than a pre-defined value, the solution is assumed to have converged and the results printed. If, however, the L_1 norm is too large, the array Z is updated with the new particular solution values and the process repeated. The iterations stop whenever the L_1 norm of the difference function is sufficiently small or when the number of iterations exceeds the specified value. In most cases, it is found that if the procedure is convergent, about five iterations are sufficient to produce a solution limited only by round-off error.

The convergence of the above process in five iterations can be explained theoretically by noting that according to equations (5.6.3) and (5.6.6), convergence of Newton's method is quadratic. This means that the accuracy of the solution in each successive iterate is twice as good as the accuracy in the preceding iterate. Therefore, if the initial function contains one correct binary digit, the fifth iterate will contain $2^5=32$ correct binary digits. Since this is more than the wordlength of the computer used, with most initial functions, five iterations are sufficient.

The operation of the program will now be illustrated by solving several representative nonlinear differential equations. Consider, first, the differential equation (5.6.9) which was used as an example in section 5.6. The sequence of linear differential equations to be solved in this case is, from (5.6.11), in DECL,

DE: D3/DX3(Y) + FN1(X) + Y

+FN2(X)*D2/DX2(Y)=FN1(X)*FN2(X) (6.3.2)

where

FN1(X)=D2/DX2(Z)

FN2(X)=Z

The boundary condition statements for this problem are

BC: Y(0.0)=0.0 BC: D/DX(Y(0.0))=0.0 (6.3.3) BC: D2/DX2(Y(0.0))=1.0





Convergence of Newton's method to the solution of (6.3.2).

- × one element
- o two elements
- Δ three elements

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TABLE 6.3.1

	<u></u>	rojective Meth	Runge-Ku	tta [49]	
	l element	2 elements)(4 iterations	3 elements)(4 iterations)	h=1.0	h=0.5
<u> X </u>	0.0553901	0.0555265	0.0557907	6	6
0.5		0.124789		0.125	0.124739
0 66667	0.220251	0.221253	0.222116	6	6
	0 489594	0.492254	0.494041	0.491425	0.491908
1.0	0.851949	0.856874	0.859315	6	6
1.5		1.06869		1.06863	1.06792
1.66667	1.28903	1.29656	1.29853	6	6
	1 77000	1.78976	1.78989	1.78738	1.78792
<u>2.0</u> 23333	2.30135	2.31563	2.31214	6	6
2.5		2.58561		2.57913	2.58121
2.66667	2.84065	2.85832	2.84920	6	6
3.0	3.38751	3.40822	3.39273	3.37957	3.40061

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Jution of a nonlinear third order differential equation.

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In this case, the initial value of z was taken to be the zero function and the problem was solved in the interval [0,3] using one, two and three ninth-order elements.

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Figure 6.3.2 contains a graph of the logarithm of the L_1 norm of the error in the approximate solution versus the logarithm of the number of iterations for this problem. Notice that the L_1 norm decreases until the sixth iterate in the one element solution and until the fourth iterate in the two and three element solutions. In this region, the convergence is seen to be approximately quadratic. Beyond these iterations, the change in the solution appears to be random and, as a result, can be attributed to round-off error. Table 6.3.1 contains the solutions obtained from the computer program in these cases. The solution values converge to an accuracy of about four significant figures as the number of elements used increases.

As a check on the results, included in Table 6.3.1 are two Runge-Kutta calculations of fifth-order accuracy for the differential equation (5.6.9)[49]. It may be observed that there is good agreement between the projective solutions and those obtained by Runge-Kutta integration.

For the next example, consider the initial value problem

$$\frac{dy}{dx} - (1 + y^2) = 0$$
 (6.3.4)

$$y(0) = 0$$

which has the solution y=tan x. In applying Newton's method to this problem, the sequence of linear initial value problems

DE: D/DX(Y) - 2.0*FN1(X)*Y = 1.0 - POW2(FN1(X)) (6.3.5) BC: Y(0.) = 0.

where

FN1(X) = Z

needs to be solved. Using zero as the initial solution function, the sequence of solutions obtained from this procedure is known to converge to the solution of (6.3.4) in the interval $0 \le x < \Pi$ [60].

Equations (6.3.5) were solved with the iterative computer program of this section in the interval [0.0, 1.5] by using two ninth-order approximating functions. In this case, since a much more rapid variation was expected in the . solution near x=1.5 than at x=0.0, the first element was chosen to span the interval from z=0.0 to x=1.0 and the second from x=1.0 to x=1.5. Figure 6.3.3 presents the graph of the solutions obtained from the program for the first five equations in the sequence (6.3.5). It can be seen that each successive solution rises above the preceding one and provides a better approximation to the exact solution y=tan x. For this problem, the L $_1$ norm of the difference function decreased until the eighth iteration and the solution obtained on this iterate is given in Table 6.3.2 along with the exact values. It is observed from this table that while the approximate solution in the first element from x=0.0 to x=1.0 is accurate to at least four significant figures, the solution in the second element from x=1.0 to x=1.5 becomes progressively less accurate until, at x=1.5, there is a 6% error. This behavior is the result of the singularity of the solution at $x=\pi/2$.



Table 6.3.2

Solution of Equation (6.3.5). Calculated

1	-0.596046E-07	2	0.111551E 00	3	0.225923E 00	4	0.346233E 00	5	0.476200E 00
6	0.620751E 00	7	0.786813E 00	8	0.984837E 00	9	0.123174E 01	10	0.155737E 01
11	0.155737E 01	12	0.173892E 01	13	0.199219E 01	14	0.230028E 01	15	0.270325E 01
16	C.324714E 01	17	0.402759E 01	18	0.526841E 01	19	0.750886E 01	20	0.131813E 02
1	0.0	2	0.111571E 00	3	0.225954E 00	4	0.346254E 00	5	0.476221E 00
6	0.620775E 00	7	0.786843E 00	8	0.984874E 00	9	0.123180E 01	10	0.155741E 01
11	0.155741E 01	12	0.176597E 01	13	0.201997E 01	14	0.233825E 01	15	0.275168E 01
16	0.331450E 01	17	0.413171E 01	18	0.543649E 01	19	0.787214E 01	20	0.141010E 02

A third example of a nonlinear initial value problem that has been solved with the iterative differential equation solving computer program is provided by

$$\frac{dy}{dx} = y - \frac{2x}{y}$$
 (6.3.6)

y(0) = 1

In this case, the quasi-linearized form of the differential equation is

DE:
$$D/DX(Y) - Y - 2.0 \times X \times POW2(FN1(X)) \times Y$$

= -4.0 \times X \times FN1(X)

where

FN1(X) = 1.0/Z

and Z is taken to be equal to 1.0 initially. Table 6.3.3 contains the solution of this problem using two equal ninthorder elements after five iterations. The execution time required for this problem was 5.6 seconds. Comparison with the exact solution $y=(2x+1)^{\frac{1}{2}}$ [49], also given in Table 6.3.3, shows that the projective solution is accurate to at least four significant figures everywhere. The error remaining is

Table 6.3.3

Solution of Equation (6.3.6).

Calculated

1	0.100000E 01	2	0.110552E 01	3	0.120182E 01	1 4	0,129096E 01	5	0.137433E 01
6	0.145292E 01	7	0.152748E 01	8	0.159855E 01	1 9	0,166660E 01	10	0.173197E 01
11	0.1731971 01	12	0.179494E 01	13	0.185578E 01	1 14	0,191468E 01	15	0.197182E 01
16	0.202734E 01	17	0.208136E 01	18	0.213401E 01	1 19	0,218536E 01	20	0.223552E 01
				Exac	t				
1	0,100000E 01	2	0.110554E 01	3	0,120185E 01	1 4	0.129099E 01	5	0.137437E 01
6	0,145297E 01	7	0.152752E 01	8	0,159861E 01	1 9	0.166667E 01	10	0.173205E 01
11	0,173205E 01	12	0.179505E 01	13	0,185592E 01	1 14	0.191485E 01	15	0.197203E 01
16	0,202759E 01	17	0.208167E 01	16	0,213437E 01	1 19	0.218581E 01	20	0.223607E 01

again attributed to round-off error accumulation on the IBM 360/75.

Consider now the problem of solving nonlinear boundary value problems by Newton's method, instead of initial value problems. In form, the procedures used are identical in both cases. As an illustration of the technique, the nonlinear boundary value problem

$$\frac{d^2 y}{dx^2} = \frac{3}{2} y^2$$
(6.3.8)
y(0) = 4
y(1) = 1

will be solved and the results compared with the exact solution [49]

$$y = \frac{4}{(1+x)^2}$$
 (6.3.9)

The quasi-linearized form of (6.3.8) is (6.3.10) DE: D2/DX2(Y) - 3.0* FN1(X)*Y=-1.5*POW2(FN1(X))

where

FN1(X) = Z

The interval [0.0, 1.0] was divided into two equal ninthorder elements. After five iterations, the solution had

Table 6.3.4

Solution of Equation (6.3.8).

Calculated

1 6 11 16	0.400000E 01 0.244981E 01 0.177756E 01 0.126554E 01	2 7 12 17	0.359003E 01 0.224987E 01 0.165288E 01 0.119001E 01	3 8 13 18	0.323999E 01 0.207343E 01 0.154087E 01 0.112106E 01	4 9 14 19	0.293874E 01 0.191697E 01 0.143987E 01 0.105793E 01	5 10 15 20	0.267762E 01 0.177756E 01 0.134849E 01 0.100000E 01
				E۶	(act				
1	0,400000E 01	2	0.359003E 01	3	0,324000E 01	4	0.293878E 01	5	0.267769E 01
6	0.244991E 01	7	0.225000E 01	8	0,207360E 01	9	0,191716E 01	10	0.177778E 01
11	0,177778E 01	12	0.165306E 01	13	0.154102E 01	14	0.144000E 01	Ĩ5	0.134860E 01
16	0.126563E 01	17	0.119008E 01	18	0,112111E 01	19	0,105796E 01	20	0,100000E 01

converged satisfactorily and this answer is presented in Table 6.3.4. The C.P.U. time used in this solution was 4.8 seconds. Once again, the solution values are accurate to four or more significant figures.

Another example of a nonlinear boundary value problem solved with the iterative computer program is

$$\frac{d^{3}y}{dx^{3}} + \frac{6y}{dx^{3}} + \frac{9y}{2} = -\frac{3}{2} \cos x \qquad (6.3.11)$$

$$y(0) = y(2\pi)$$

$$\frac{dy(0)}{dx} = \frac{dy(2\pi)}{dx}$$

Although the exact solution of (6.3.11) is not known, Collatz has expanded the solution in terms of a Fourier series and obtained the following approximate solution [49]

 $y^{\sim} - \frac{3}{320000}$ (805.5 + 32240 cos x + 2418 cos 2x (6.3.12)

 $-240 \cos 3x - 2.7 \cos 4x$)

In solving (6.3.11) using Newton's method, the following sequence of linear boundary value problems need to be solved

Table 6.3.5

Solution of Equation (6.3.12).

Calculated

7 **	1 6 11 16	-0.330338E 00 0.671440E=01 0.269717E 00 -0.400692E=01	2 7 12 17	-0.307883E 00 0.157045E 00 0.257938E 00 -0.149697E 00	3 8 13 18	-0.244180E 00 0.221138E 00 0.221139E 00 -0.221139E 00 -0.244181E 00	4 9 14 19	-0.149698E 00 0.257938E 00 0.157048E 00 -0.307883E 00	5 10 15 20	-0,400719E-01 0,269718E 00 0,671473E-01 -0,330338E 00
					E	xact				
	1 6 11	-0,330195E 00 0,673786E-01 0,269805E 00 -0,398399E-01	2 7 12 17	-0.307809E 00 0.157144E 00 0.257985E 00 -0.149604E 00	3 8 13 18	-0.244174E 00 0.221149E 00 0.221150E 00 -0.244173E 00	4 9 14 19	-0.149605E 00 0.257984E 00 0.157145E 00 -0.307809E 00	5 10 15 20	-0.398413E-01 0.269805E 00 0.673801E-01 -0.330195E 00

DE: D2/DX2(Y) + 6.0*Y + FN1(X)*(Y) (6.3.13)

= POW2(FN1(X)) - 1.5*COS(X)

BC: Y(0.) = Y(6.283184)

BC: D/DX(Y(0.)) = D/DX(Y(6.283184))

where FN1(X) =Z.

Using two equal ninth-order elements and the zero function as Z on the initial iterate, the sequence of solutions of these boundary value problems converge very rapidly with the L_1 norm of the difference between iterates decreasing to 3.6 x 10^{-7} on the fourth iterate. Table 6.3.5 contains the values of the solution from the fourth iterate as well as the values of the function (6.3.12) on the set of interpolation points. It can be seen that the two approximate solutions agree very closely.

6.4 PARTIAL DIFFERENTIAL EQUATIONS

The technique and required subprograms for the solution of linear two-dimensional partial differential equations have already been described. A main computer program which connects these subprograms together in the proper order is presented in Figure 6.4.1. When this main program is used with the overlay procedure in Figure 6.4.2, the entire two-dimensional program package requires 194K bytes of core memory. In this way, the arrays DN, C and COM which contain the large twodimensional coefficient matrices are enlarged without increasing the memory requirements of the program.

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Figure 6.4.1

	DESD	5
NUMTREAL OWNERS AN ITTEN BY ZULTAN CSENDES OF MCGILL UNIVERSITY,	DE2D	ï0
RADITAL DIFERENTIAL SOLVE ARBITRARY LINEAR TWO-DIMENSIONAL	DE2D	15
ANITAL DIFFERENTIAL EQUATIONS.	DE2D	20
NTHENETON HD. FROM HERE AND A RECEIPTION	DEZD	25
DIMENSION YD(500) MTY(100), IELCX(100), IELCY(100), YBC(300)	DE2D	30
DIMENSIUM INST(5(C), CUNST(500), XD(500), WTX(100), DN(8000), C(8000),	DF2D	35
CUM(19000); VR(1); YI(1); RDDTR(100); RDDTI(100); XBC(300)	DF2D	40
	DE20	45
DECL 60-95		50
	DEZD	55
1 CALL INST2D(INST, CONST, NUMDE, XBC, YBC, XD, YD, WTX, WTY, IELCX, IELCY,		60
* NCCOND, NELMTS, NRANGX, NRANGY, NDOM, LAMDA, RITEP, RITEG, PLOTP,		65
* PLOTG, NEIGWE, NEIGPL)		70
NRX = NRANGX + 1		75
NKY = NRAHGY + 1		90
ND = NDDM + 1		85
CALL DIFF2D(DN,MCN,WTX,WTY,NELMTS,INST,CONST,XD,YD,NRX,NRY,ND,		00
* NULLD, NULLV, LAMDA)		05
		100
DECL 150-175		105
	DE2D	110
CALL CONT2D (PN, NFX; NRY; ND; MDN; NULLV; C; NDPTS; NDNHOM; NULLV;		115
* NCCDND, IELCX, IELCY, NELMTS, WTX, WTY, LAMDA, RITEG)		120
		125
DECL 190-210		130
		125
DU 5 I = 1,NULV	DF2D	140
GALL PICI(XU)YD)(JNDPTS)I)	DE2D	145
2 CALL BUUN2D (C, NDFTS, NONHOM, NULLD, NULLV, DN, NWBC, NRX, NRY, ND,	DE2D	150
* INST, CONST, NUMDE, XD, YD, NELMTS, WTX, WTY, LAMDA, RITEP, YI)	DE2D	155
	DESD	160
D ^E CL 250-320	DE2D	165
	DE2D	170
AF(GAMDA) HWRC = NEIGPL	DE2D	175
	DE2D	180
(ALL PICICXD, YD, DH, NDPT, I)	DE2D	185
	DESD	190
END	DE2D	195

Figure 6.4.2

ENTRY MAIN 1

INSERT MAIN, NULL, RITE

UVEPLAY ALPHA

- 2345 INSERT INST2D, CHARTE, ENUM, PICT, SCALE UVERLAY ALPHA
- 6
- INSERT DJFF2D, CONT2D, BOUN2D, OPG2D1, OPG2D2, OPRATR, FCN2D, KRON, DPFU1, ALPHA 7
- INSERT EIGEN, RILMAT, EIGQR, DQRT 8

С С С

C C C

C C C

С С С

In order to illustrate the operation of the program, consider the elliptic partial differential equation

$$\nabla^2 z = -2 \sin x \sin y \qquad (6.4.1)$$

with the boundary conditions

$$z(0,y) = z(\pi,y) = z(x,0) = z(x,\pi) = 0 \quad (6.4.2)$$

Using one seventh-order element, a working DECL computer program to solve this problem is

* 7	POISSON'S EQUATION IN A RECTANGLE
DE:	D2/DX2(Z)+D2/DY2(Z)=-2*SIN(X)*SIN(Y)
BC:	Z(0.0, 0.0)=0.0
BC:	Z(0.0, 0.449) = 0.0
	•

BC: Z(3.14, 3.14)=0.0 REC: X=0.0, 31.41592, Y=0.0, 3.141592

In this program, the boundary conditions (6.4.2) have been translated into twenty-eight separate boundary condition statements specifying the solution value at each of the twenty-eight nodal points on the perimeter of the element. Due to the interpolatory nature of the approximating functions b_{ij} the twenty-eight boundary conditions in the above statements are linearly independent and completely equivalent to (6.4.2) in the approximating space. Although it would be a desirable extension of the DECL language, present program capabilities doe not include generation of the above boundary condition statements from compact statements such as

BC: Z(0.0, Y) = 0.0

The above DECL program was run on the IBM 360/75 and required 18.1 seconds of C.P.U. time for execution. The nontrivial solution values are given in Table 6.4.1 for one half-quadrant of the solution region along with the corresponding values of the exact solution.

$z=\sin x \sin y \qquad (6.4.3)$

Taking advantage of the symmetry in the above problem, equation (6.4.1) was also solved using one seventh-order element in the rectangle $x=[0,\pi/2],y=[0,\pi/2]$ with Neumann boundary conditions along the lines $x=\pi/2$ and $y=\pi/2$ and this solution is also presented in Table 6.4.1. Notice that the solution obtained without using the symmetry properties of the problem is accurate to about three significant figures and that the solution obtained by using symmetry is accurate to about four significant figures.

Table 6.4.1

		Variaes for FUIs		in a rectangle
x	У	Numerica]	Numerical	Exact
		without		
	ι	ising symmetry	using symmetry	
π/7	π/7	0.188149	0.188239	0.188256
π/7	2π/7	0.339105	0.339204	0.339224
π/7	3π/7	0.422871	0.422955	0.423005
2π/7	2π/7	0.611182	0.611237	0.611261
2π/7	3π/7	0.762153	0.762205	0.762230
3π/7	3π/7	0.950418	0.950470	0.950485

Solution point values for Poisson's equation in a rectangle.

Table 6.4.2

Point values of the solution of parabolic partial differential equations.

Point	coor	dinates	(6.4.4)	(6.4.6)	
x	У	Numerical	Exact	Numerical	Exact
	π/7	0.27704	0.27699	0.67997	0.67965
π/7	4π/7	0.072070	0.072066	2.6128	2.6123
π/7	π	0.018739	0.018750	10.002	10.040
2π/7	π/7	0.49917	0.49912	1.2251	1.2247
2π/7	4π/7	0.12989	0.12986	4.7081	4.7071
2π/7	π	0.033774	0.033786	18.049	18.092
3π/7	π/7	0.62235	0.62239	1.5273	1.5272
3π/7	4π/7	0.16203	0.16193	5.8704	5.8697
3π/7	π	0.042142	0.042131	22.546	22.561

Figure 6.4.3 contains the graph produced of the solution in the first case by the subroutine PICT. Smooth lines have been drawn by hand on this figure to round out the surface contours between point values.

The next example presented is the two-dimensional parabolic partial differential equation

DE:D2/DX2(Z) = D/DY(Z) (6.4.4)

with

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$$z(0,y)=z(\pi,y) = 0$$

 $z(x,0)=sin x$

A single seventh-order element was used in the rectangle $x=[0,\pi], y=[0,\pi]$. The DECL computer program for this problem contained twenty-two boundary condition statements and



required 15.3 seconds of C.P.U. time to produce a solution. Selected point values of this solution are given in Table 6.4.2 along with values of the exact solution -u (6.4.5)

 $z=sin \ x \ e^{-y}$ (6.4 This table also contains the values of the solution of

the complementary problem

DE: D2/DX2(Z) = -D/DY(Z) (6.4.6)

which has the exact solution

$$z=sin \ x \ e^{y} \tag{6.4.7}$$

Both numerical solutions agree with the exact solutions to three or four significant figures. The graph produced by the program with the solution of equation (6.4.6) is shown in Figure 6.4.4.

For the third example, consider the eigenvalue boundary value problem

DE:
$$D2/DX2(Z) + D2/DY2(Z) = K*Z$$
 (6.4.8)

$$z(0,y) = z(x,0) = \frac{d}{dx}z(1,y) = \frac{d}{dy}z(x,1) = 0$$
 (6.4.9)

The eigenvalues of this problem are known to be

$$K = \left(\frac{\pi}{2}\right)^{2} \left[\left(2n-1\right)^{2} + \left(2m-1\right)^{2} \right] \quad n, m=1, 2, \dots \quad (6.4.10)$$

and the eigenfunctions are

$$z=sin\left(\frac{n\pi}{2} x\right)sin\left(\frac{m\pi}{2} y\right) \tag{6.4.11}$$

Equation (6.4.8) was solved using one sixth-order element and required 25.5 seconds of C.P.U. time for solution. It is of interest to note that in this problem the solution was comprised of 49 independent point values but that with the procedure of section 5.5, the order of the matrix eigenvalue problem solved was only 25. Table 6.4.3 contains the values of the first eight eigenvalues produced by the program as





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seen that the first three eigenvalues are accurate to four significant figures but that the higher eigenvalues become progressively less accurate with a 1.5% error in the eighth eigenvalue. Computer plots of the first and fourth modes of this problem are shown in Figures 6.4.5 and 6.4.6.

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Figenvalues	of	the	Helmholtz	equation	in	a	rectaŋgle.
LIGGHIMIMUT	-						the second se

Eigenvalue	Numerical	Exact	
number			
]	4.93452	4.93480	
2	24.6784	24.6740	
3	24.6793	24.6740	
4	44.426	44.413	
5	65.49	64.15	
6	65.49	64.15	
7	85.24	83.89	
, 8	85.24	83.89	

All of the examples shown up to this point have been two-dimensional partial differential equations in rectangular regions for which the locations of the boundary conditions coincided with the edges of the elements. Consider now the problem of solving Helmholtz's equation (6.4.8) in a circular region where

$$z(x^2 + y^2 = 1) = 0 \tag{6.4.12}$$



Figure 6.4.6.



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In this problem, the projective solution region will not match the region defined by the boundary conditions for any rectangle or combination of rectangles. Consequently, a method for specifying boundary conditions at locations other than the edges of the elements must be employed.

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As pointed out in section 5.3, equations (5.3.3) and (5.3.4) may be used to determine the numerical conditions on a projective solution for it to satisfy any boundary condition in any geometric location. However, as was the case with rectangular regions, when working with an irregular region it is wise to take advantage of the orthogonality properties of the interpolation polynomials b_{ij} in choosing boundary condition point locations. In the case of the circular boundary condition (6.4.12), for example, a sixth-order element may be super-imposed on the circle as shown by the



One quadrant of a sixth order element super-imposed on a circular boundary.

quadrant in Figure 6.4.7. The straight lines drawn on this figure indicate the locations of the zeros of the interpolation polynomials b_{ij} and the crosses denote the points of intersection of these lines with the circular boundary. In analogy with the rectangular boundary value problem (6.4.9), it is natural to take these points of intersection to be the locations on the circle where the approximate solution is made equal to zero. This procedure has the effect of shifting the four boundary condition locations A,B,D and E used in the rectangular case to the locations A', B', D' and E' on the circular boundary. However, a simple count of the number of arbitrary constants in the general solution reveals that one additional boundary condition needs to be specified besides the six imposed at the crosses in Figure 6.4.7.

The addition of this extra boundary condition is, however, unusual in the sense that it is not used to restrict the behavior of the solution on the circular boundary but rather as a means of redefining the basis functions in the domain space of the differential operator. In the above problem, the complete set of basis functions b_{ij} which interpolate over the entire square region in Figure 6.4.7 does not constitute a satisfactory basis set with the circular boundary conditions (6.4.12) because it contains an extraneous function. The extraneous function in this case is the interpolation polynomial which has a unit value at point C in Figure 6.4.7 and governs the solution behavior in the corners of the element. By defining a new set of basis functions in which this interpolation polynomial does not

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exist, a basis is formed in which the number of arbitrary constants in the general solution equals the number of intersection points of the nodal lines with the circle.

It is important to note in this analysis that the new set of basis functions is derived from the established set by eliminating some of its members. In the present case, this is most easily accomplished by setting the coefficient of the unwanted corner interpolation polynomial equal to zero with the boundary condition statement

BC: Z(1.0, 1.0)=0.0 (6.4.13) In this way, the solution of problem (6.4.12) can be obtained directly from the DECL computer program without performing any program modifications.

There are, of course, some disadvantages to using the elimination procedure described above to solve problem (6.4.12) with rectangular elements instead of defining special circular elements which fit the geometry exactly. Solutions obtained with the above procedure are optimized over the entire element and not in the smaller circular region. Furthermore, the DECL computer program will perform unnecessary arithmetic operations with function coefficients which will be eliminated. However, in most engineering applications, the resulting inefficiencies are not serious enough to justify solving the problem with special techniques.

A DECL computer program was written to solve problem (6.4.12) using the boundary condition statements described above. The program required 27.1 seconds of C.P.U. time for execution and the first eight eigenvalues produced are given in Table 6.4.4. The problem was also solved in one quarter

of the region using a sixth-order element and Neumann boundary conditions along the axes. These results are also presented in Table 6.4.4 along with the exact eigenvalues. A comparison of the answers shows that the lower order eigenvalues are accurate to two or three significant figures. Bearing in mind that the solutions are not stationary in the eigenvalues, this accuracy is quite acceptable. It is somewhat lower than was the case with the rectangular problem (6.4.9), however, since in that case a larger and more suitable basis set was used. Plots of the first and sixth modes for the circular Helmholtz boundary value problem are shown in Figures 6.4.8 and 6.4.9.

Table 6.4.4

teenvolues of the Helmholts equation in a sincle

Eigenvalues of	LITE HEIMHOILZ EL	uacion in a circ	16.	
Eigenvalue	Numerical	Numerical	Exact	
number	full circle	quarter circl	8	<u></u>
1	2.389	2.393	2.405	
2	3.877	-	3.832	•
3	3.877	-	3.832	
4	5.110	5.137	5.135	
5	5.260	-	5.135	
6	5.638	5.524	5.520	
7	5.712	-	5.520	
8	7.705	7.609	7.016	

In order to demonstrate the feasibility of solving two-dimensional partial differential equations in complicated regions with the DECL program, Helmholtz's equation was also











sixth-order element was used in the approximation with the boundary condition point locations chosen as indicated by the circles in Figure 6.4.10. Here Neumann conditions were used for points along the coordinate axes and Dirichlet conditions were imposed elsewhere. The first eigenvalue for this problem was equal to 2.027 with the corresponding eigenvector behaving similarly to the dominant eigenvector in the circular case. These results agree with an analysis of this problem using ordinary finite elements to approximately three significant figures [71].



Cigar-shaped region showing boundary condition point locations.

Figure 6.4.10

Finally, consider the following two-dimensional Bessel-type equation

DE: POW2(X) * (D2/DX2(Z) + D2/DY2(Z)) (6.4.14)

DE: +X*D/DX(Z) -Z = K*POW2(X)*Z

which governs the behavior of the azimuthal electric and magnetic fields in a rotationally symmetric resonant cavity. The boundary conditions of interest are

$$z(boundary)=0.0$$
 (6.4.15)

for electric fields and for magnetic fields

э.,

$$x \frac{dz}{dn}(boundary) = -z(boundary)\frac{dz}{dn}$$
 (6.4.16)

where n is the direction normal to the boundary. This boundary value problem has been solved with the DECL program and some of these results will be presented here. Consider first a rectangular toroidal cavity with x=[2,5] and $y=[-\frac{1}{2},\frac{1}{2}]$.

Table 6.4.5

Resonant frequencies of a rectangular toroidal cavity with $x=[2,5], y = [\frac{1}{2}, \frac{1}{2}].$

Eigenvalu number	e Electr	Electric field		Magnetic field		
	Numerical	Exact*	Numerical	Exact*		
]	3.32117	3.32150	1.03668	1.03660		
2	3.78948	3.78526	2.08585	2.08868		
3	4.475	4.451	3.14177	3.13768		
4	6.397	6.375	3.30820	3.30820		
5	6.655	6.629	3.77096	3.77250		
6	7.068	7.030	4.588	4.440		
* These 1	numbers were com	puted by A. K	lonrad.			

The resonant frequencies of the electric and magnetic fields obtained for this cavity from the DECL program using one sixth-order element are presented in Table 6.4.5 along with the exact values. There is excellent agreement with all low order eigenvalues. In addition, Figures 6.4.11 and 6.4.12 contain two of the field plots produced by the DECL program.

Now consider a right cylindrical cavity in which the solution of equation (6.4.14) is required in the region x=[0,3], $y=[-\frac{1}{2},\frac{1}{2}]$. Unfortunately, equation (6.4.14) has a singularity along the line x=0.0 due to the $1/x^2$ coefficient in Bessel's equation and function values along this line cannot appear in the solution vector. Consequently, in order to obtain meaningful results, the solution region must be made to overlap the line x=0.0. It is not necessary to avoid the line x=0.0 when specifying boundary conditions however, since the boundary conditions (6.4.15) and (6.4.16) are valid in any location.

The calculation of the electric and magnetic fields in the right cylindrical cavity has been performed by the DECL program using a sixth-order element in the region x=[-0.05, 3.0], y = [-0.5, 0.5] and the eigenvalues produced are given in Table 6.4.6. In this calculation, the boundary conditions were applied along the lines x=0.0, x=3.0, y=-0.5and y=0.5. The eigenvalues have approximately the same accuracy as in the previous example. Figures 6.4.13 and 6.4.14 contain examples of the field plots produced by the DECL program for this problem.

Figure 6.4.11



Figure 0.4.12



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Resonant frequencies of a right cylindrical cavity with $x=[0,3], y=[-\frac{1}{2},\frac{1}{2}]$.

Eigenvalue number	e Electric field		Magnetic field		
	Numerical	Exact*	Numerical	Exact*	
1	3.39112	3.39130	0.80129	0.80160	
2	3.92052	3.91649	1.8201	1.8400	
3	4.847	4.623	2.9437	2.8845	
4	6.433	6.412	3.242	3.152	
5	6.728	6.704	3.631	3.641	
6	6.975	7.140	4.305	3.931	
* These n	umbers were comput	ted by A. Kon	by A. Konrad.		

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<u>CHAPTER Z</u>

CONCLUSIONS

In this thesis, two different but nevertheless interrelated themes have been developed. The first and perhaps primary theme has been the introduction of a discretization procedure which uses different orders of polynomials to approximate the domain and the range spaces of differential equations. The notion of assembling complicated operators from simple ones, using projection operators between different spaces, and employing generalized matrix inversion to solve the resulting rectangular matrix equations are all logical consequences of this numerical approach. The second theme has been the development of differential equation solving programs with versatile and easy to use input and output characteristics. This development essentially constitutes the creation of a new special purpose computer language which makes it possible to solve many differential equations automatically.

The advantages of using the discretization procedure developed in this thesis are numerous. First of all, it is computationally efficient. The method relies primarily on matrix addition and multiplication for its execution, a trait well-suited to the operating characteristics of digital computers. Second, the method is highly accurate. In addition to providing numerical procedures of up to ninth-order sectional polynomial accuracy, the solutions are optimized in a least squares sense in their respective approximation spaces. Third,

the method produces general solutions of differential equations as well as particular ones. With these general solutions, it is possible to solve a differential equation separately in many regions using relatively small matrices, the connecting inter-element conditions being added subsequently. Further, unlike most **ordinary** projective methods where the boundary conditions are located at the edges of the region considered, general solutions allow boundary conditions to be specified anywhere.

As for the development of the automatic differential equation solving programs in this thesis, these too have been highly successful. As shown by the examples in Chapter 6, the DECL computer programs provide a practical alternative for solving many ordinary and two-dimensional partial differential equations to established methods. The program is sufficiently accurate for most engineering applications and their computational times are also acceptable, considering the general nature of the programs. Most importantly, however, the programs have a simple input and output structure that allows many differential equations to be solved with a minimum of analysis and data preparation.

It must be mentioned, however, that although the theory and the computer programs in this thesis have many attractive features, they also have several serious shortcomings. On the theoretical side, the most prominent of these is the neglect of this thesis to provide a rigorous error bound for the solutions

generated. As a result, in those cases where the accuracy of the solution mustable defined precisely, the user is required to examine the convergence of the solution using previously established procedures. In addition to requiring a further knowledge of numerical analysis, these procedures often magnify the computation required to treat a differential equation by an order of magnitude.

Some of the other mathematically unavoidable deficiencies in the method are the large sizes of matrices which result with high-order, two-dimensional elements, the failure of Newton's method to converge in all cases, and the necessity of avoiding singular points in the calculations. Fortunately, however, while these problems are troublesome, they do not preclude alternative formulations of the differential equation to avoid the occurence of these problems.

The DECL computer programs also have some limitations which are not based on theoretical reasons. Among these program limitations are the cumbersome technique used to specify two-dimensional boundary conditions, the lack of double precision versions of the programs, and the necessity of evaluating the quasi-linear form of a nonlinear differential equation by hand, instead of having the DECL compiler do it automatically. In terms of future development, these areas indicate some of the features of the DECL computer programs and of the DECL computer language which should be improved. Hopefully, in a continuing evolutionary development, newer, more

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sophisticated procedures will eventually replace the inadequate ones existing presently in DECL. Thus, while the DECL computer programs given in this thesis are insufficient for some purposes, they do constitute an important first step in the development of a general computer language for the automatic solution of differential equations.

Even in their present form, however, the overall performance of the DECL programs in treating most ordinary and two-dimensional partial differential equations accurately and efficiently indicates that these programs can be a valuable tool in engineering applications. By taking advantage of the ability of these programs to solve a large variety of problems without modification, it is often possible for practicing engineers to bypass the difficulty and expense of writing a special computer program to solve each separate type of differential equation. In this way, the valuable skills of many engineers who out of necessity are presently part-time numerical analysts and programmers can be deployed more efficiently.

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